



Modern Methods in Heterogeneous Catalysis Research

Electron Paramagnetic Resonance

November 21st, 2014

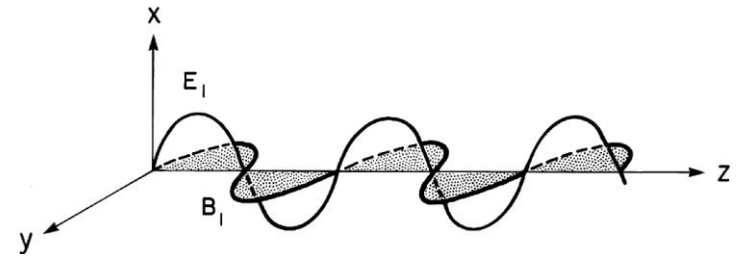
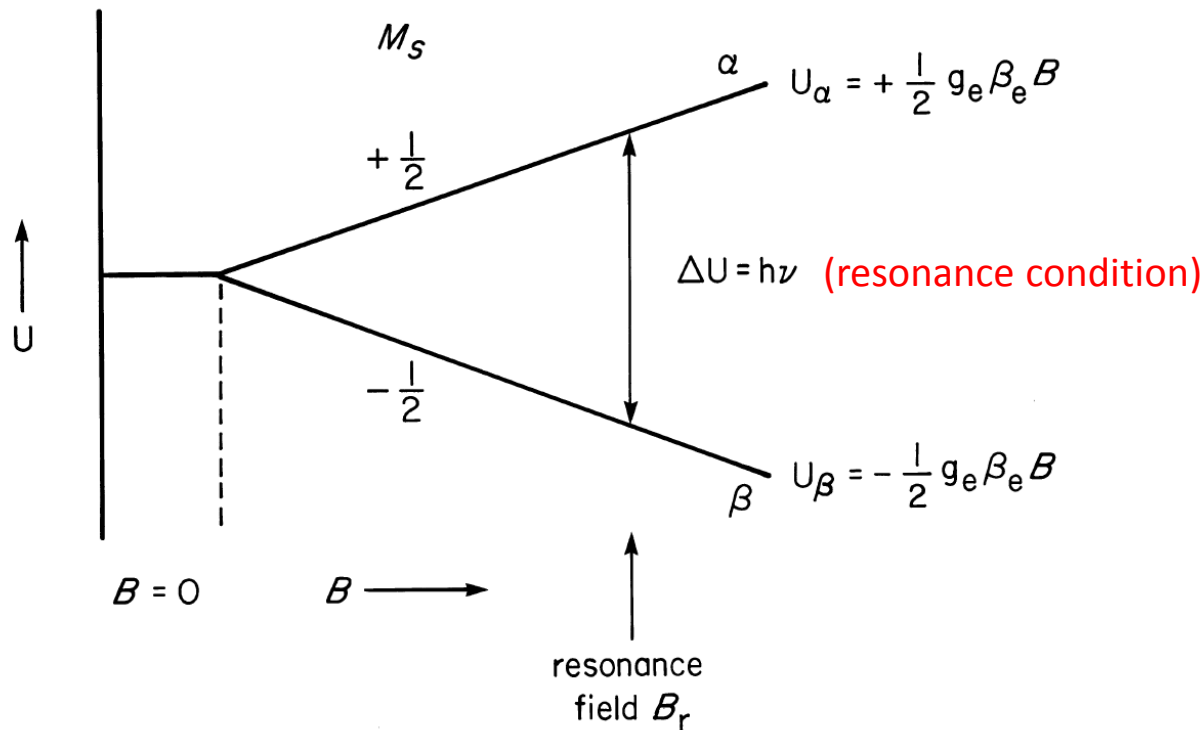
Outline

1. Basic Principles
2. Electron-Nucleus Interactions (Hyperfine Coupling)
3. Anisotropy
4. Electron-Electron Interactions (Fine Coupling)
5. Linewidths
6. Literature

1. Basic principles

- **Electron paramagnetic resonance (EPR) = Electron spin resonance (ESR)** spectroscopy
- Same underlying physical principles as in **nuclear magnetic resonance (NMR)**

One unpaired (free) electron:



Zeeman effect:

Selection rule: $\Delta M_S = \pm 1$

$$\Delta U = g\beta_e B$$

$$g = \frac{h\nu}{\beta_e B}$$

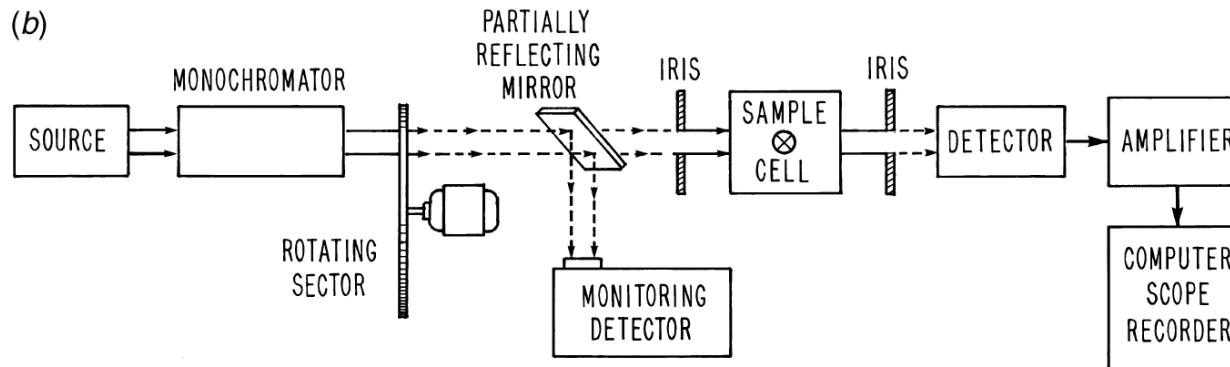
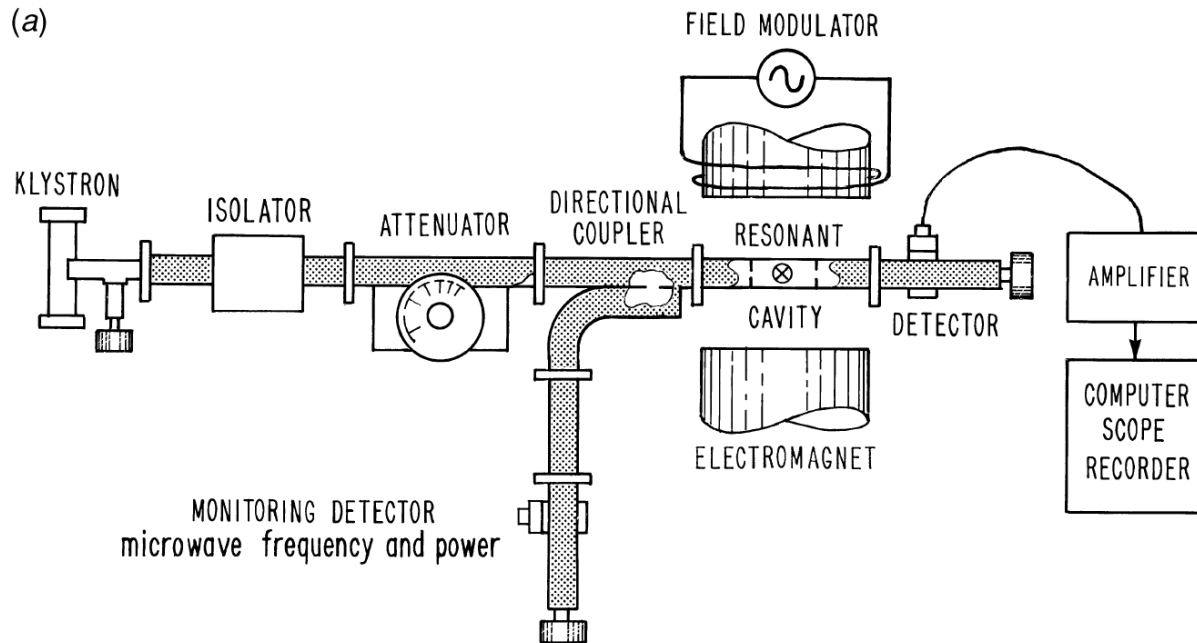
g : g factor

for free electron: $g_e = 2.0023$

β_e : Bohr magneton

1. Basic principles

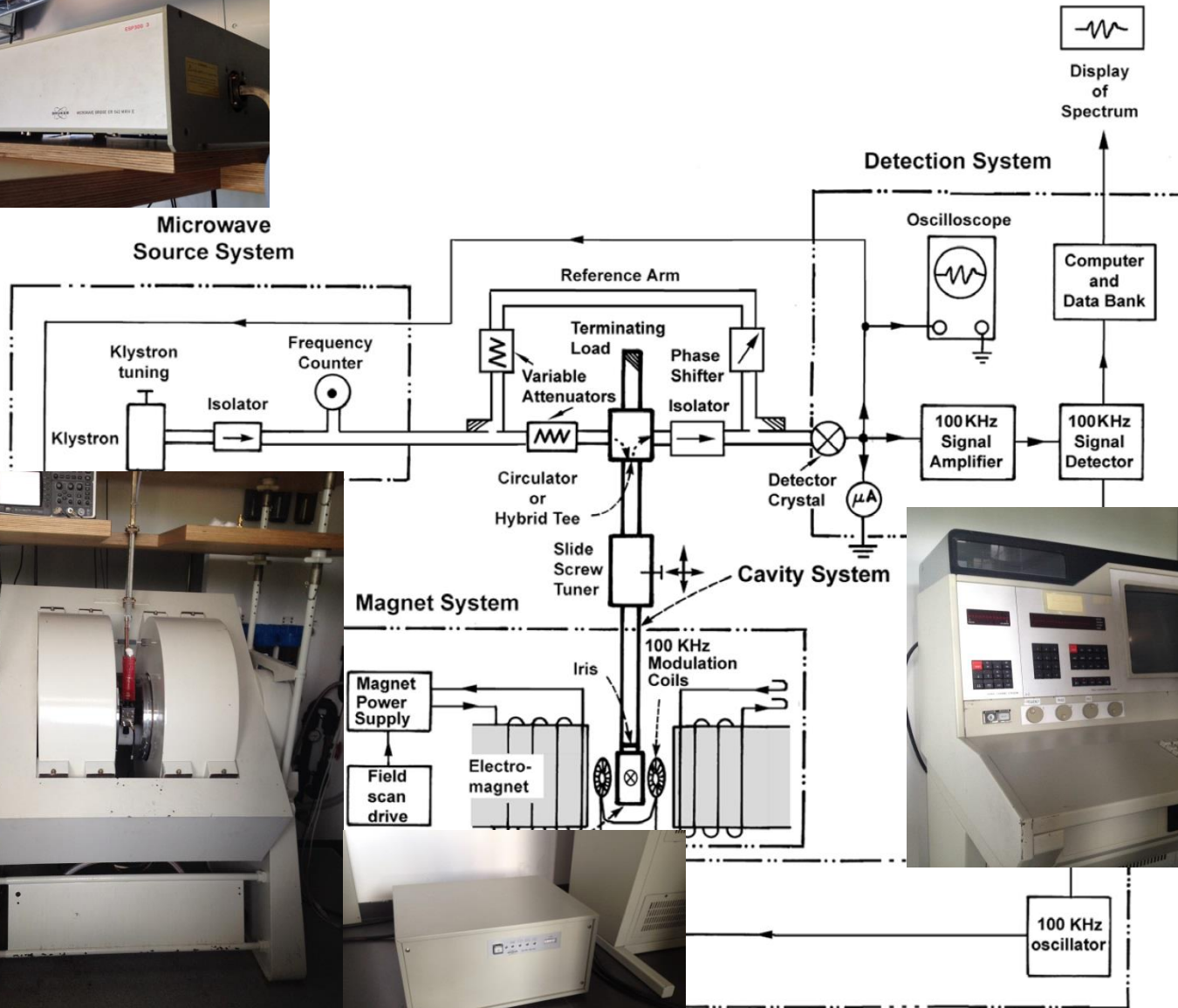
➤ A continuous wave (cw) EPR spectrometer:



1. Basic principles



Microwave Source System



Magnet System

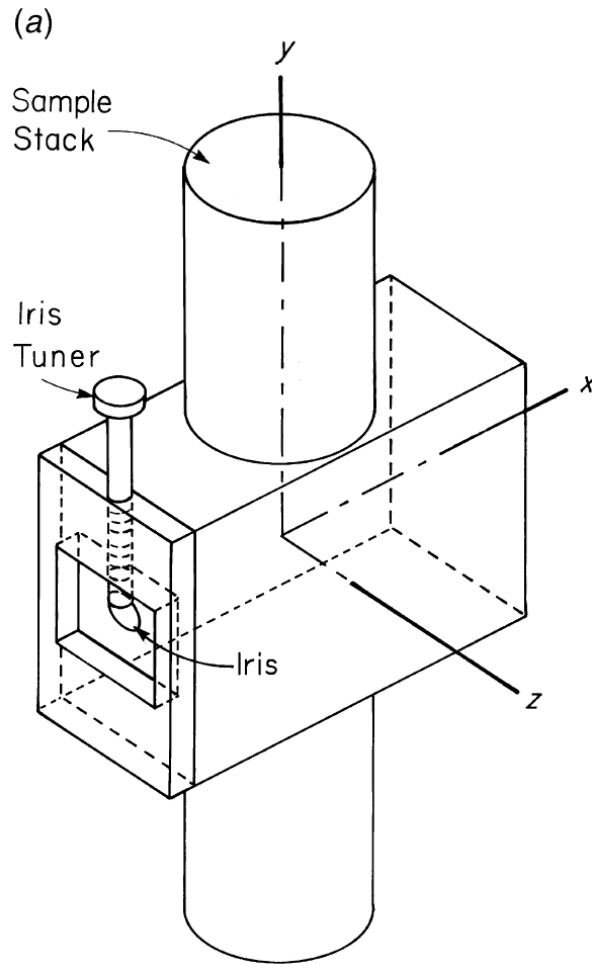


Field Modulation System

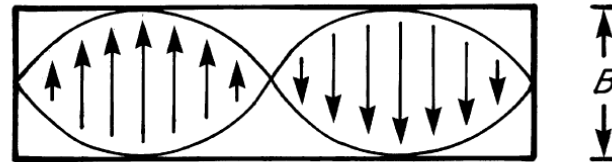
1. Basic principles

➤ The EPR resonator (cavity)

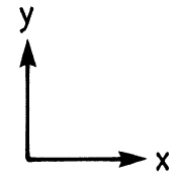
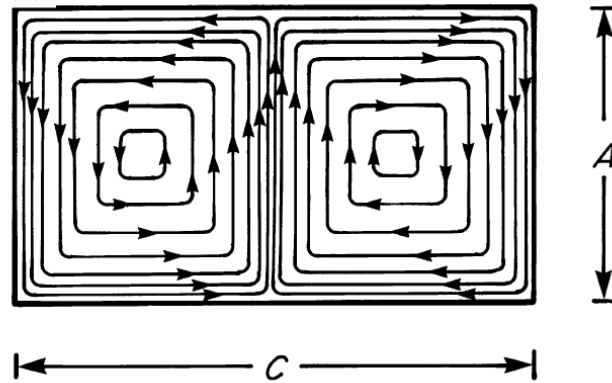
Rectangular TE_{102} cavity



(b) Electric field:



(c) Magnetic field:



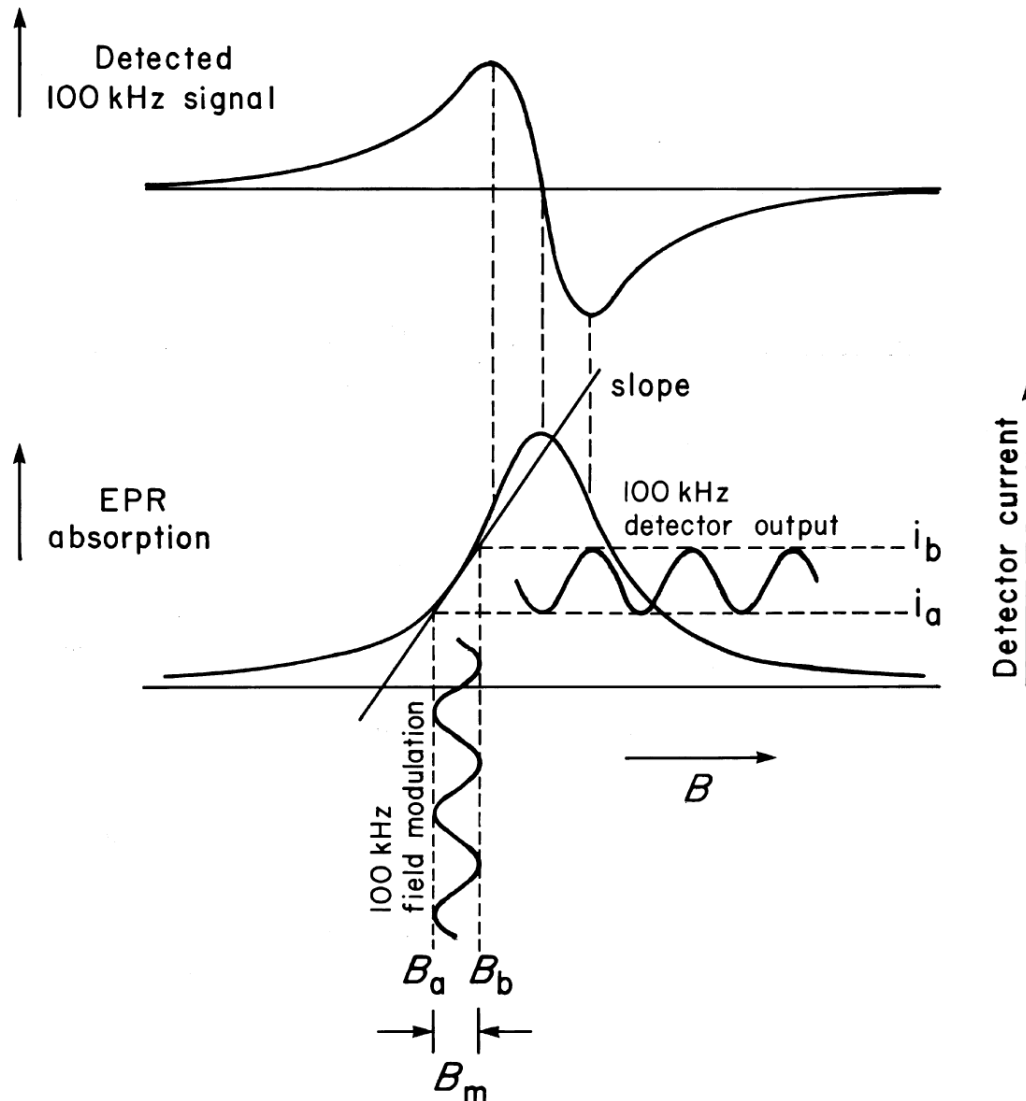
1. Basic principles

➤ Typical frequency and magnetic induction ranges in EPR

Microwave band	Microwave frequency (GHz)	B_0 (for $g = 2$) gauss
L	1	390
S	3	1070
X	9	3380
K	24	8560
Q	35	12,480
W	94	33,600

1. Basic principles

➤ (Phase) Sensitive detection by field modulation



1. Basic principles

➤ Samples that can be principally measured by EPR:

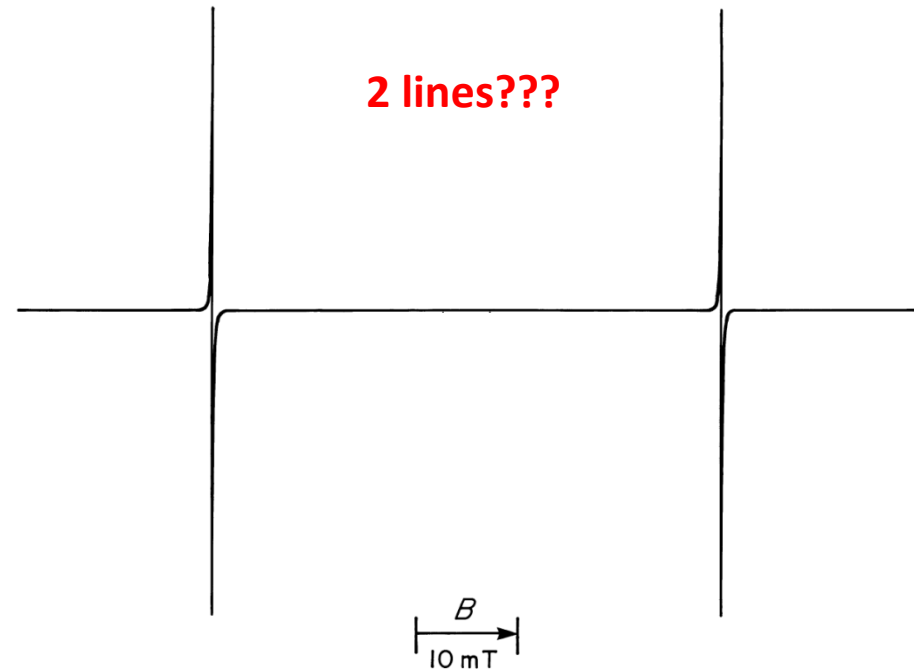
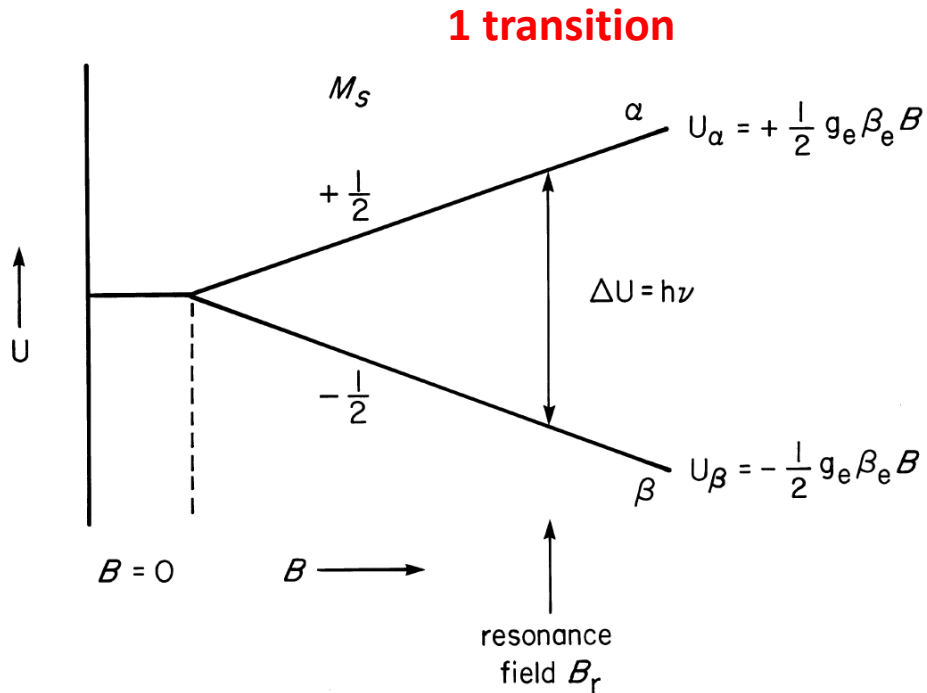
- ✓ Free radicals in solids, liquids or in the gas phase
- ✓ Transition metal ions with unpaired electron(s)
- ✓ Point defects in solids
- ✓ Systems with more than one unpaired electrons, e.g. triplet systems, biradicals, multiradicals
- ✓ Systems that temporarily generate states with unpaired electrons by excitation with, e.g., light
- ✓ Systems with conducting electrons

2. Electron-Nucleus Interactions

2. Electron-Nucleus Interactions

Example: Hydrogen atom

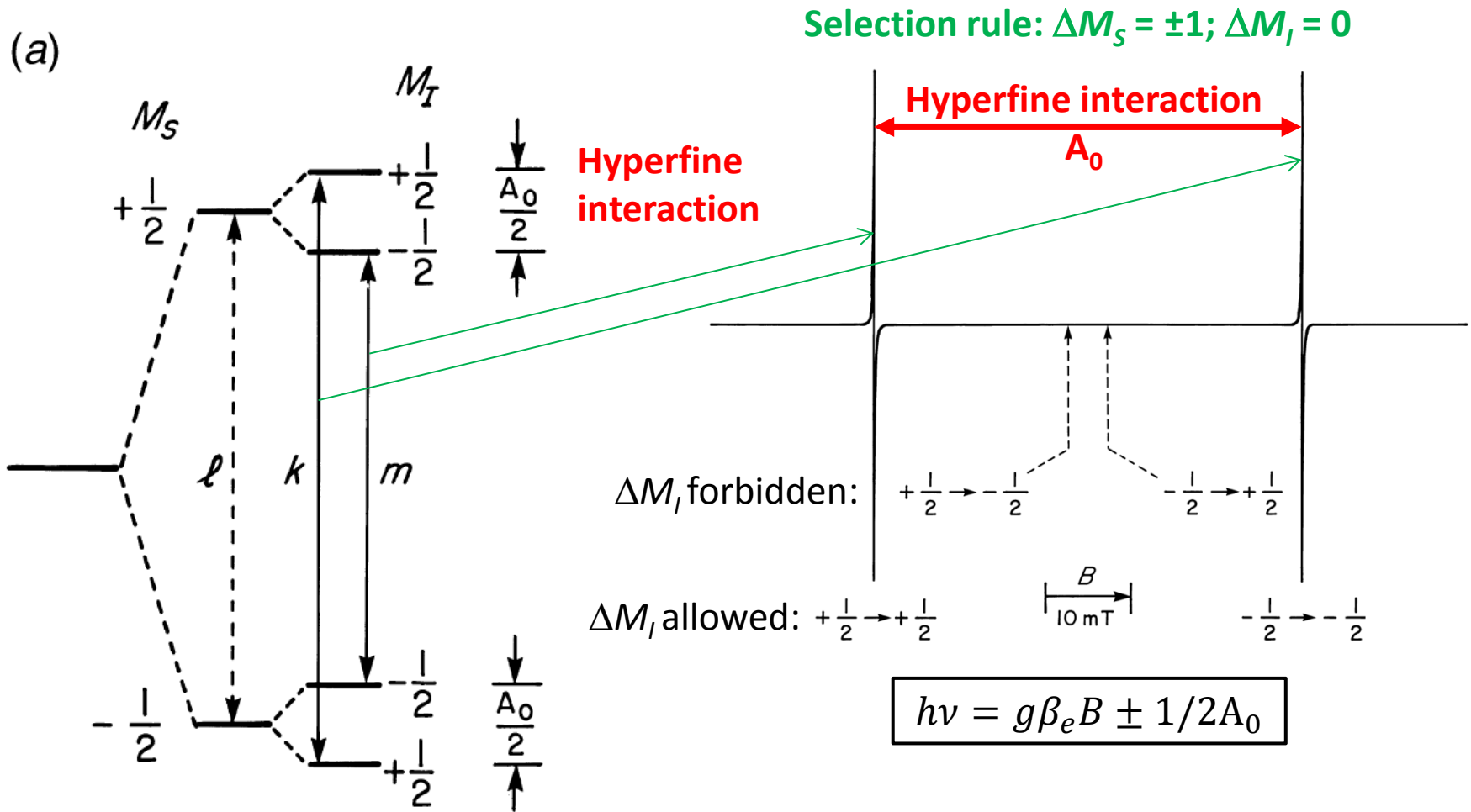
➤ One unpaired electron → electron spin $S = \frac{1}{2}$



2. Electron-Nucleus Interactions

Example: Hydrogen atom

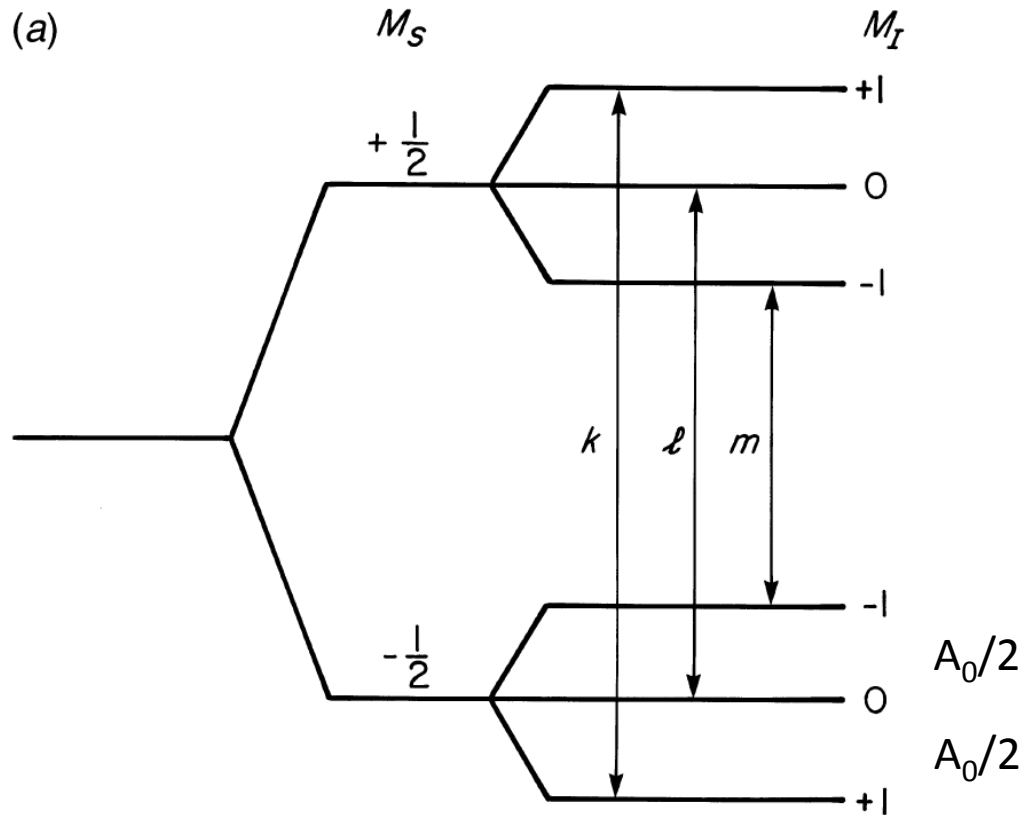
- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- H atom has a **nuclear spin**: $I = \frac{1}{2}$, $M_I = \pm\frac{1}{2}$



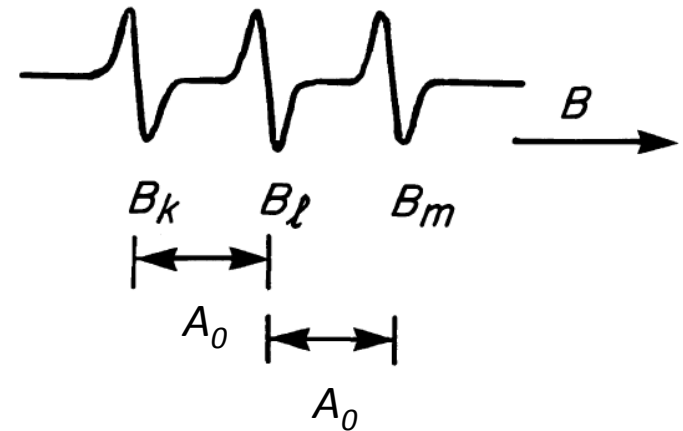
2. Electron-Nucleus Interactions

Example: Deuterium atom

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- Nuclear spin: $I = 1, M_I = -1, 0, +1$



Selection rule: $\Delta M_S = \pm 1; \Delta M_I = 0$

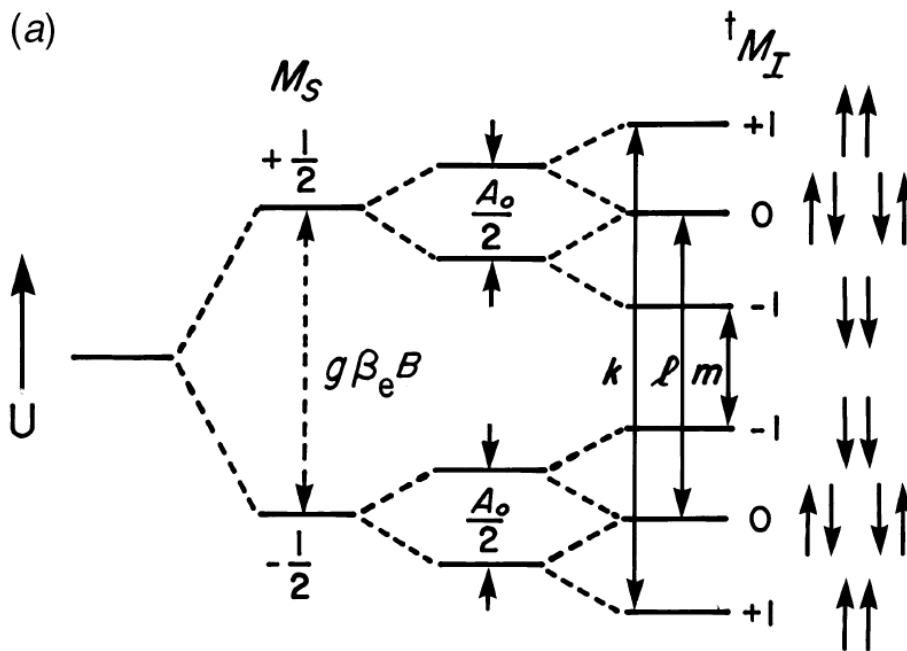


In general for single nucleus with spin I interacting with one electron: $2I+1$ lines of equal intensity, separated by hyperfine splitting A_0

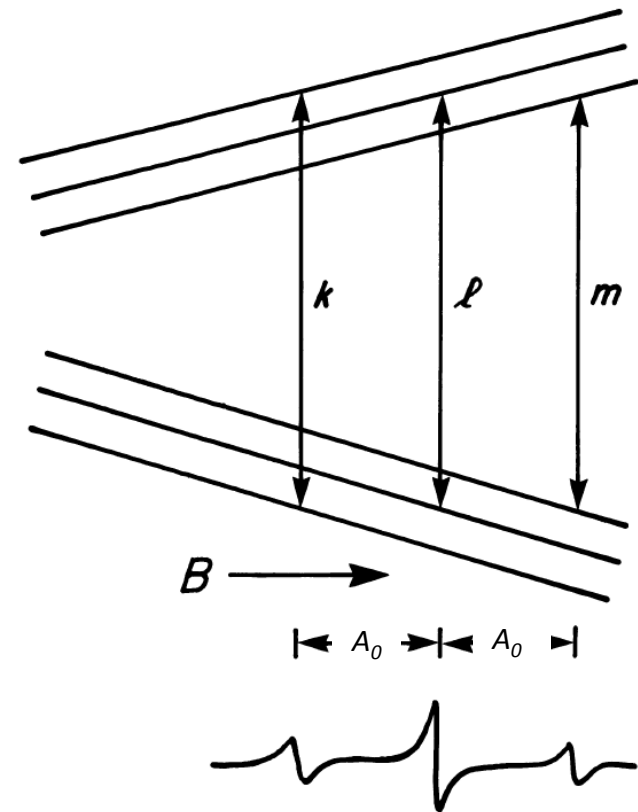
2. Electron-Nucleus Interactions

Interaction of multiple nuclei with one electron

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- Two equivalent nuclei with $I = \frac{1}{2}$, $M_I = -\frac{1}{2}, +\frac{1}{2}$



(b) Selection rule: $\Delta M_S = \pm 1$; $\Delta M_I = 0$



2. Electron-Nucleus Interactions

Interaction of multiple nuclei with one electron

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- n equivalent nuclei with

(a) $I_i = 1/2 \quad (i = 1, 2, \dots, n)$				(b) $I_i = 1 \quad (i = 1, 2, \dots, n)$			
No. of nuclei n	No. of lines $2I + 1$	First-order line positions/ a_0 M_I	Binomial intensity ratios	No. of nuclei n	No. of lines $2I + 1$	First-order line positions/ a_0 M_I	Multinomial intensity ratios
0	1	0	1	0	1	0	1
1	2	$\frac{1}{2} \quad \frac{1}{2}$	1 1	1	3	$\bar{1} \quad 0 \quad 1$	1 1 1
2	3	$\bar{1} \quad 0 \quad 1$	1 2 1	2	5	$\bar{2} \quad \bar{1} \quad 0 \quad 1 \quad 2$	1 2 3 2 1
3	4	$\frac{3}{2} \quad \frac{1}{2} \quad \frac{1}{2} \quad \frac{3}{2}$	1 3 3 1	3	7	$\bar{3} \quad \bar{2} \quad \bar{1} \quad 0 \quad 1 \quad 2 \quad 3$	1 3 6 7 6 3 1
4	5	$\bar{2} \quad \bar{1} \quad 0 \quad 1 \quad 2$	1 4 6 4 1	4	9	$\bar{4} \quad \bar{3} \quad \bar{2} \quad \bar{1} \quad 0 \quad 1 \quad 2 \quad 3 \quad 4$	1 4 10 16 19 16 10 4 1
5	6	$\frac{5}{2} \quad \frac{3}{2} \quad \frac{1}{2} \quad \frac{1}{2} \quad \frac{3}{2} \quad \frac{5}{2}$	1 5 10 10 5 1				
6	7	$\bar{3} \quad \bar{2} \quad \bar{1} \quad 0 \quad 1 \quad 2 \quad 3$	1 6 15 20 15 6 1				
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮

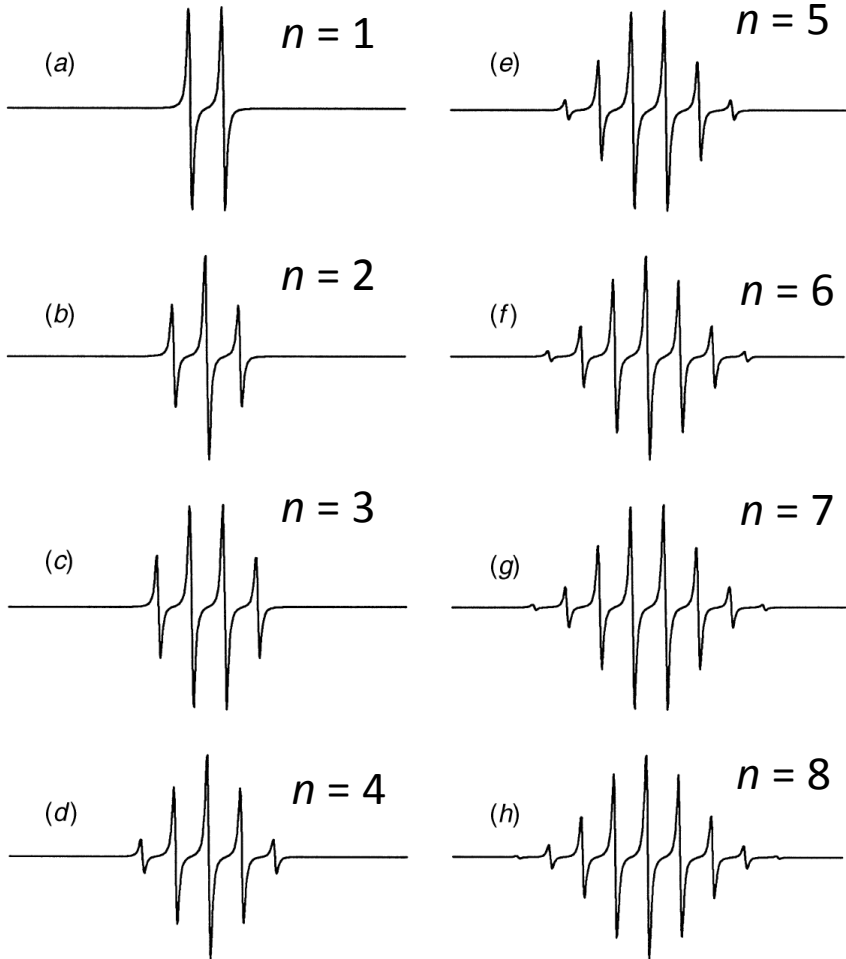
In general for n equivalent nuclei with spin I interacting with one electron: $2nI+1$ lines with multinomial intensity ratios ("Pascal's triangle"), separated by hyperfine splitting A_0

2. Electron-Nucleus Interactions

Interaction of multiple nuclei with one electron

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- n equivalent nuclei with $I = \frac{1}{2}$

$$2nI+1$$



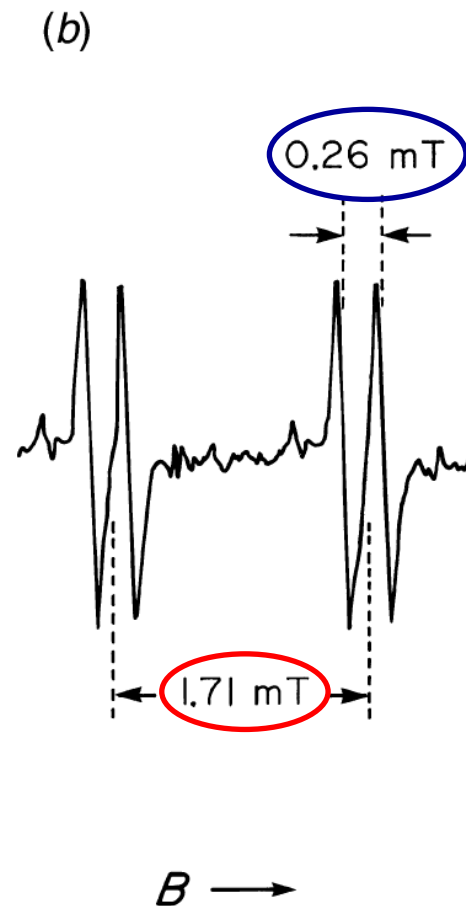
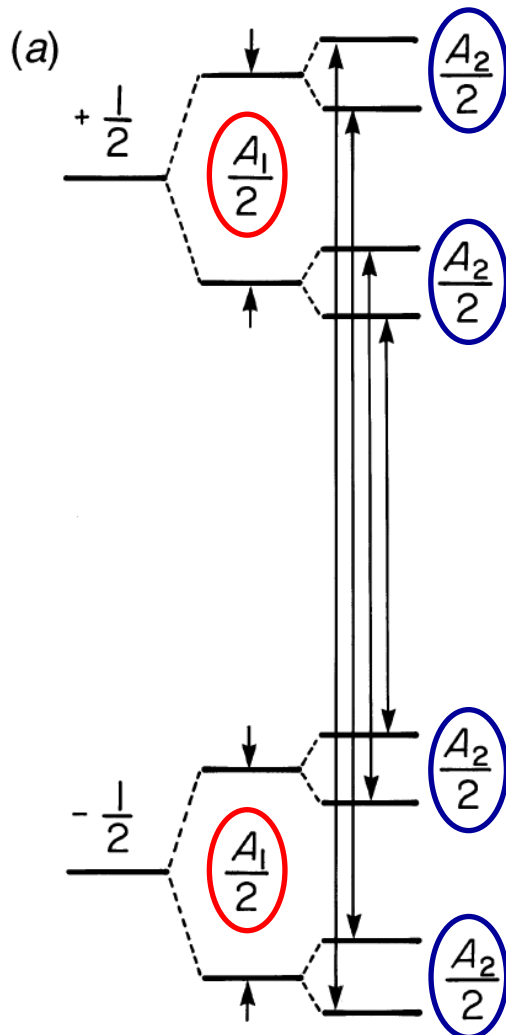
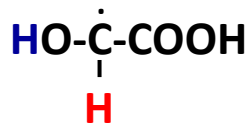
$B \rightarrow$



2. Electron-Nucleus Interactions

Interaction of multiple nuclei with one electron

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$
- Two **inequivalent** nuclei with $I = \frac{1}{2}$



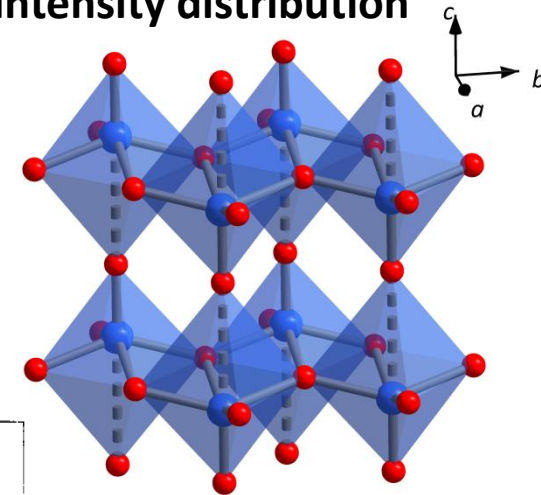
2. Electron-Nucleus Interactions

EPR spectrum of V_2O_5 single crystal

- One unpaired electron \rightarrow electron spin $S = \frac{1}{2}$ (V^{4+} defect)
- n (?) equivalent nuclei with $I = 7/2$

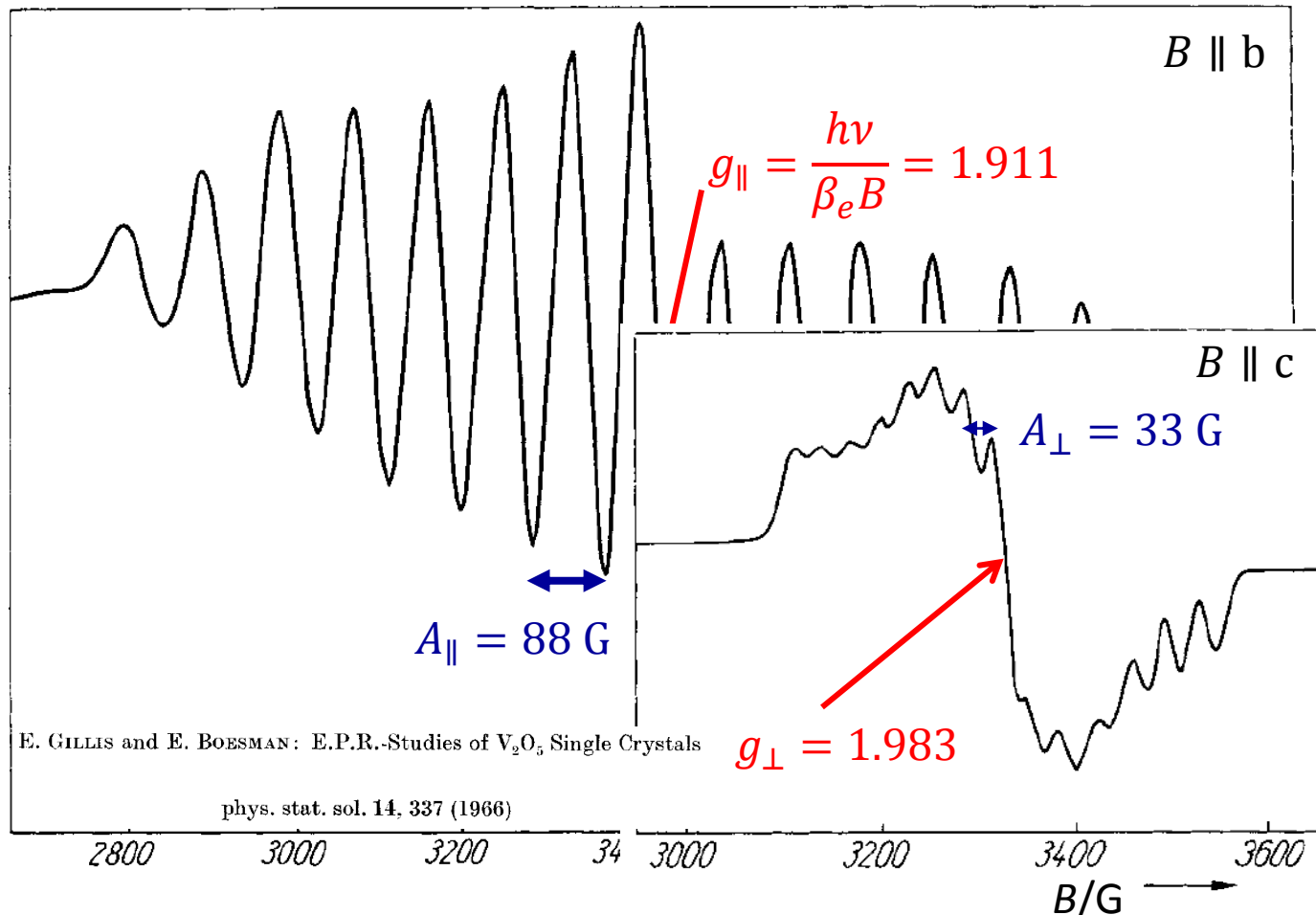
Number of lines = $2nI+1 = 2 \times 2 \times 7/2 + 1 = 15$ equally spaced lines with intensity distribution

1 : 2 : 3 : 4 : 5 : 6 : 7 : 8 : 7 : 6 : 5 : 4 : 3 : 2 : 1



Orthorhombic space group:
 $a=3.564(2) \text{ \AA}$ $b=11.519(6) \text{ \AA}$
 $c=4.373(2) \text{ \AA}$

Conclusion:
 electron localized
 at two vanadium
 atoms



3. Anisotropy

3. Anisotropy

Electron in environment with:

- Cubic symmetry (cubic, octahedral, tetrahedral coordination)
- (Uni)Axial Symmetry
- Orthorhombic Symmetry

Spin Hamiltonian:

$$\hat{H} = \beta_e (g_x B_x \hat{S}_x + g_y B_y \hat{S}_y + g_z B_z \hat{S}_z) = \beta_e [B_x \quad B_y \quad B_z] \cdot \begin{matrix} \text{Row vector} \\ \begin{bmatrix} g_x & 0 & 0 \\ 0 & g_y & 0 \\ 0 & 0 & g_z \end{bmatrix} \end{matrix} \cdot \begin{matrix} \text{column} \\ \text{vector} \\ \begin{bmatrix} \hat{S}_x \\ \hat{S}_y \\ \hat{S}_z \end{bmatrix} \end{matrix} = \beta_e \mathbf{B}^T \cdot \mathbf{g} \cdot \hat{\mathbf{S}}$$

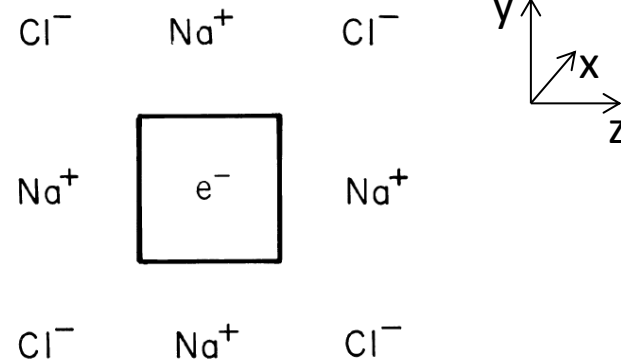
➤ **Cubic Symmetry:**

Isotropic g factor (independent on magnetic field direction): $g_x = g_y = g_z$

→ g is a scalar constant

$$\hat{H} = \beta_e g (B_x \hat{S}_x + B_y \hat{S}_y + B_z \hat{S}_z)$$

F center in NaCl:



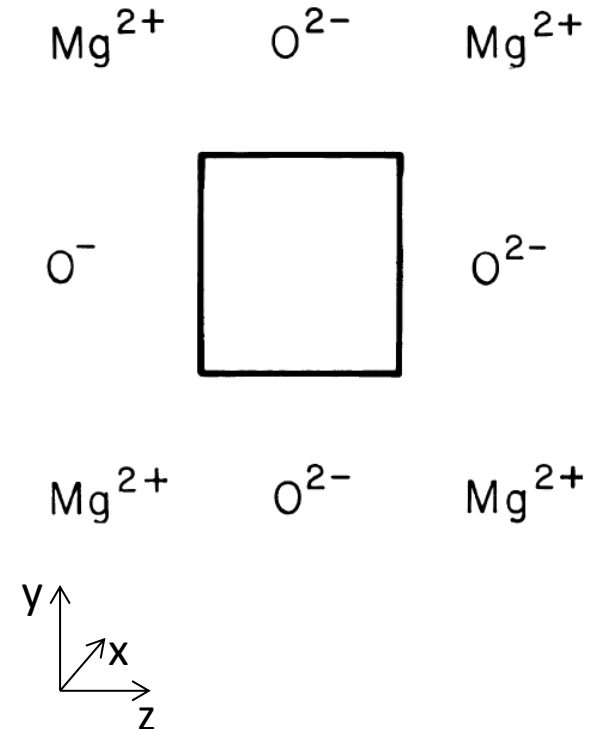
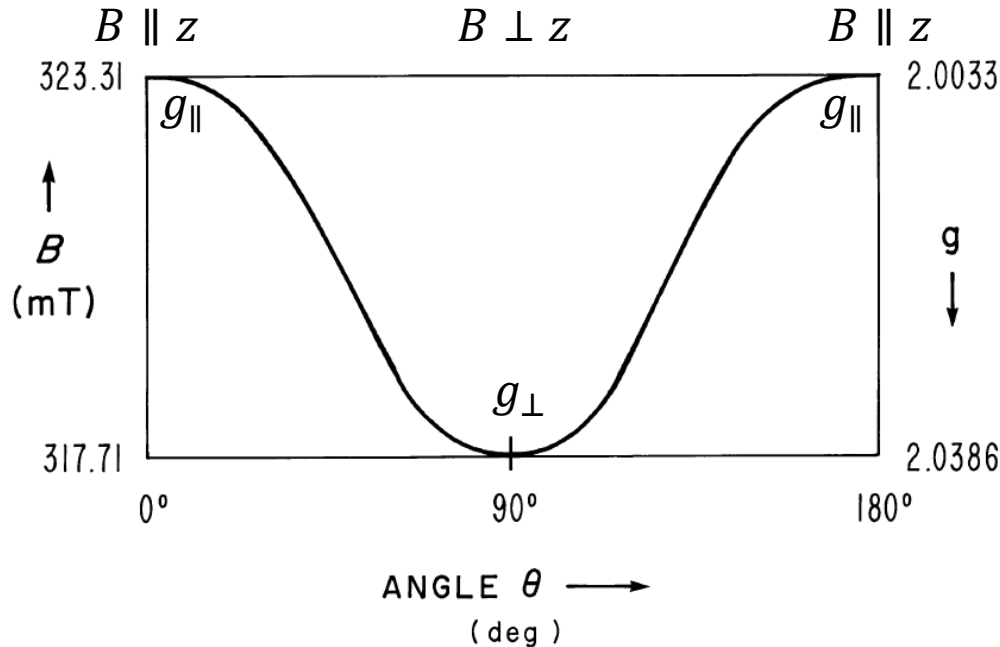
3. Anisotropy

➤ (Uni)Axial Symmetry

Anisotropic g factor (dependent on magnetic field direction): $g_x = g_y = g_{\perp} \neq g_z = g_{\parallel}$

Spin Hamiltonian: $\hat{H} = \beta_e(g_x B_x \hat{S}_x + g_y B_y \hat{S}_y + g_z B_z \hat{S}_z)$
 $= \beta_e[g_{\perp}(B_x \hat{S}_x + B_y \hat{S}_y) + g_{\parallel} B_z \hat{S}_z]$

V center in MgO

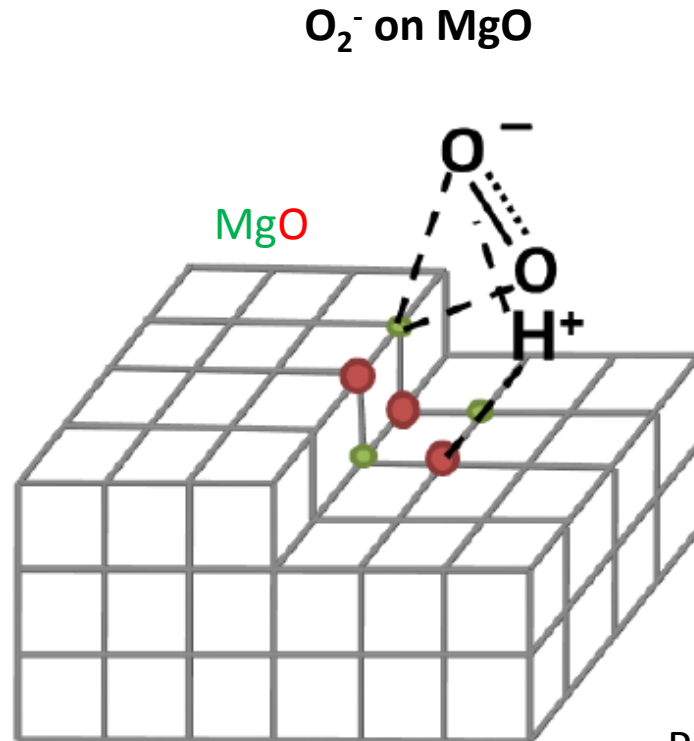


3. Anisotropy

➤ Orthorhombic Symmetry

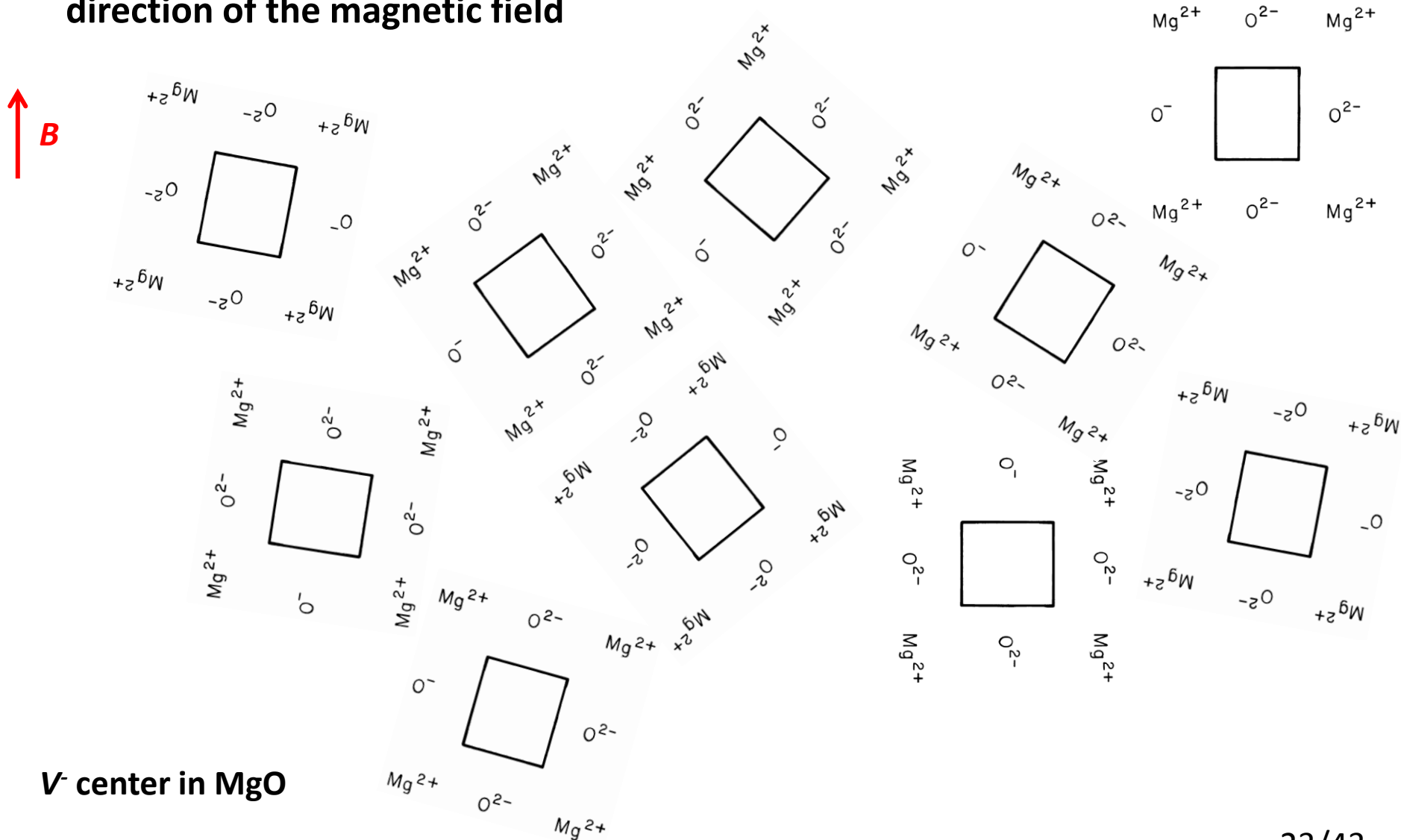
Spin Hamiltonian: $\hat{H} = \beta_e (g_x B_x \hat{S}_x + g_y B_y \hat{S}_y + g_z B_z \hat{S}_z)$

Anisotropic g factor (dependent on magnetic field direction): $g_x \neq g_y \neq g_z$



3. Anisotropy

- Crystalline **powders**: principal axis has all possible orientations relative to the direction of the magnetic field



3. Anisotropy

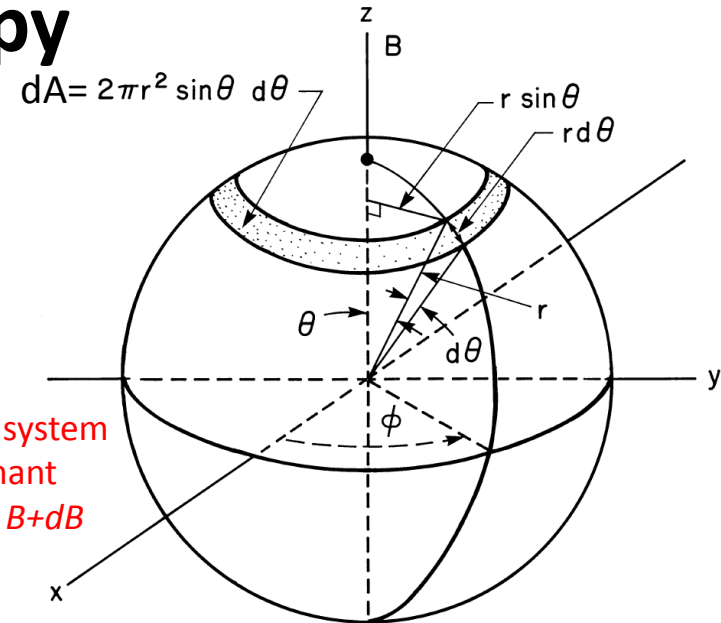
- Crystalline powder
- $S = \frac{1}{2}, I = 0$
- Uniaxial local symmetry

Fraction of symmetry axes between θ and $\theta+d\theta$

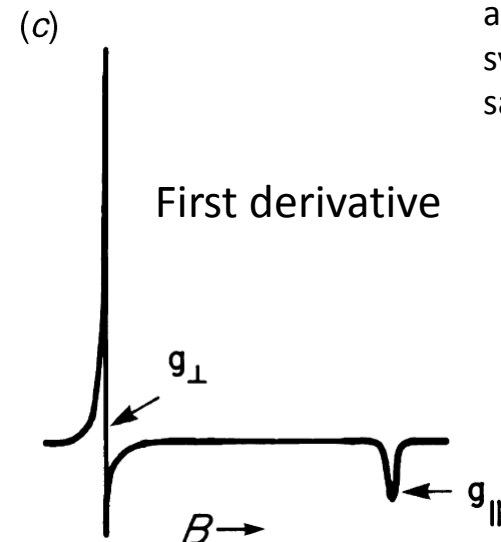
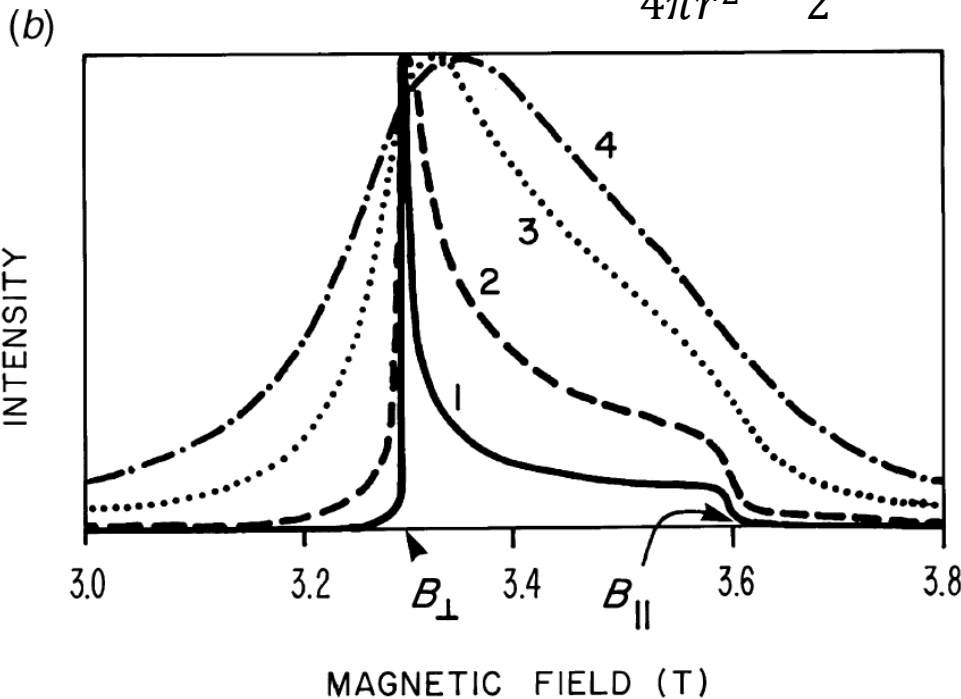
$$P(\theta)d\theta = \frac{dA}{A_{sphere}} \sim P(B)dB$$

Probability of a spin system experiencing a resonant field between B and $B+dB$

$$P(\theta)d\theta = \frac{dA}{4\pi r^2} = \frac{1}{2} \sin\theta d\theta$$

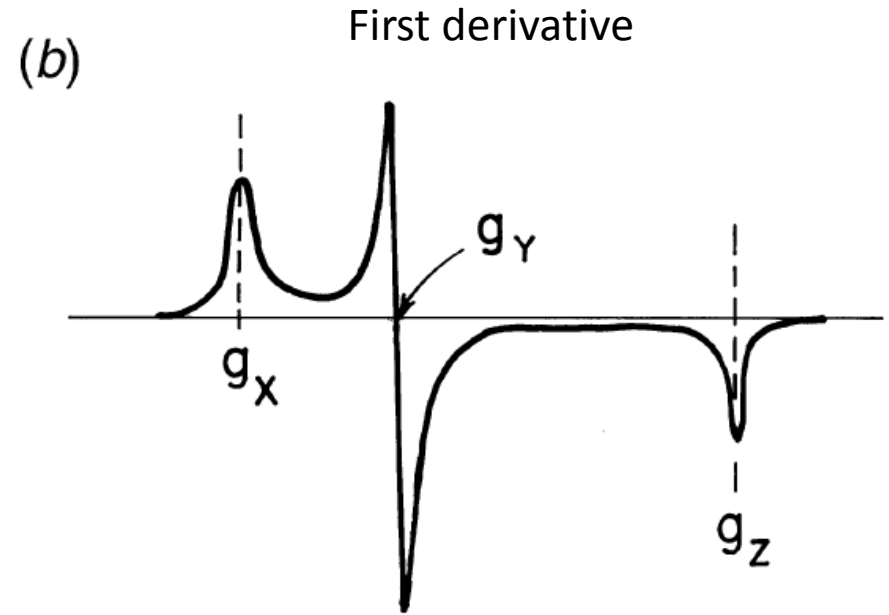
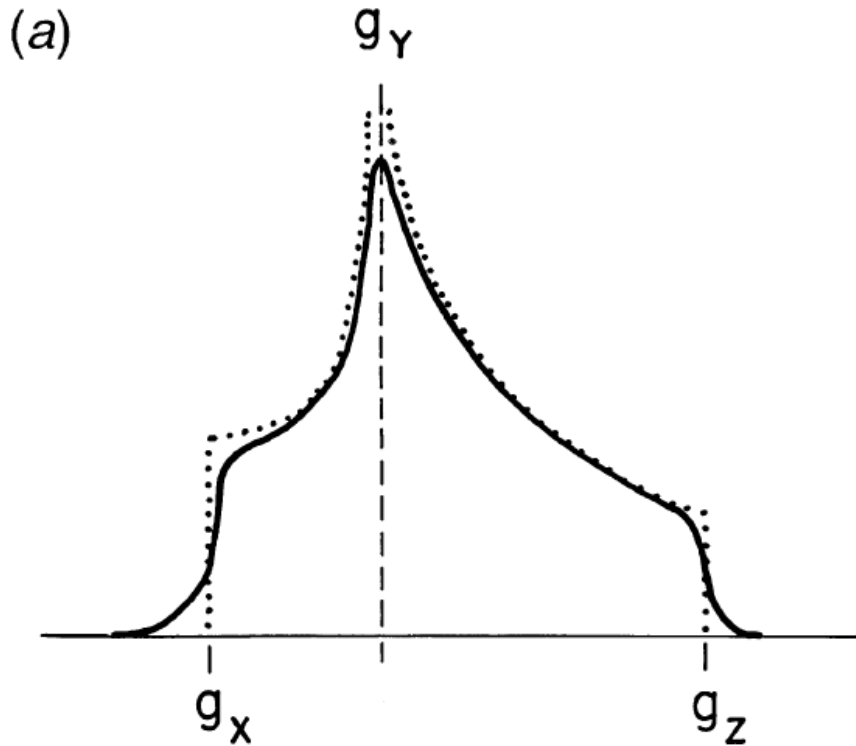


θ : Angle between magnetic field and principal symmetry axis of any spin system in the sample



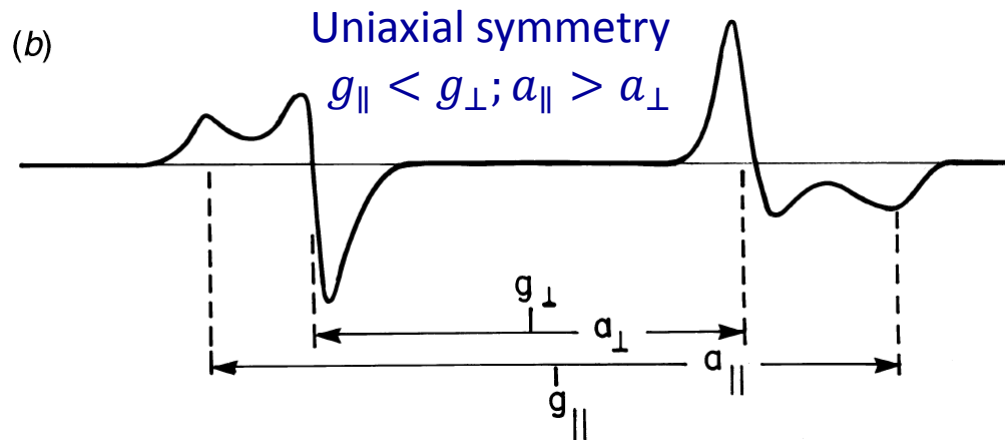
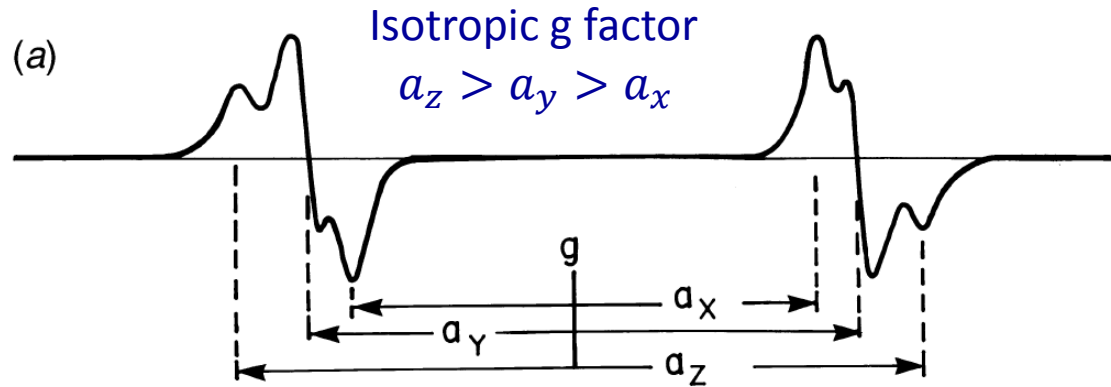
3. Anisotropy

- Crystalline powder
- $S = \frac{1}{2}, I = 0$
- Orthorhombic local symmetry



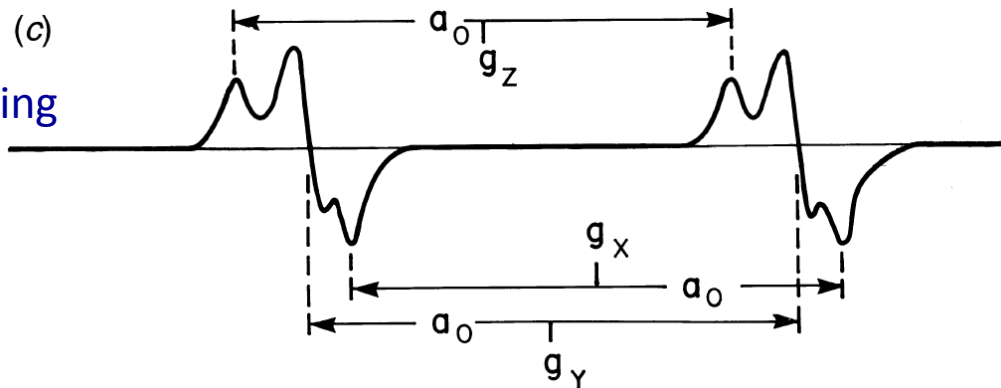
3. Anisotropy

- Crystalline powder
- $S = \frac{1}{2}, I = \frac{1}{2}$
- Hyperfine anisotropy



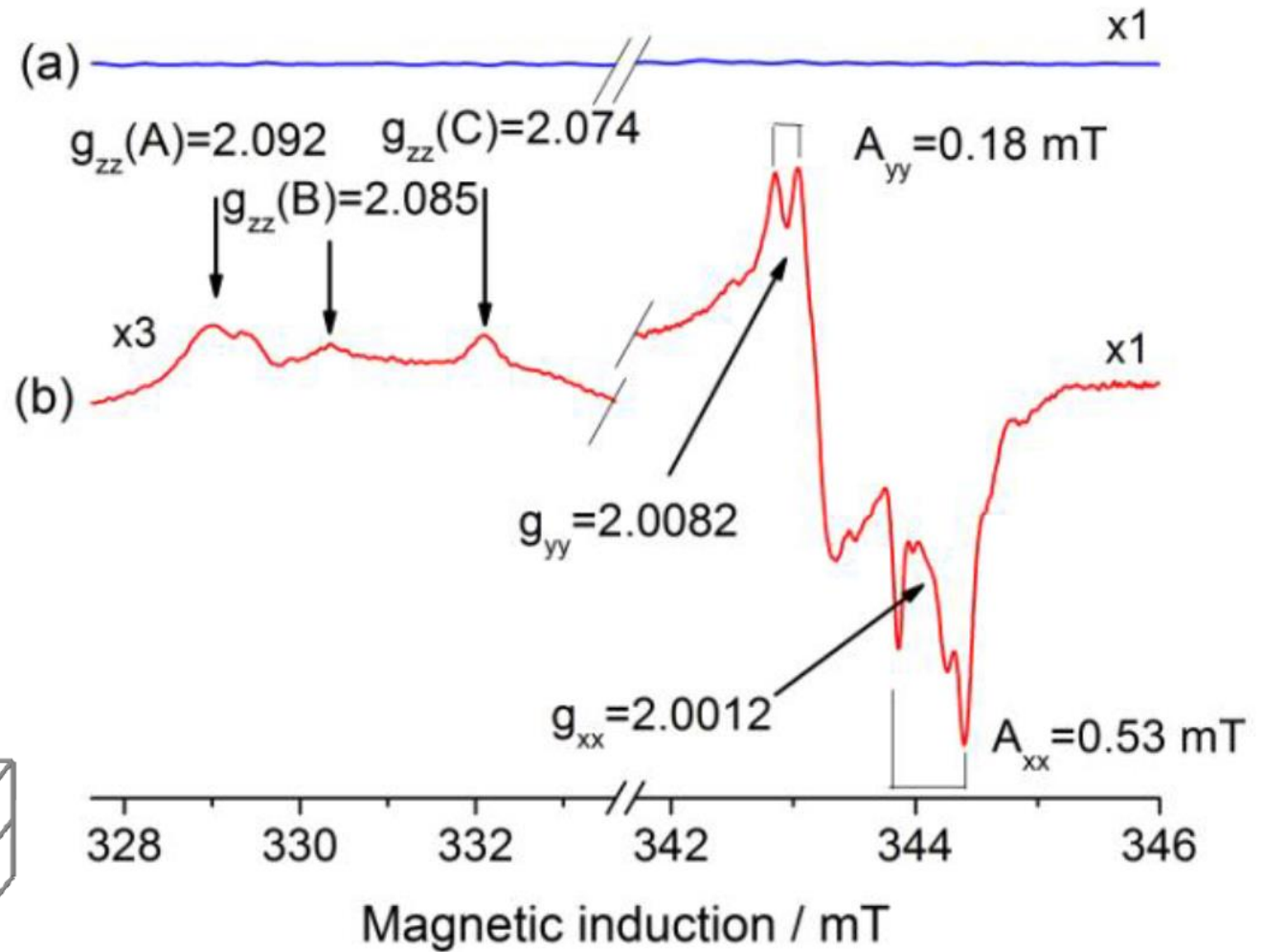
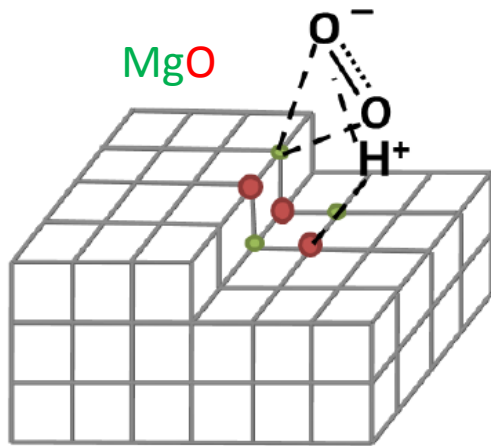
Isotropic hyperfine splitting

$$g_z > g_y > g_x$$



3. Anisotropy

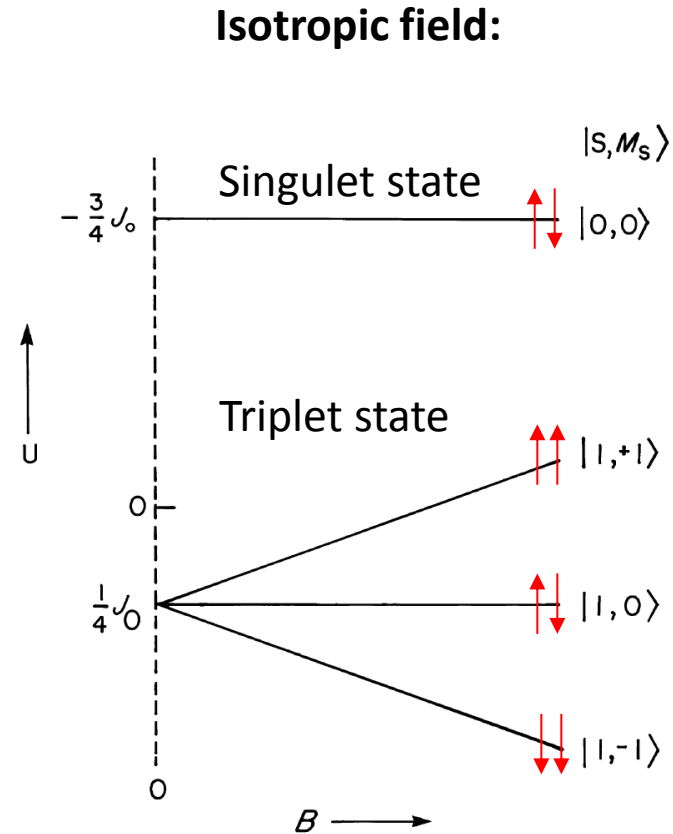
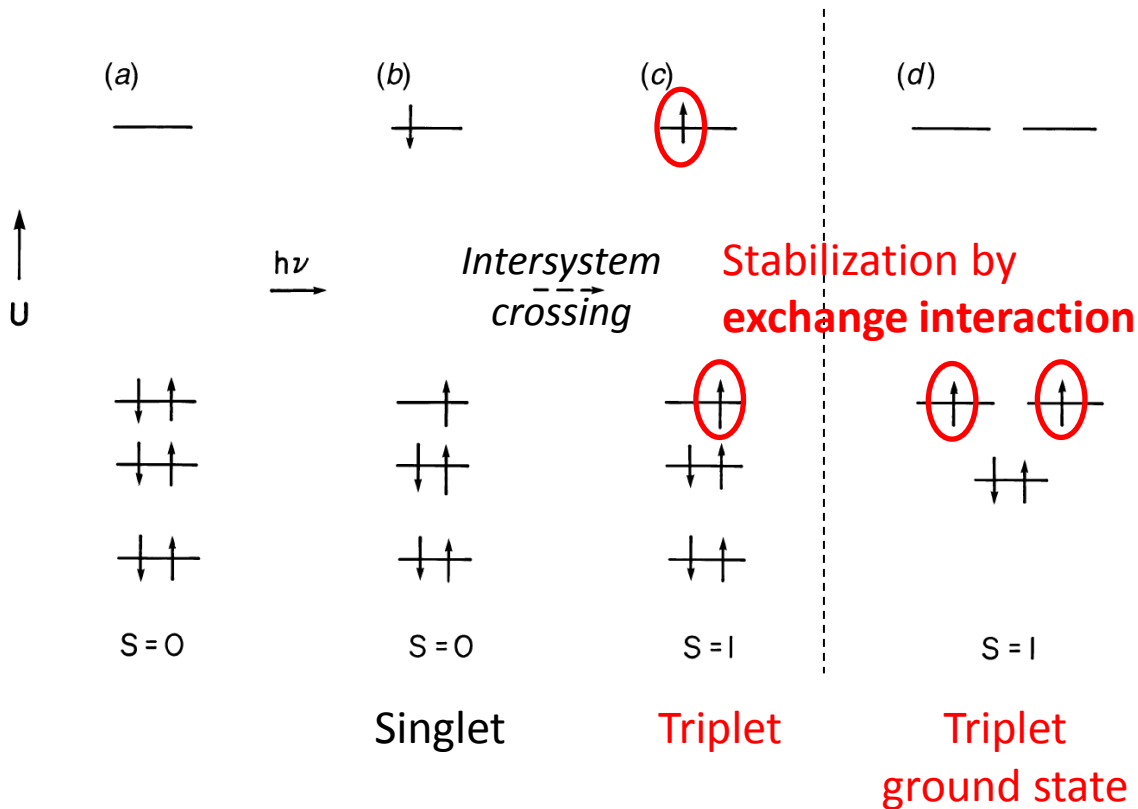
- MgO: H^+/O_2^-
- $S = \frac{1}{2}$, $I(\text{H}) = \frac{1}{2}$
- Orthorhombic:



4. Electron-Electron Interactions

4. Electron-Electron Interactions

- E.g. O_2 , V^{3+} , Ni^{2+} , Fe^{3+}
- Electron exchange interaction



4. Electron-Electron Interactions

➤ Electron-electron dipole interaction

Only Zeeman effect ($S = \frac{1}{2}$, isotropic g factor, $l = 0$): $\hat{H} = g\beta_e(B_x\hat{S}_x + B_y\hat{S}_y + B_z\hat{S}_z)$

With electron-electron dipole (anisotropic) interaction (fine coupling) ($S = 1$, isotropic g factor, $l = 0$):

$$\hat{H} = g\beta_e(B_x\hat{S}_x + B_y\hat{S}_y + B_z\hat{S}_z) + D_x\hat{S}_x^2 + D_y\hat{S}_y^2 + D_z\hat{S}_z^2$$

$$\hat{H} = g\beta_e\mathbf{B}^T \cdot \hat{\mathbf{S}} + D(\hat{S}_z^2 - \frac{1}{3}\hat{S}^2) + E(\hat{S}_x^2 - \hat{S}_y^2)$$

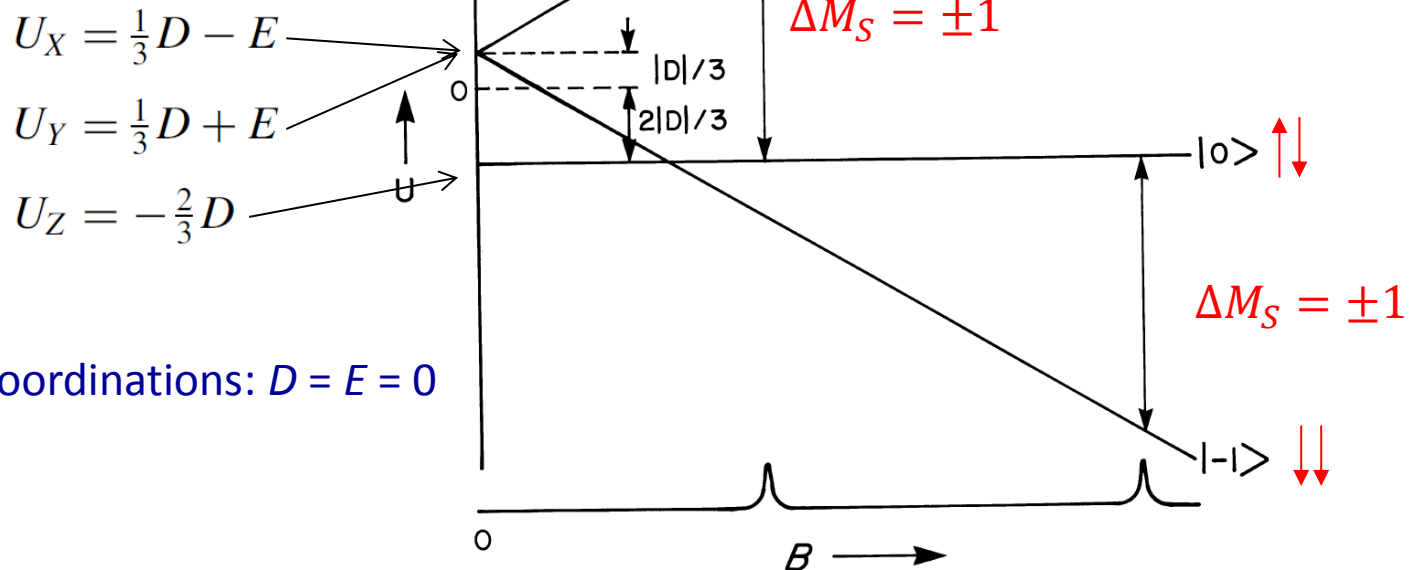
e.g. $E = 0$

Zero-field splitting

$$U_X = \frac{1}{3}D - E$$

$$U_Y = \frac{1}{3}D + E$$

$$U_Z = -\frac{2}{3}D$$



In cubic coordinations: $D = E = 0$

4. Electron-Electron Interactions

➤ Electron-electron dipole interaction

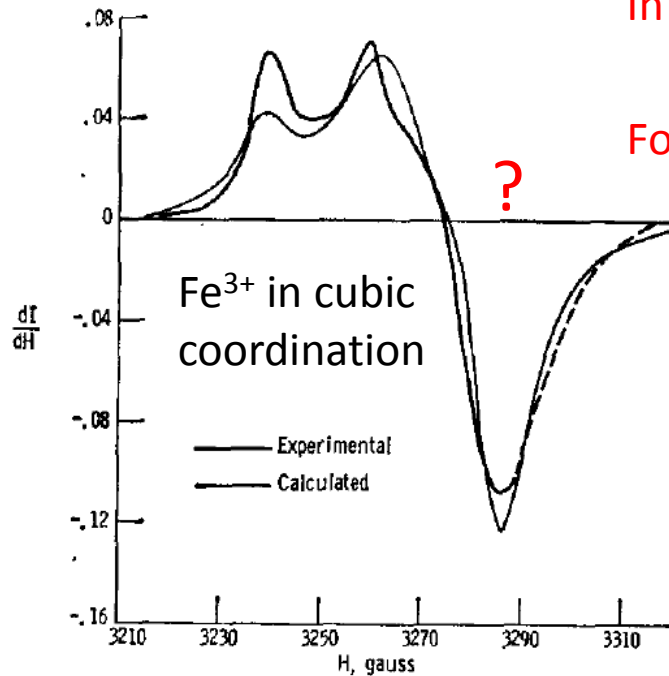
With electron-electron dipole interaction (fine coupling) ($S = 1$, isotropic g factor, $I = 0$):

Second order term $\sim S^2$

$$\hat{H} = g\beta_e \mathbf{B}^T \cdot \hat{\mathbf{S}} + D(\hat{S}_z^2 - \frac{1}{3}\hat{S}^2) + E(\hat{S}_x^2 - \hat{S}_y^2)$$

In cubic coordinations: $D = E = 0$

For high spin systems, e.g. Fe^{3+} : $S = 5/2$



$$\mathcal{H}_s = g\beta\mathbf{H} \cdot \mathbf{S}$$

$$+ \frac{1}{6}a[S_x^4 + S_y^4 + S_z^4 - \frac{1}{5}S(S+1)(3S^2 + 3S - 1)]$$

a : Fourth-order (high spin) parameter for cubic coordination

Polycrystalline EPR Spectrum of Fe^{3+} Ions in MgO

JACK H. LUNSFORD

Lewis Research Center, National Aeronautics and Space Administration, Cleveland, Ohio

(Received 22 June 1964)

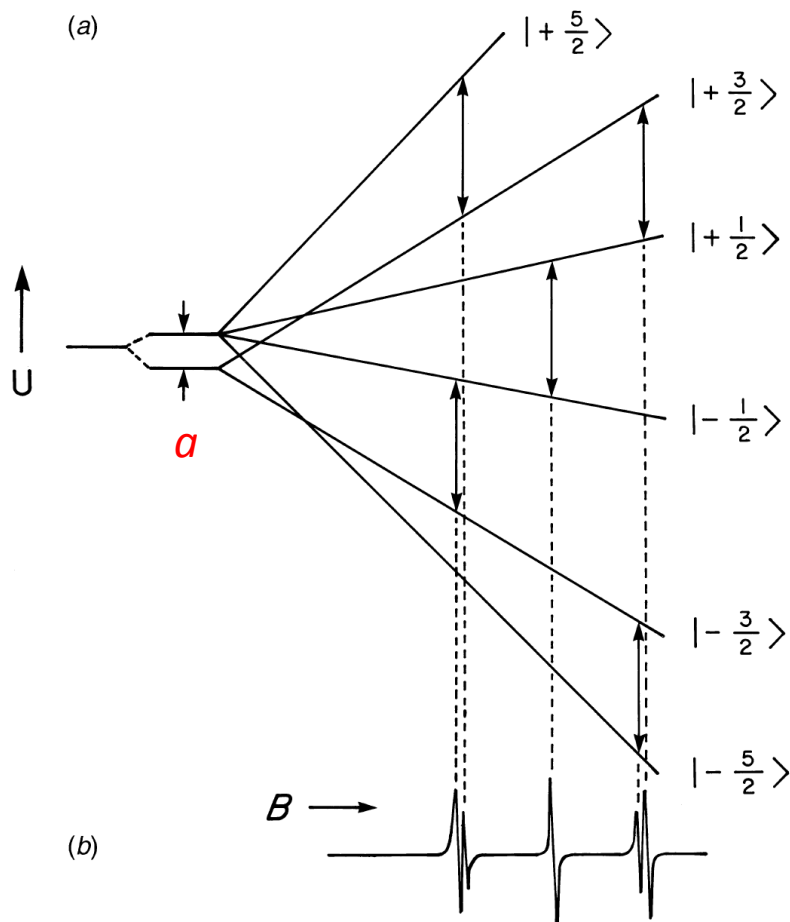
FIG. 1. Experimental and calculated derivative spectra for Fe^{3+} ions in polycrystalline MgO. The experimental curve is for Fisher reagent-grade MgO powder and is extrapolated in the region of an overlapping Cr^{3+} spectrum (dashed portion).

4. Electron-Electron Interactions

➤ Electron-electron dipole interaction

With electron-electron dipole interaction (fine coupling) ($S = 5/2$, isotropic g factor, $I = 0$):

Single crystal d^5 ion in octahedral (cubic) crystal field with B parallel to principal axis



$$\mathcal{H}_s = g\beta\mathbf{H} \cdot \mathbf{S} + \frac{1}{6}a[S_x^4 + S_y^4 + S_z^4 - \frac{1}{5}S(S+1)(3S^2 + 3S - 1)]$$

For (uni)axial symmetries: Axial fourth-order parameter **F**

➤ The complete (simplified) **Spin Hamiltonian**

$$\hat{\mathcal{H}} = g\beta_e \mathbf{B}^T \cdot \hat{\mathbf{S}} + \hat{\mathbf{S}}^T \cdot \mathbf{D} \cdot \hat{\mathbf{S}} + \hat{\mathbf{S}}^T \cdot \mathbf{A} \cdot \hat{\mathbf{I}}$$

Zeeman splitting: *g tensor*

Fine coupling (high spin
electron-electron
interactions): *D, E, a, F*

Nuclear hyperfine
coupling (electron-
nucleus
interactions): *A
tensor*

Spectrum Simulation with EasySpin®

Mn

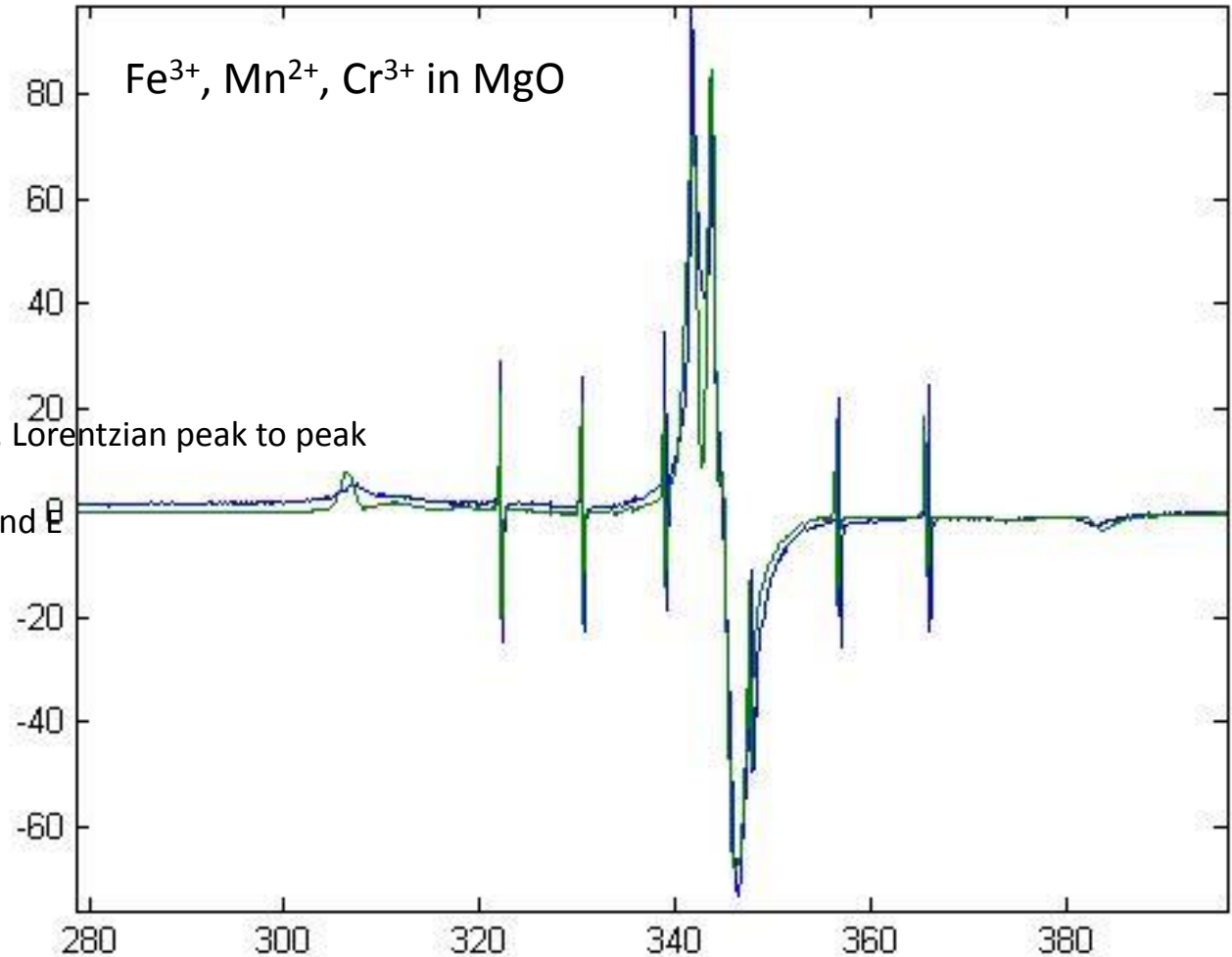
Sys1.S = 5/2;
Sys1.g = 2.0007;
Sys1.Nucs = '55Mn';
Sys1.A = -244;
Sys1.AStrain = 0;
Sys1.aF = [60 0];
Sys1.lwpp = 0.1;
Sys1.D = 90;
Sys1.DStrain = 120;
Sys1.weight = 100;

Fe

Sys2.S = 5/2;
Sys2.g = 2.0027; % isotropic g
Sys2.lwpp = [0.5 0.0]; % Gaussian, Lorentzian peak to peak
width, mT
Zero-field splitting in terms of D and E
Sys2.D = [120 0]; %in MHz
Sys2.DStrain = [600 0];
Sys2.aF=[650 0];
Sys2.aFStrain = [0 0]
Sys2.weight=16000

Cr

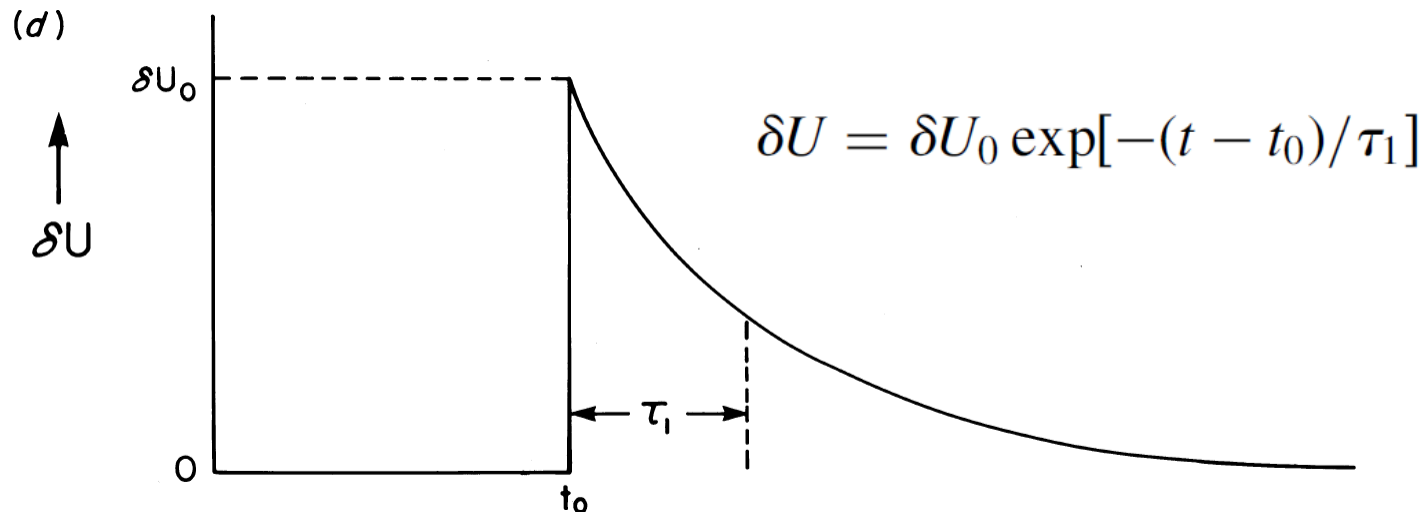
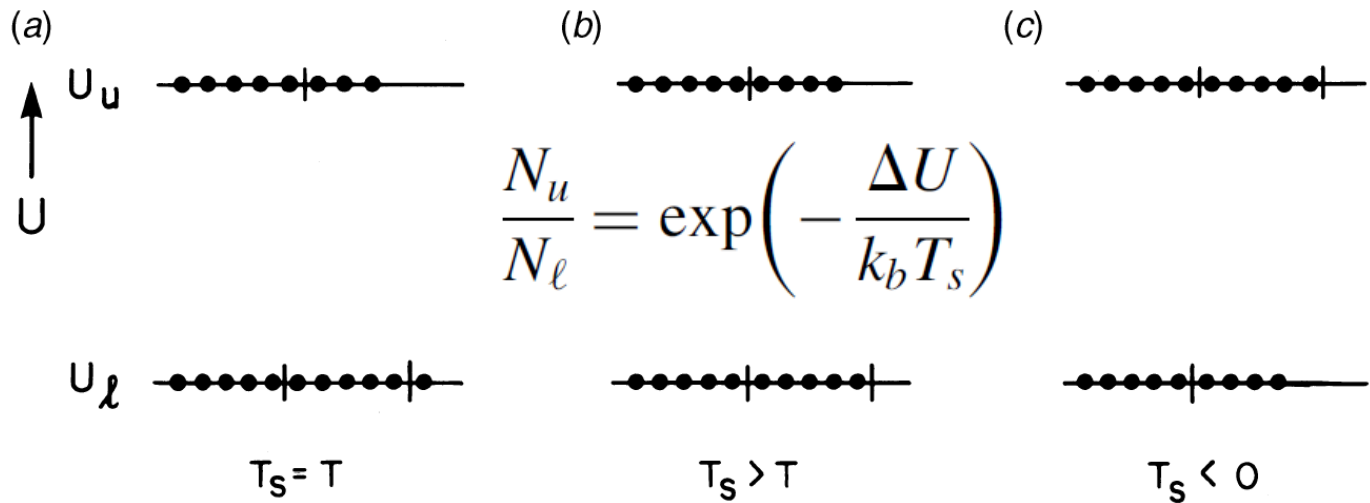
Sys3.S=3/2;
Sys3.g=1.98;
Sys3.Nucs='Cr';
Sys3.A=3;
Sys3.lwpp=0.2;
Sys3.weight=5;



5. Linewidths

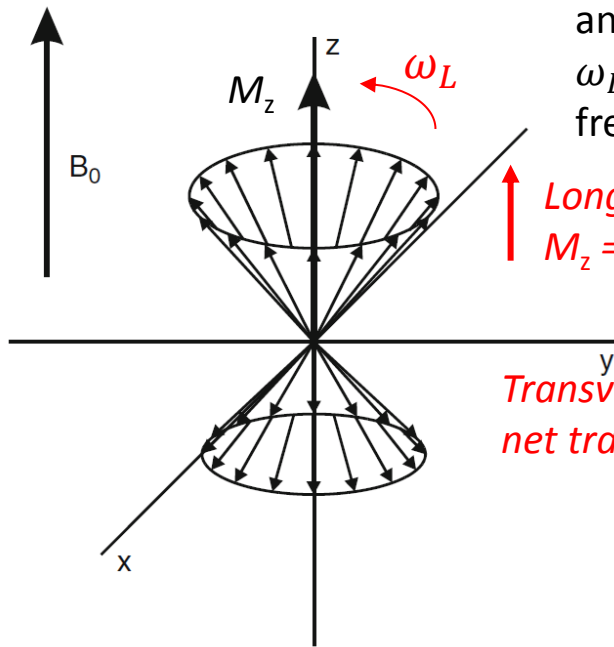
5. Linewidths

- **Spin relaxation: spin-lattice relaxation time τ_1** (spin interaction with surroundings, longitudinal relaxation time)



5. Linewidths

- **Spin relaxation: spin-spin relaxation time τ_2** (spin spin interaction, transversal relaxation time)



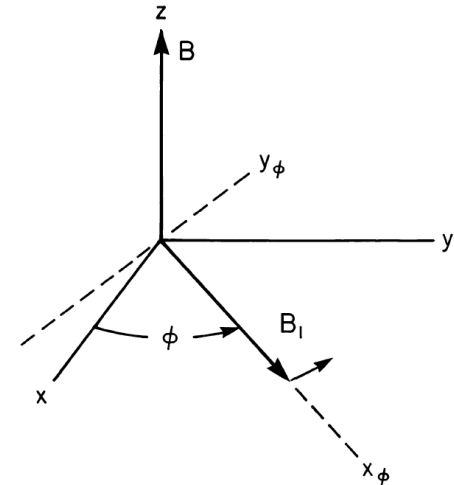
\mathbf{M} precesses about \mathbf{B}_0 with angular frequency $\omega_L = -\gamma_e \mathbf{B}_0$ (Larmor frequency)

Longitudinal magnetization $M_z = \text{const.}$

Transversal M_x, M_y oscillate, no net transversal magnetization

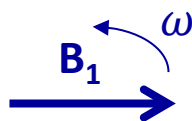
Net transversal magnetization (all spins rotate in phase with ω)

Frame rotates with the angular frequency ω

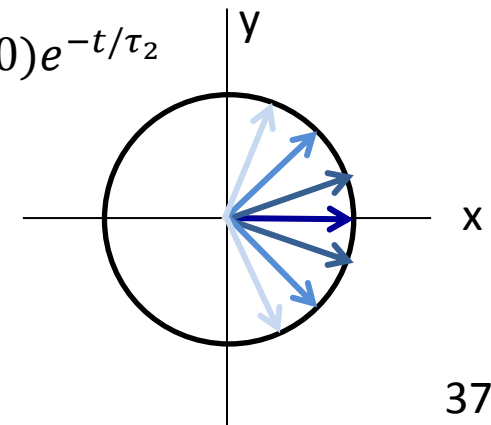
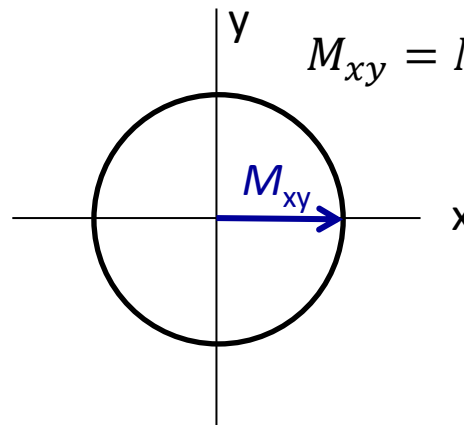


Transversal magnetization decays with τ_2

Superposition of a rotating perpendicular field B_1



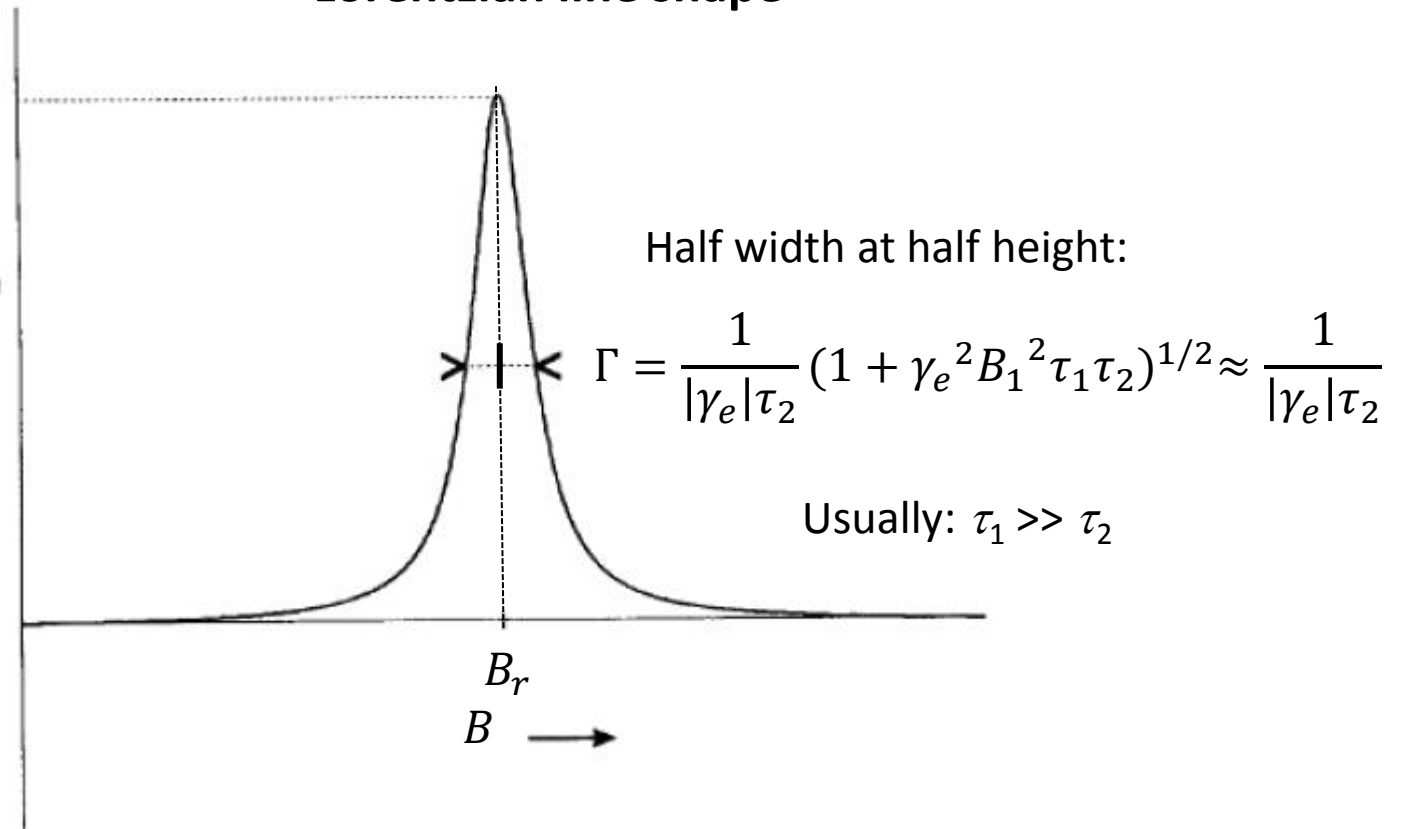
At resonance: $\omega = \omega_L$



5. Linewidths

- Linewidths: homogeneous broadening by τ_1 and τ_2

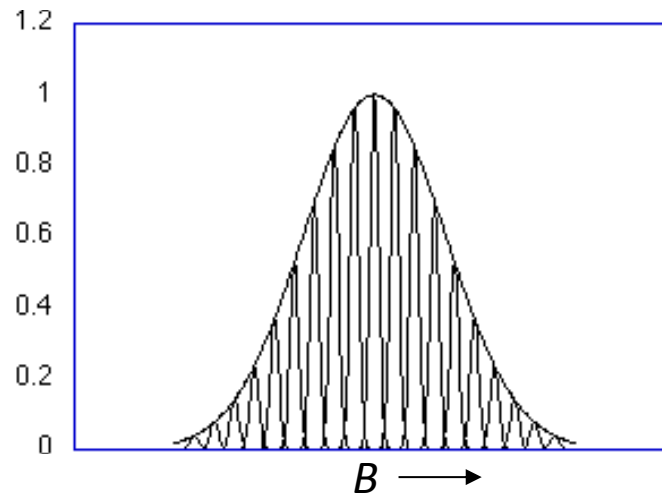
Lorentzian line shape



5. Linewidths

- **Linewidths: inhomogeneous broadening by superposition of spectra from individual equivalent spins**

Gaussian(-Lorentzian) line shape



Caused by

- An inhomogeneous external magnetic field
- Unresolved hyperfine structure
- Anisotropic interactions
- Dipolar interactions

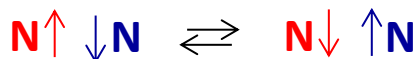
5. Linewidths

➤ Example for line broadening: electron-spin exchange



$^{14}\text{N}: I = 1 \rightarrow M = 2I + 1 = 3$ (Triplet)

in EtOH

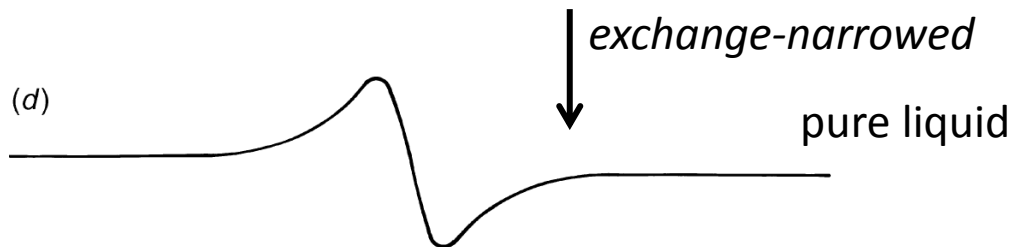
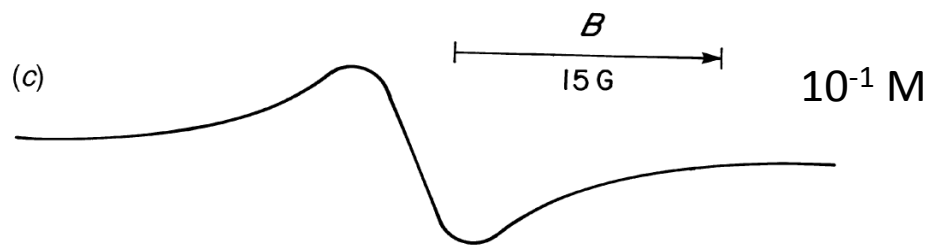
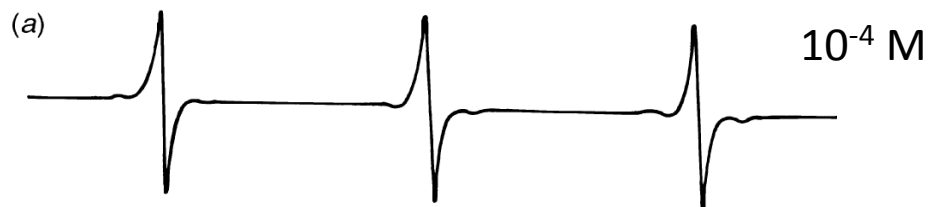


Exchange of spin orientations between two nuclei
 → Decreases interaction time τ between (same) nuclear and electron spin state

Spin exchange rate:

$$k_{\text{exchange}} = \frac{1}{2\tau[\text{conc.}]}$$

At high concentrations τ can become so small, that the time-averaged hyperfine field is close to zero and hfs coalesces to a single, narrowed line



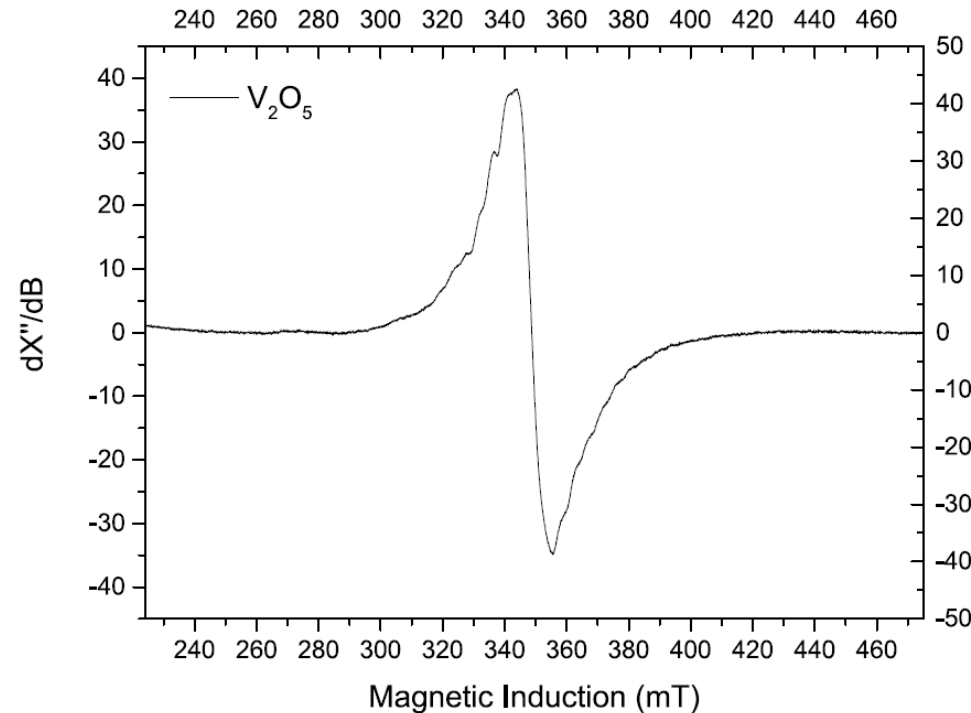
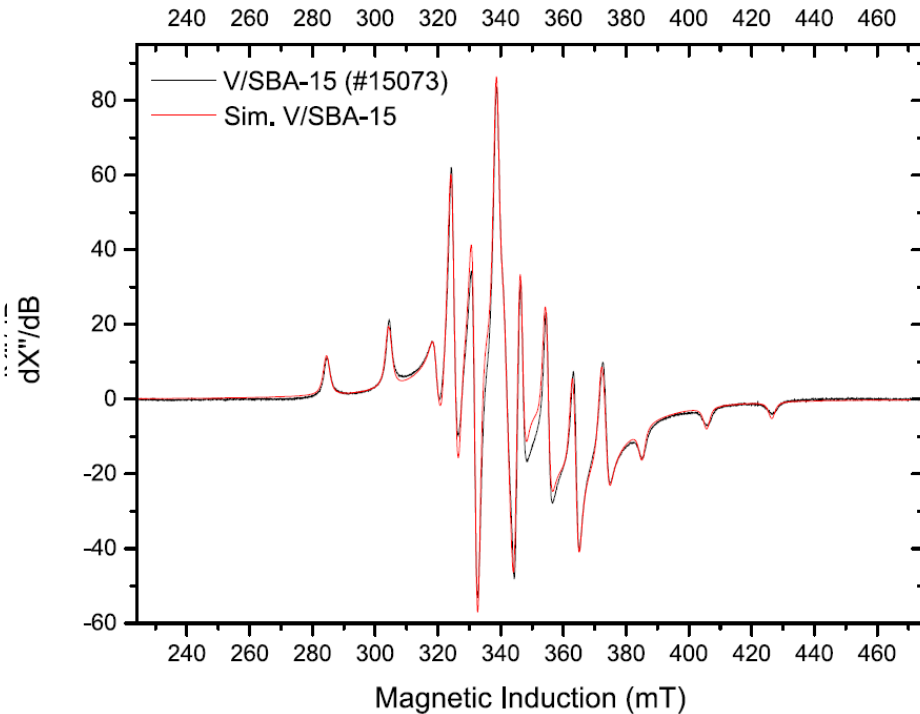
5. Linewidths

➤ Example for line broadening: electron-spin exchange

V/SBA-15

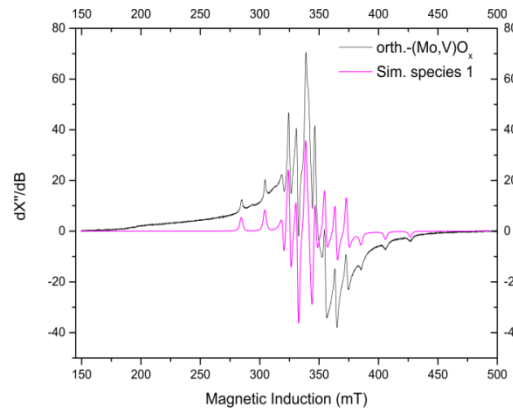
^{51}V : $I = 7/2$

V/SBA-15	Spin S	Nucleus	g Tensor			A Tensor (mT)			ΔB_{pp} (mT)
			g_{xx}	g_{yy}	g_{zz}	A_{xx}	A_{yy}	A_{zz}	
Narrow species	$\frac{1}{2}$	^{51}V	1.979	1.980	1.935	7.22	7.72	20.3	1.0
Broad species	$\frac{1}{2}$			1.969		/			18

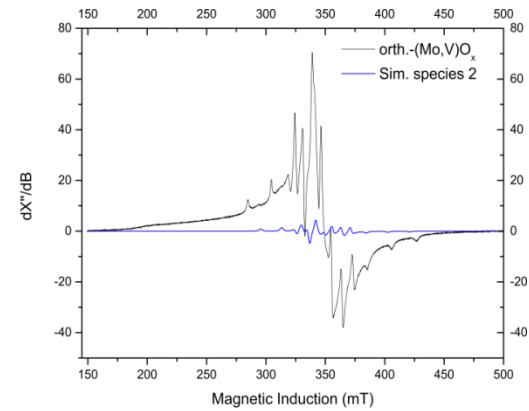


Differentiation between isolated and strongly interacting V atoms possible

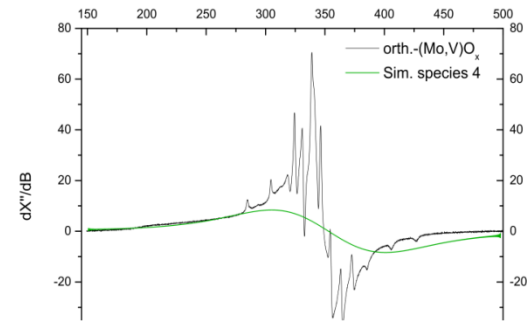
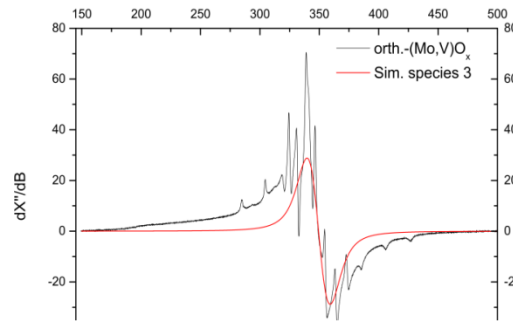
Orthorhombic MoV oxide



a)



b)



Species	Nucleus	Description and Assignment	Ref
1	^{51}V (V^{4+})	Narrow isolated V^{4+} -species, axially symmetric g and A tensors, “isolated V^{4+} ion in an axial crystal field (tetragonal distortion) ... typical for a vanadyl ion VO^{2+} with aqua or comparable ligands” ^a , similar to $[\text{VO}(\text{H}_2\text{O})_5]^{2+}$ (like heteropoly compound $\text{HVOPVMo}_{12}\text{O}_{40}$ of Lee et al. [3])	Lee et al. [3], Luca et al. [4]
2	^{51}V (V^{4+})	Same as “species 1” but with lower intensity, (like heteropoly compound $\text{H}_4\text{PVMo}_{11}\text{O}_{40}$ of Lee et al. [3])	Lee et al. [3], Luca et al. [4]
3	V	Broad species centred at $g = 1.97$ indicating interacting V^{4+} -centres (line broadening caused by spin-spin interactions): bulk species	Luca et al. [4]
4	V (or Mo)	Broad species centred at $g = 1.95$ indicating strongly interacting paramagnetic species (V^{4+} or Mo^{5+}): bulk species	
5	/	/	/

6. Literature

- ✘ J. A. Weil, J. R. Bolton, “Electron paramagnetic resonance“, John Wiley & Sons **2007**
(comprehensive book about cw-EPR theory, also as e-book available)
- ✘ G. R. Eaton, S. S. Eaton, D. P. Barr, R. T. Weber, “Quantitative EPR”; Springer **2010**
- ✘ A. Schweiger, “Pulsed Electron Spin Resonance Spectroscopy: Basic principles, Techniques, and Examples of Applications“, Angew. Chem. Int. Ed. **1991**, *30*, 265-292
- ✘ A. Schweiger, G. Jeschke, “Principles of pulse electron paramagnetic resonance“, Oxford Univ. Press **2001**
- ✘ P. Rieger, “Electron Spin Resonance: Analysis and Interpretation“, Royal Soc. of Chemistry **2007** (also as e-book available)