



Modern Methods in Heterogeneous Catalysis Research

Charge transport in catalysis

10th January 2014

Outline

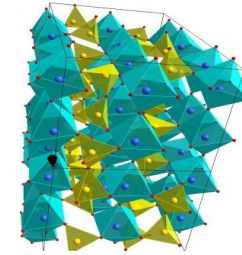
1. Introduction: Some practical examples
2. Semiconductors
3. Surface states and space charge
4. Bulk-surface charge transfer
5. Literature

1. Introduction

Experimental hints for the participation of “lattice” oxygen (oxygen covalently bound to the active transition metal ion) and “lattice” charge carriers in the reaction:

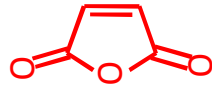
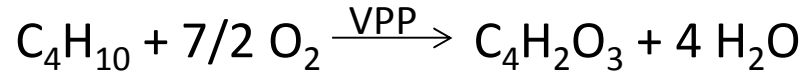
- 1) Reaction runs (awhile) without gas phase oxygen present (riser reactor concept)
- 2) Different oxidation states of the transition metal are monitored in dependence on the partial pressures of reaction gases or on contact time
- 3) By using $^{18}\text{O}_2$ in isotope exchange experiments, first ^{16}O (from the catalyst) is found in the reaction product, and only after some time ^{18}O
- 4) After long $^{18}\text{O}/^{16}\text{O}$ exchange, ^{18}O is found in the catalyst lattice (e.g. in ToF-SIMS or Raman experiments)
- 5) Conductivity changes upon reaction conditions, correlation between conductivity and activity/selectivity for differently doped semiconductors

1. Introduction

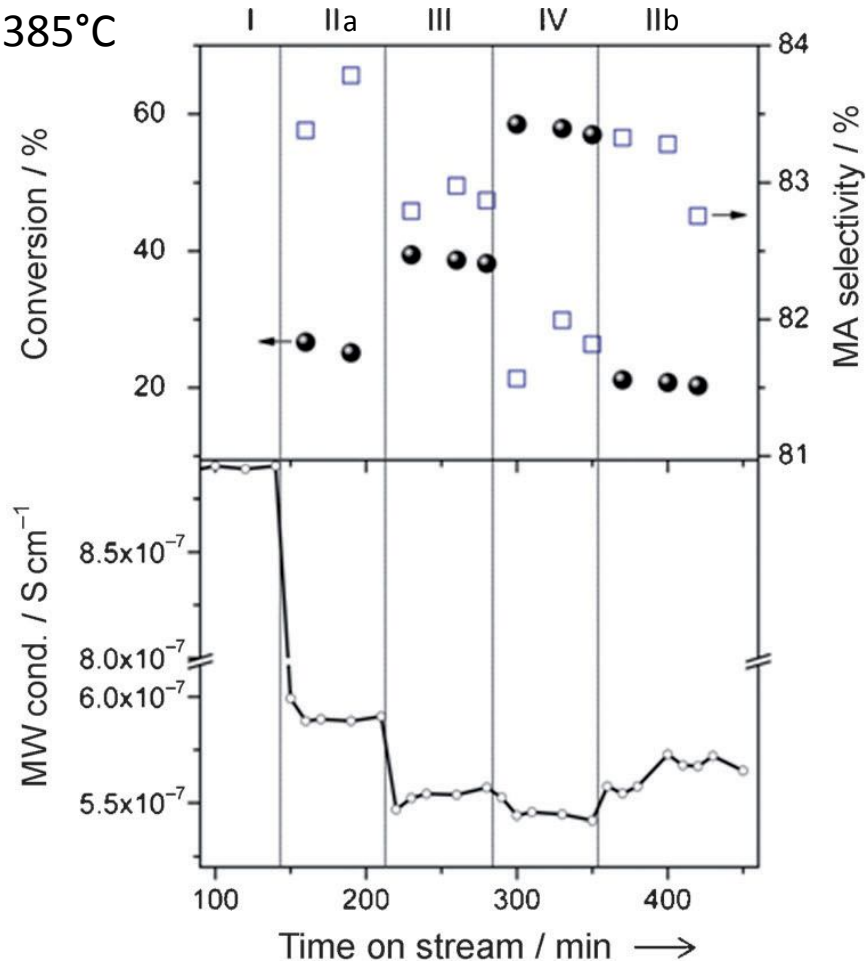


n-butane

maleic anhydride



385°C

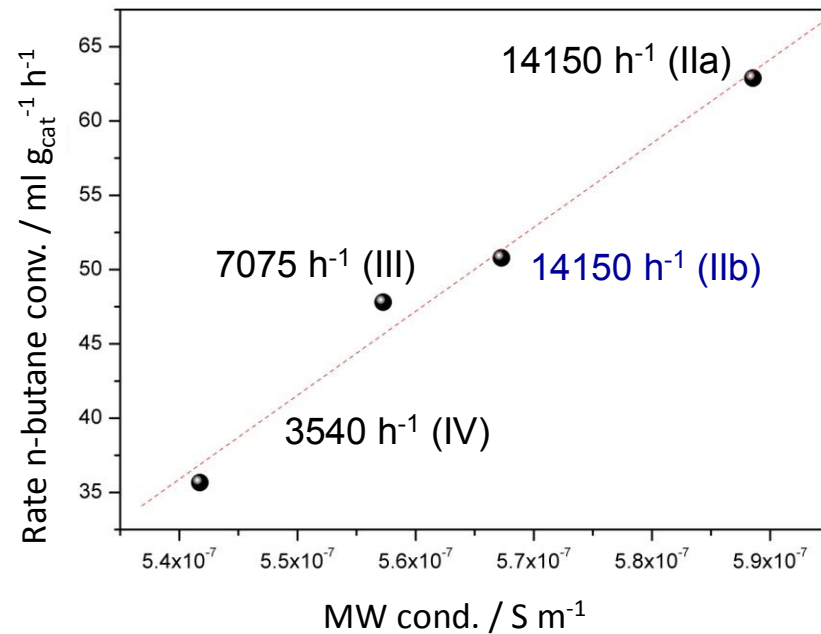


I: 100% N₂, GHSV: 14150 h⁻¹

II: 2% n-C₄H₁₀, 20% O₂ in N₂, GHSV: 14150 h⁻¹ (a,b)

III: 2% n-C₄H₁₀, 20% O₂ in N₂, GHSV: 7075 h⁻¹

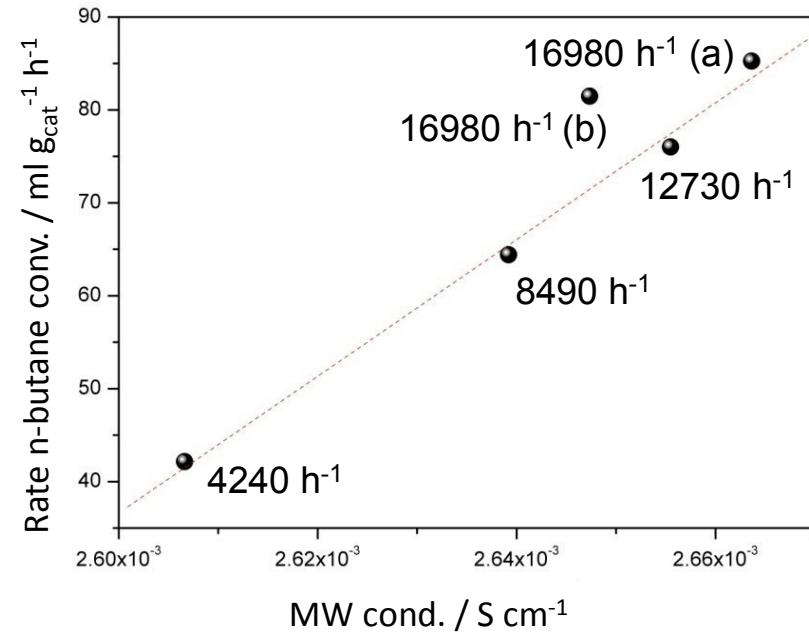
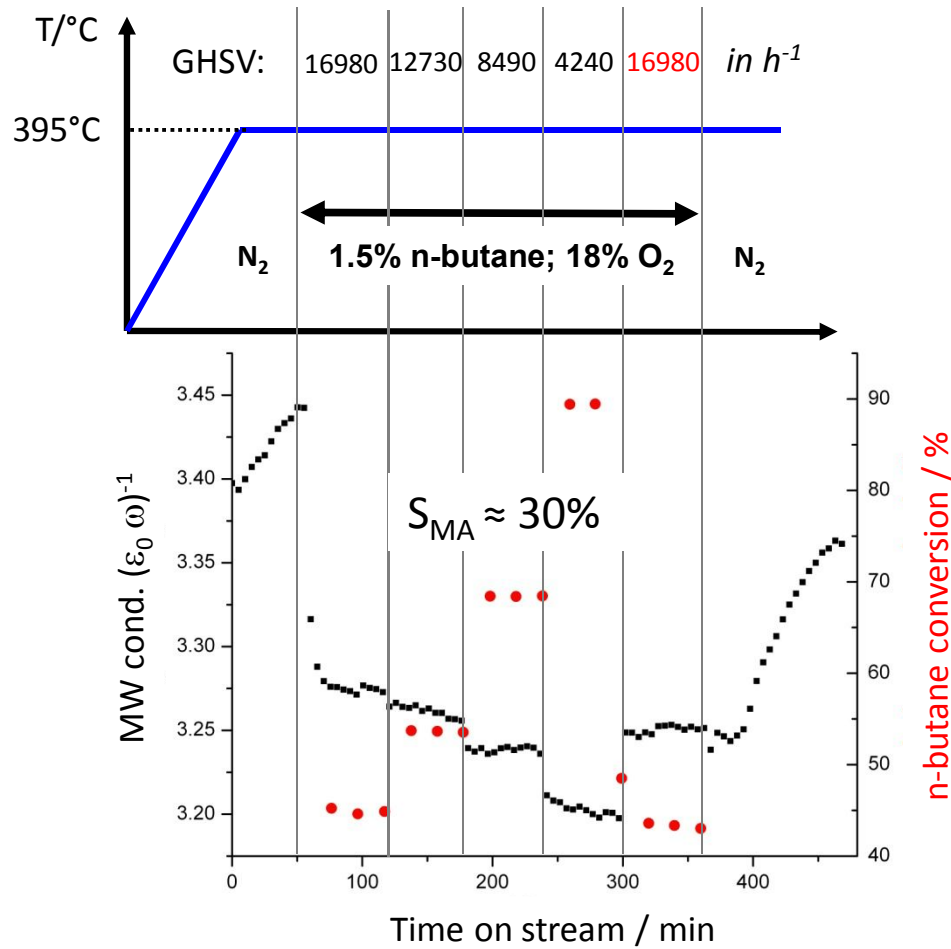
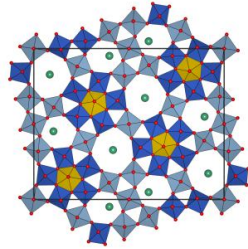
IV: 2% n-C₄H₁₀, 20% O₂ in N₂, GHSV: 3540 h⁻¹



→ Relationship between conductivity and activity

1. Introduction

MoVTenbO_x M1



→ Relationship between conductivity and activity

1. Introduction

ANGEWANDTE CHEMIE

HERAUSGEGEBEN VON DER GESELLSCHAFT DEUTSCHER CHEMIKER

67. Jahrgang · Nr. 17/18 · Seite 433–540 · 7. September 1955

FORTSETZUNG DER ZEITSCHRIFT »DIE CHEMIE«

Neuere Gedanken zur Natur der heterogenen Katalyse

Von Prof. Dr. GEORG-MARIA SCHWAB¹⁾

Physikalisch-Chemisches Institut der Universität München

In letzter Zeit sind nähere Aufschlüsse über die Natur der Grenzflächenkatalyse durch die Anwendung der Festkörperphysik erreicht worden. Versuche an Legierungen haben gezeigt, daß die Aktivierung in einem Übergang von Elektronen entweder zum oder vom Katalysator besteht. Inzwischen haben diese Gesichtspunkte sich auch auf halbleitende Oxydkatalysatoren übertragen lassen. Die Chemisorption ist somit als Elektronenübergang deutbar. Verschiedene neuere Einzelfälle werden beschrieben.

1. Introduction

Dehydrogenation of formic acid with different bronze alloys: **activation energy** and **resistivity**

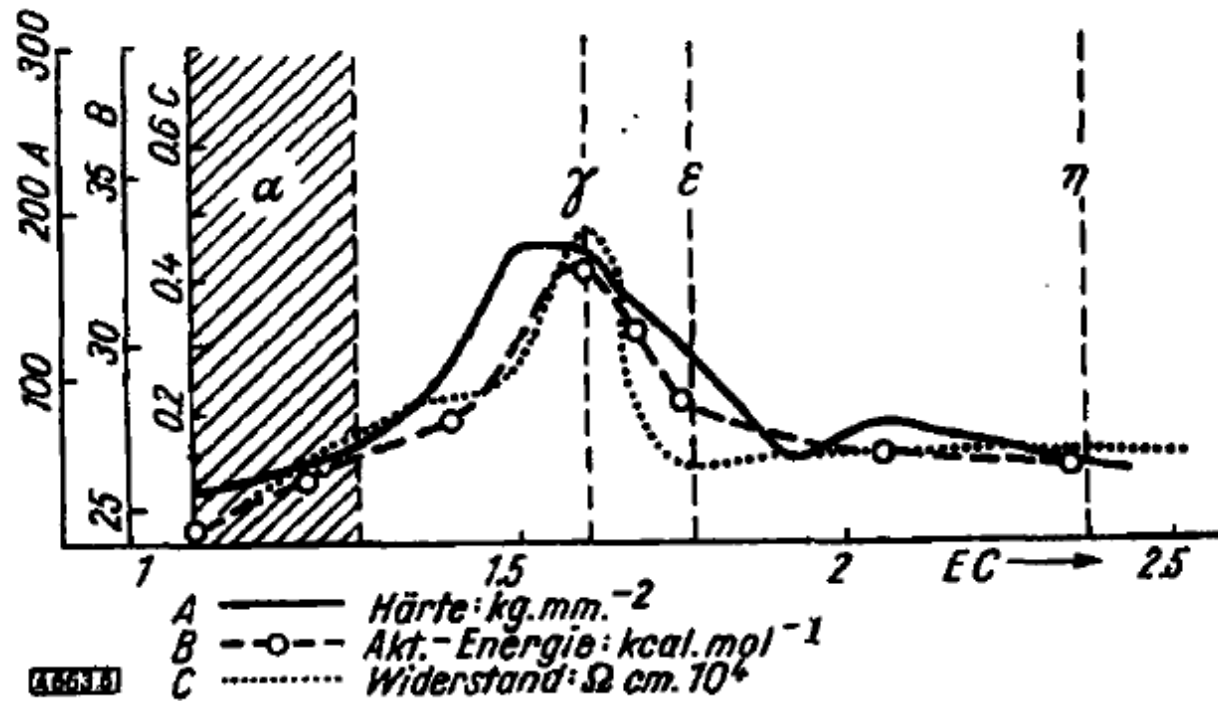
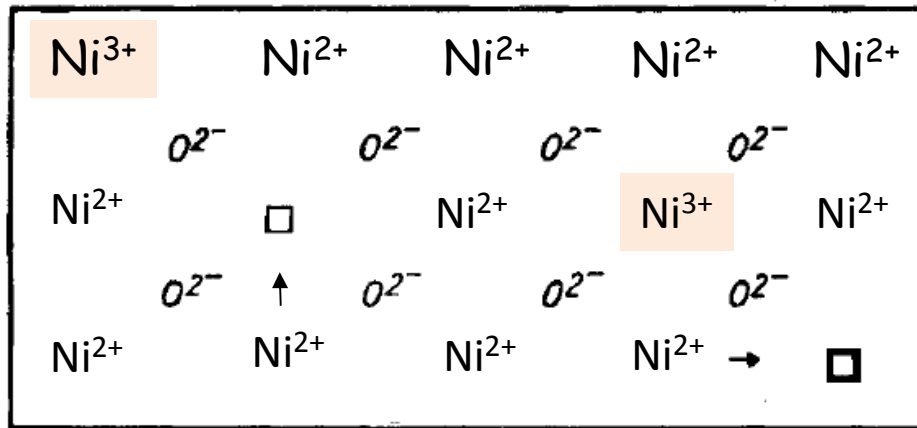


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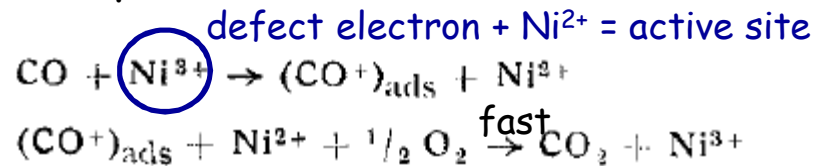
Aktivierungsenergie, Widerstand und Härte bei Bronzen
verschiedener Elektronenkonzentration EC

1. Introduction

Defect electrons (holes) in NiO (with O excess): **p-type semiconductor**

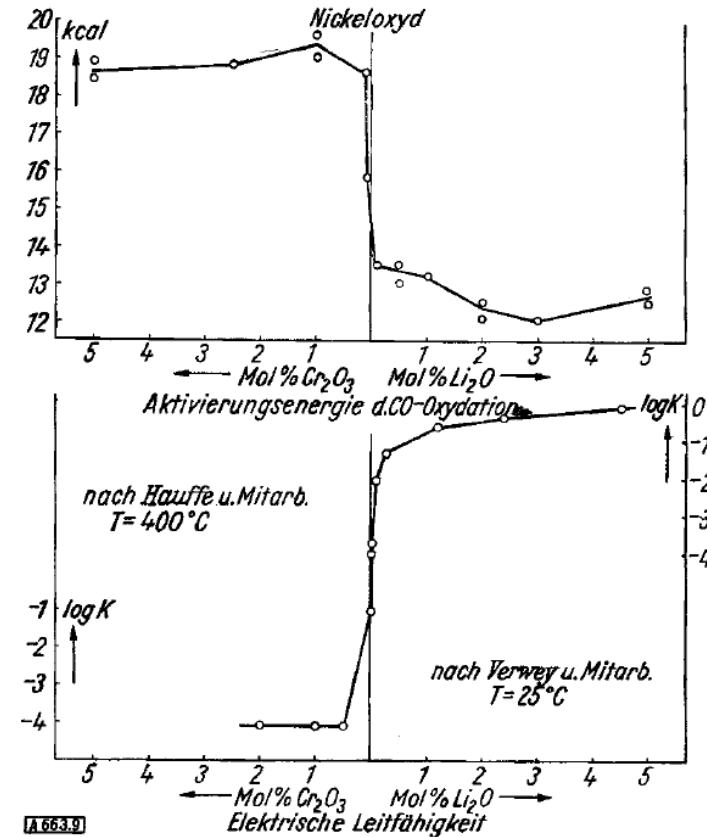


Explanation: donor reaction



→ CO chemisorption rate-determining

Increase of defect electrons by Li₂O, decrease by Cr₂O₃ addition:
 $\text{CO} + \frac{1}{2} \text{O}_2 \rightarrow \text{CO}_2$

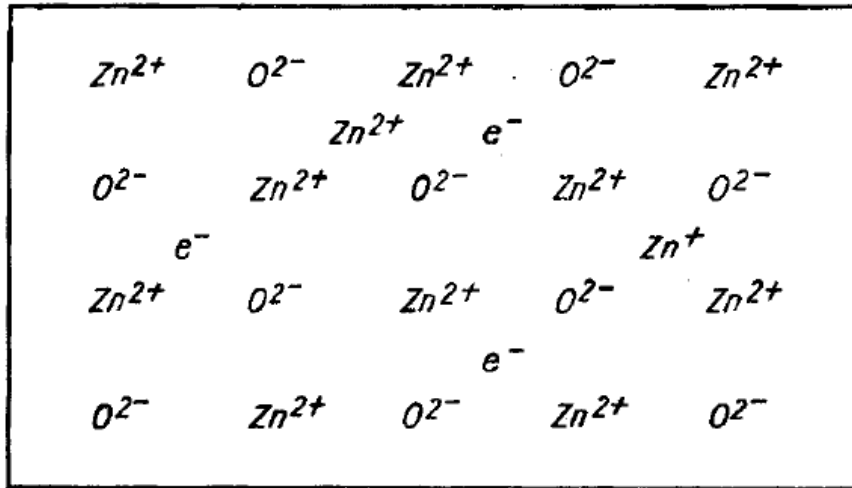


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Leitfähigkeit und Aktivierungsenergie an dotiertem NiO

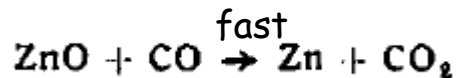
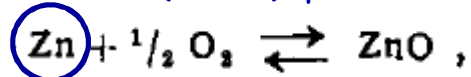
1. Introduction

Excess electrons in ZnO (with Zn excess): **n-type semiconductor**



Explanation: acceptor reaction

$Zn^{2+} + (1 \text{ or } 2) \text{ quasi-free electrons} = \text{active site}$



→ Oxidation of active site rate-determining (strong influence of oxygen on rate)

Increase of quasi-free electrons by Ga_2O_3 , decrease by Li_2O addition:
 $CO + \frac{1}{2} O_2 \rightarrow CO_2$

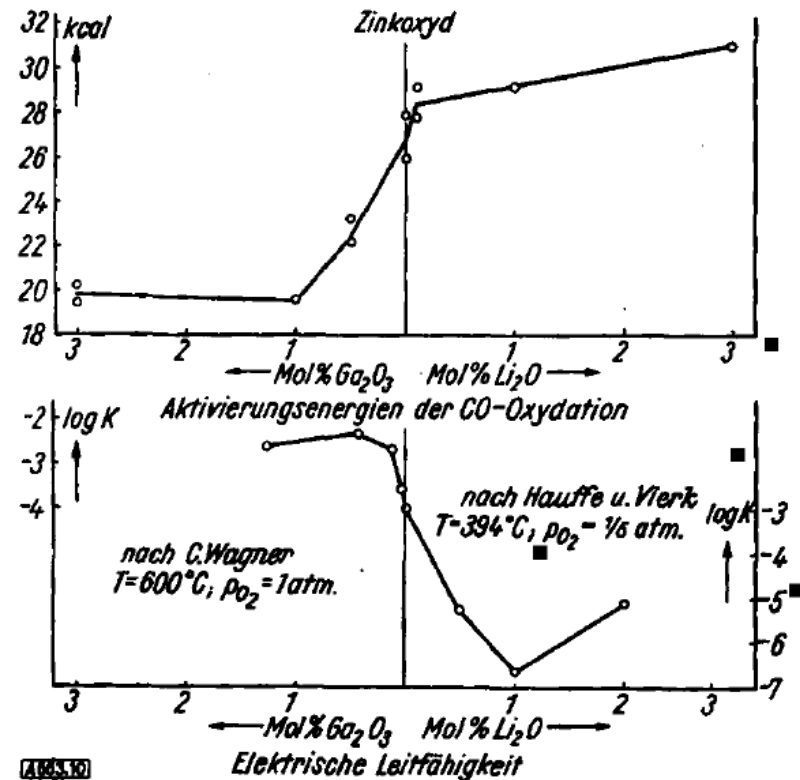


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Leitfähigkeit und Aktivierungsenergie an dotiertem ZnO

1. Introduction

✘ Oxidation catalysts are most often metal oxides, and most oxides are semiconductors!

- n-type semiconductor oxides contain anionic vacancies (V_O^{2-}) associated with deficit anionic oxygen
- p-type semiconductor oxides contain positive holes (h^+) as charge carriers associated with excess anionic oxygen

✘ Semiconductors: electrical conductivity varies exponentially with increasing temperature:

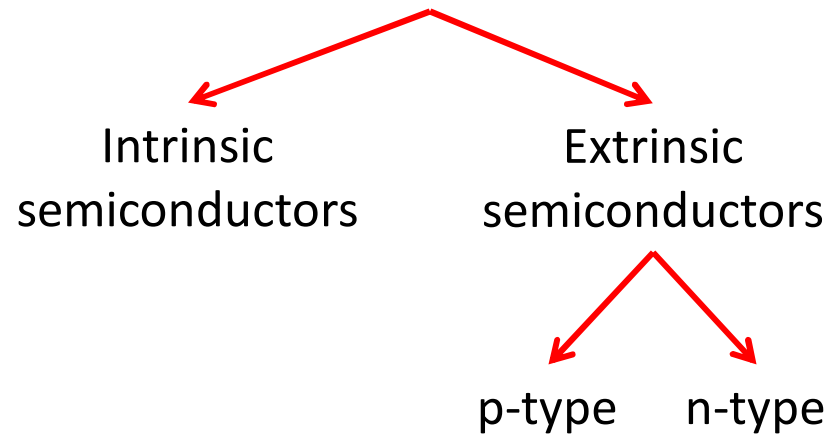
$$\sigma = \sigma_0 \exp(-E_c/RT)$$

1. Introduction

Electric conductivity

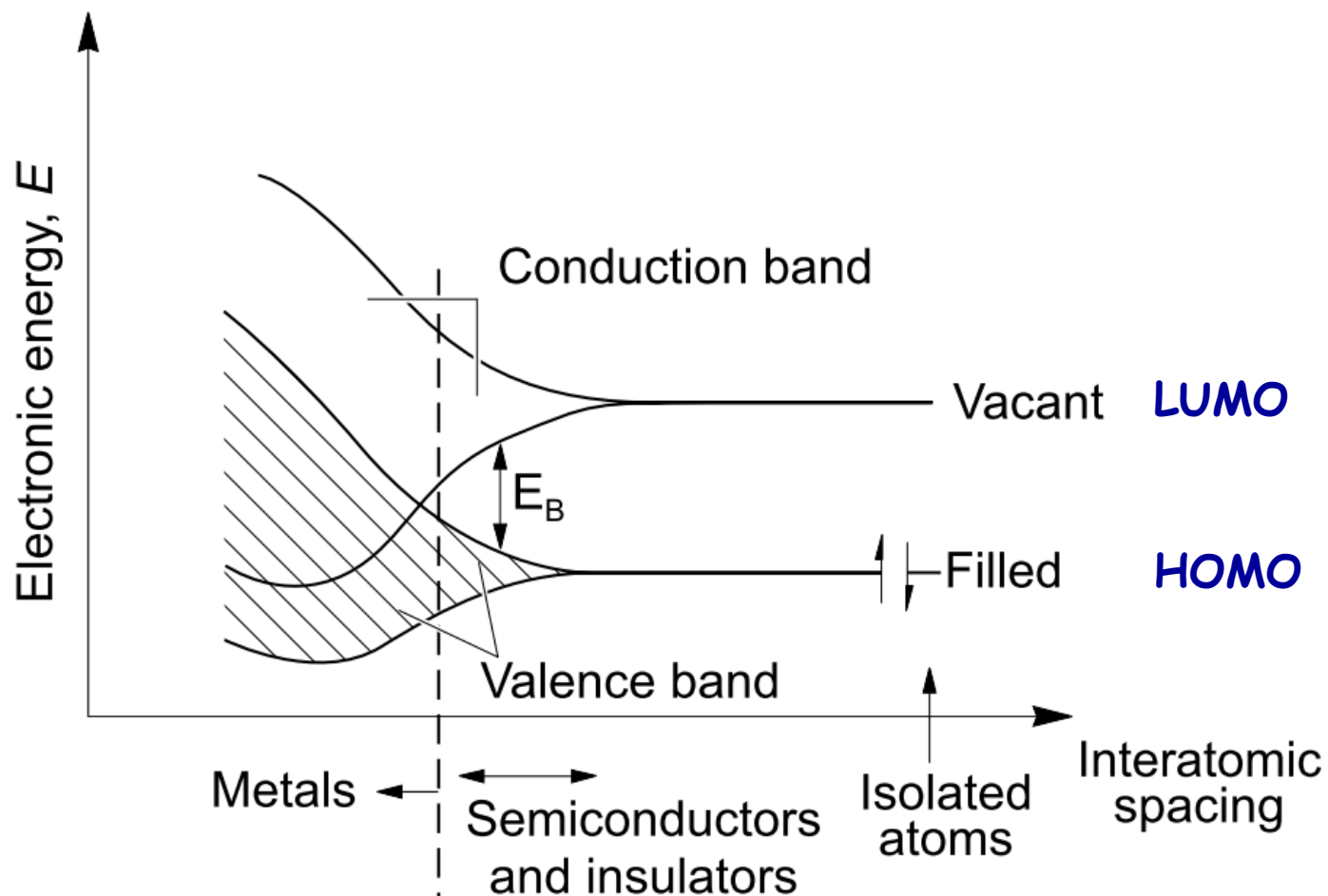


Metal oxide catalysts



Some semiconductor theory

2. Semiconductors



A. W. Bott, *Current Separations* 1998, 17, 87

2. Semiconductors

Occupancy of electrons in a band is determined by **Fermi-Dirac statistics**

Fermi-Dirac distribution (for an electron gas):

$$f(E) = \frac{1}{\exp\left(\frac{E - \mu}{k_B T}\right) + 1}$$

E ...Energy

k_B ...Boltzmann constant

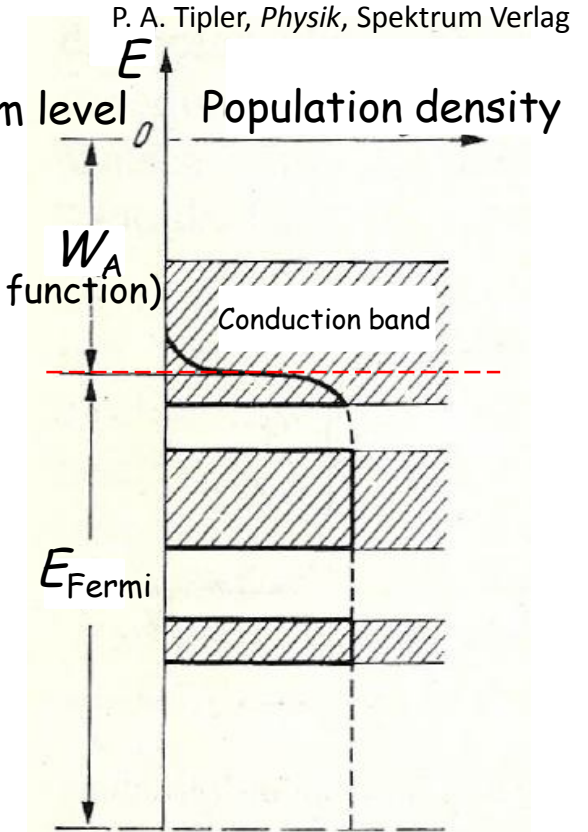
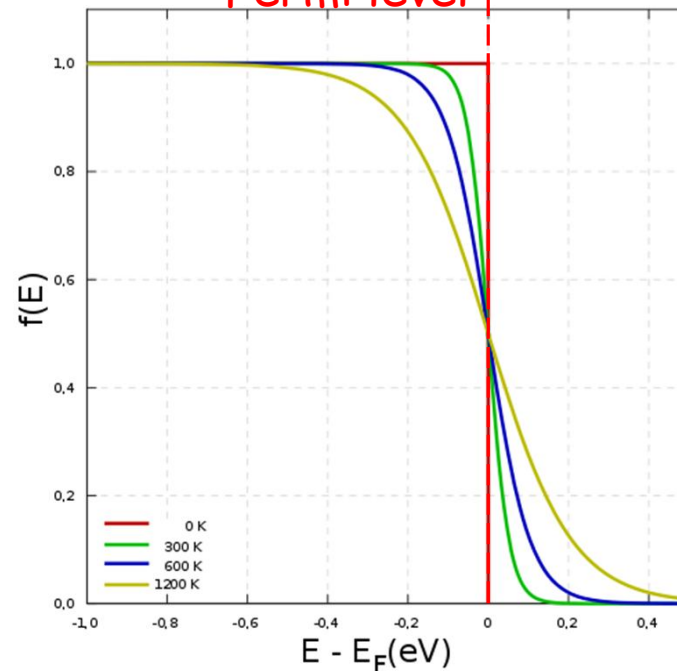
T ...Temperature

μ ...(Electro-)Chemical potential

$$\mu(T = 0) = E_{\text{Fermi}}$$

$$f(E = \mu) = 1/2$$

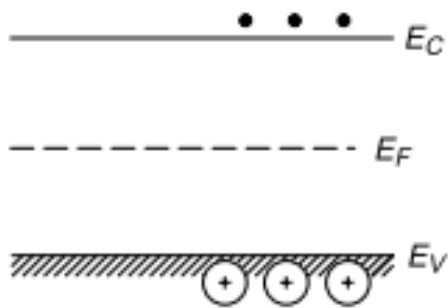
= Fermi level !



The Fermi curve determines the population of occupied states, independent of the existence of states in the regarded E region

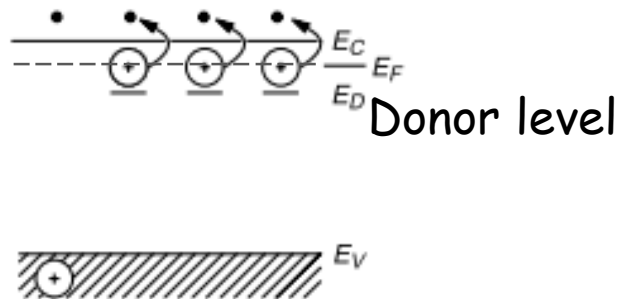
2. Semiconductors

Intrinsic semiconductor



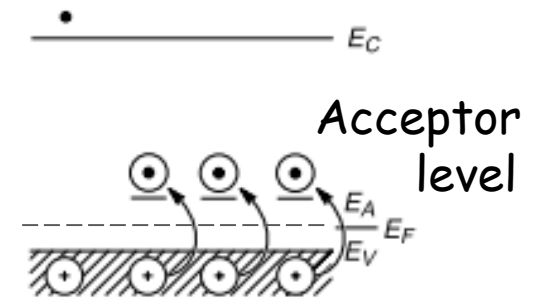
e.g. Si (band gap at 300 K = 1.12 eV)

(Extrinsic) n-type semiconductor

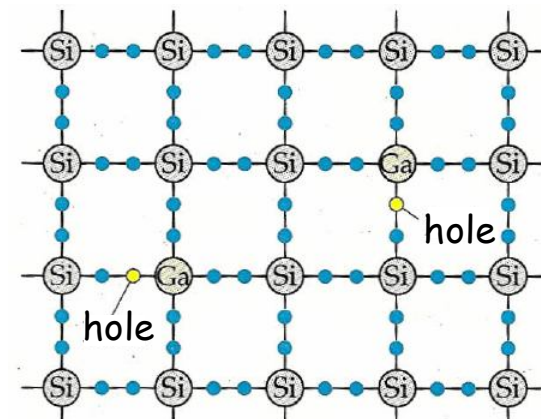
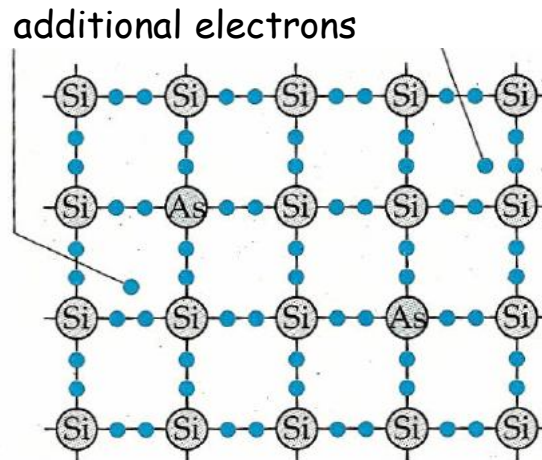


e.g. As-doped Si

(Extrinsic) p-type semiconductor



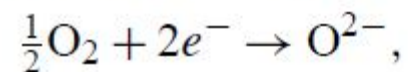
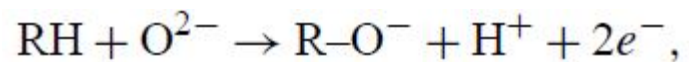
e.g. Ga-doped Si



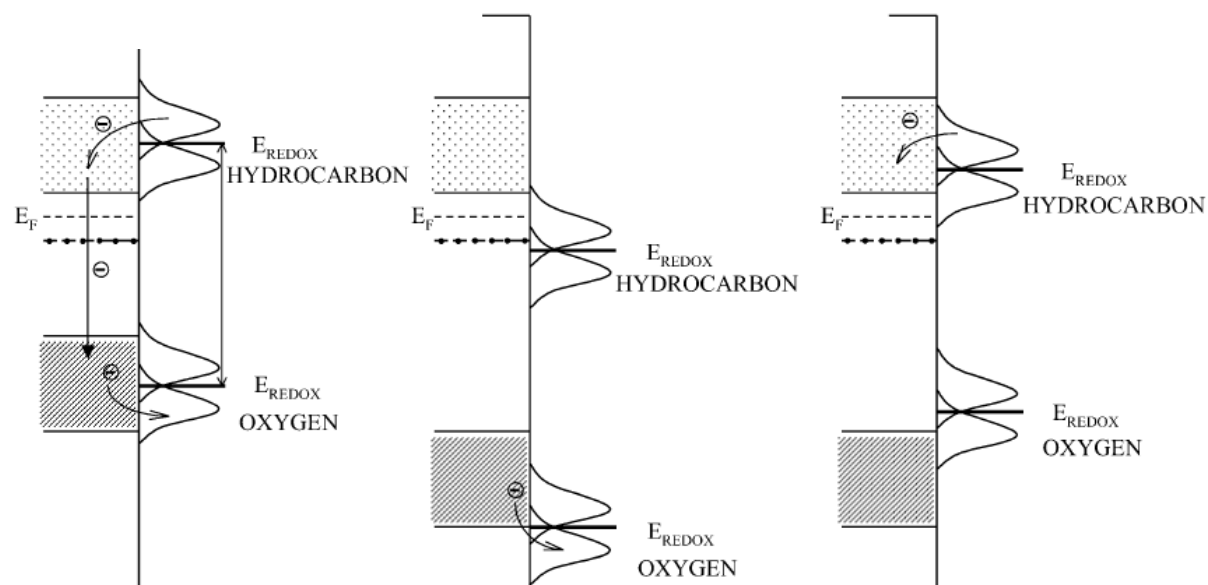
P. A. Tipler, *Physik*, Spektrum Verlag

2. Semiconductors

Selective oxidation of hydrocarbons, consideration of bands and frontier orbitals:



Rigid band assumption (no local surface states)



CATALYTIC OXIDATION OF HYDROCARBON MOLECULE CAN PROCEED

CATALYTIC OXIDATION OF HYDROCARBON MOLECULE CANNOT PROCEED, BECAUSE THE MOLECULE IS NOT ACTIVATED

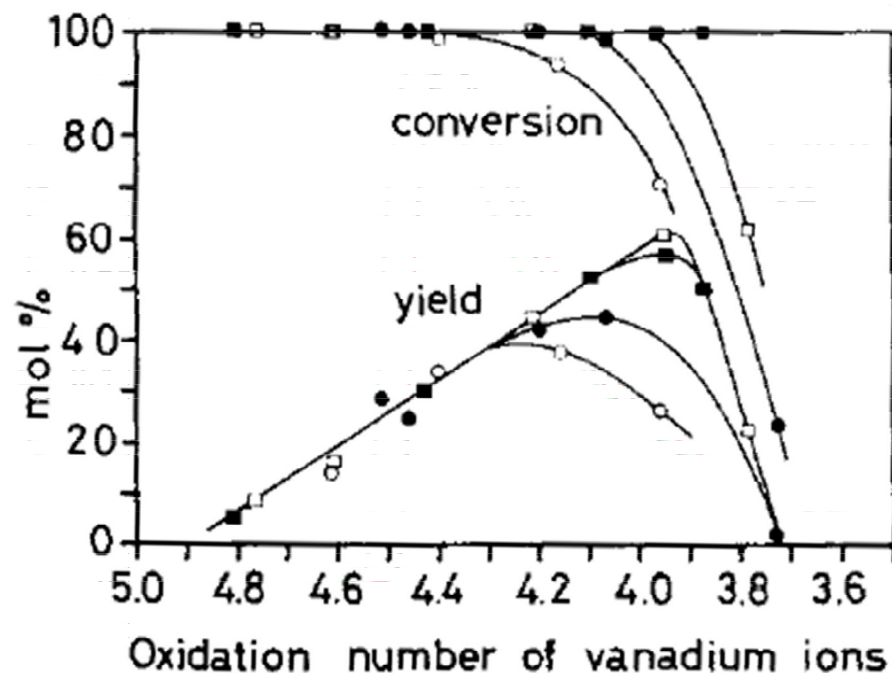
CATALYTIC OXIDATION OF HYDROCARBON MOLECULE CANNOT PROCEED, BECAUSE THE CATALYST IS NOT REOXIDIZED

J. Haber, M. Witko, *J. Catal.* **2003**, 216, 416-424

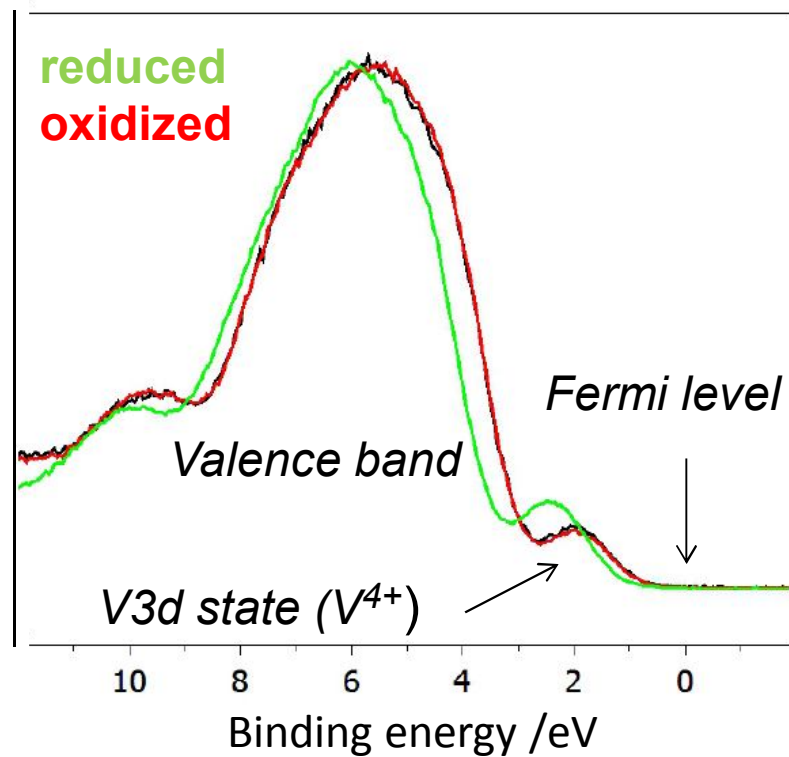
2. Semiconductors

V_2O_5/P_2O_5 1-butene \rightarrow maleic anhydride

Nakamura et al., *J. Catal.* 1974, 34, 345



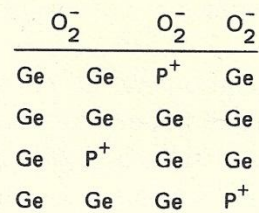
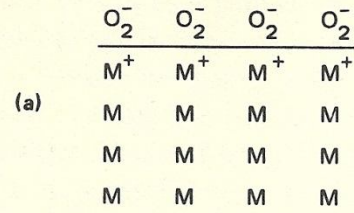
XP spectrum of VPP:



- Rate limiting: reduction of O_2
- High V^{4+}/V^{5+} ratio \rightarrow V^{4+}/V^{5+} band almost completely occupied \rightarrow highest Fermi energy \rightarrow most rapid O_2 reduction
- Empty levels needed for electron injection from olefin

Catalysis occurs at interfaces: the formation of surface states and space charge regions

Atomistic model:



surface charge Q

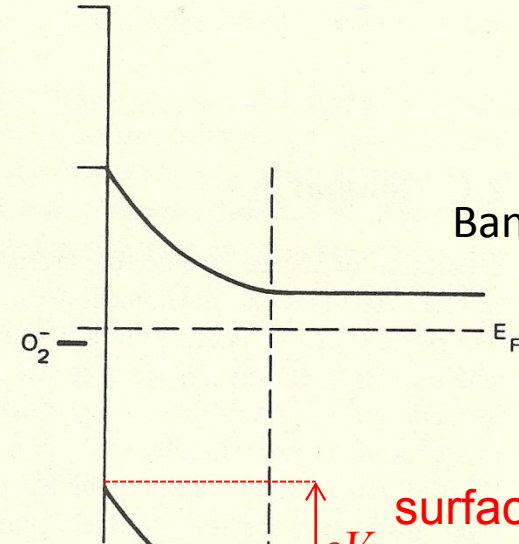
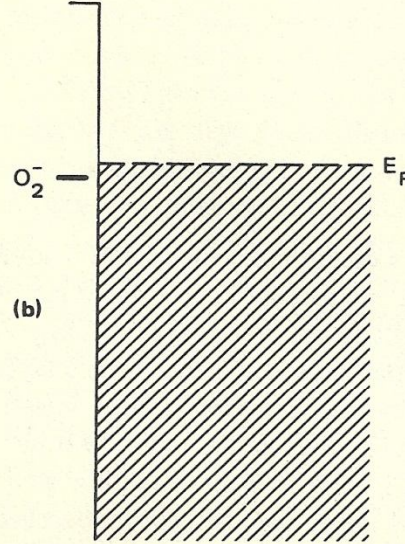
(space) charge density ρ

Double layers:

parallel plate

space charge model

Band model:



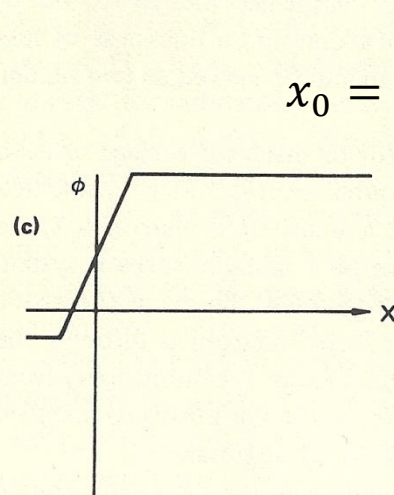
Band bending function:

$$V = \phi_{bulk} - \phi$$

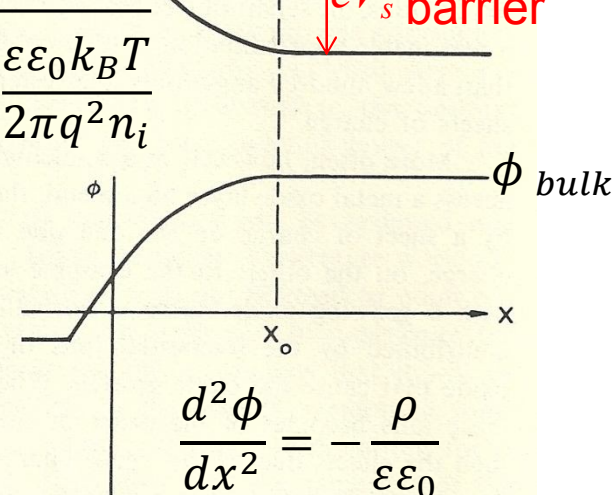
$$x = 0; V = V_s:$$

$$\frac{dV}{dx} = \frac{Q}{\epsilon\epsilon_0}$$

Potential:



$$x_0 = \sqrt{\frac{\epsilon\epsilon_0 k_B T}{2\pi q^2 n_i}}$$



$$\frac{d^2\phi}{dx^2} = -\frac{\rho}{\epsilon\epsilon_0}$$

$$x = 0:$$

$$\frac{d\phi}{dx} = -\frac{Q}{\epsilon\epsilon_0}$$

3. Surface states and space charge

Models for the adsorbate/solid bond

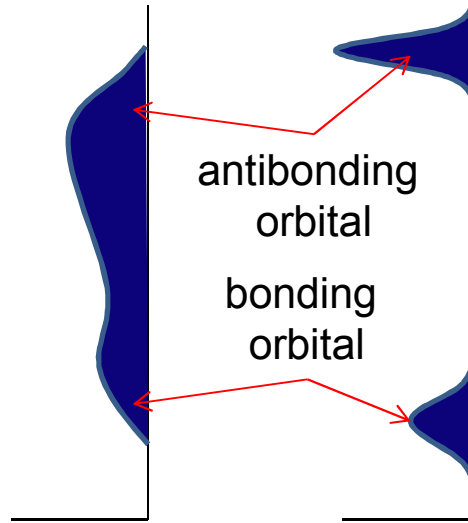
Weak Intermediate Strong
Interaction

Rigid band model

Atomistic model



Electronic state of an adsorbate atom



antibonding orbital
bonding orbital



Local density of states (LDS) of a surface atom on solid

E
DOS

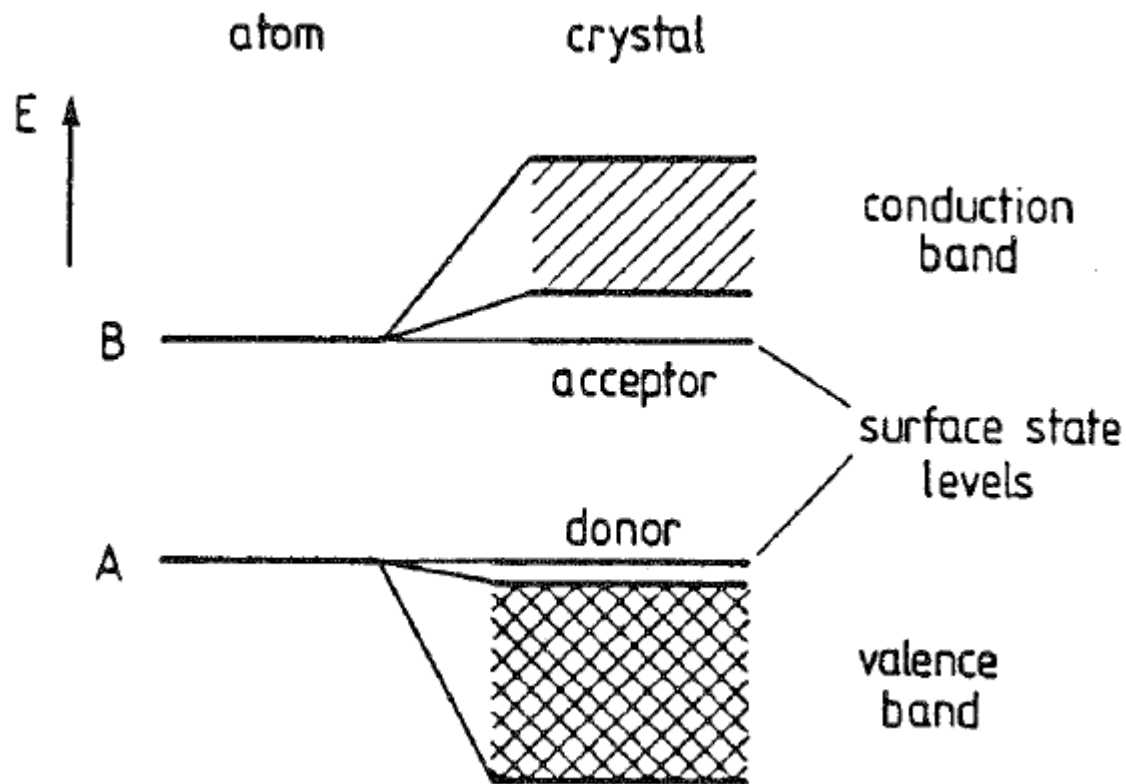
3. Surface states and space charge

Broadening of molecular orbitals of adsorbing species into (apparent) bands can occur through several mechanisms:

- 1) Heterogeneity of surface sites
- 2) Overlap of the orbitals of the adsorbate
- 3) Temporal fluctuations due to the presence of polar species
- 4) Interaction of the molecular orbitals of the adsorbate with the bands of the solid

3. Surface states and space charge

Surface states of a 3D crystal

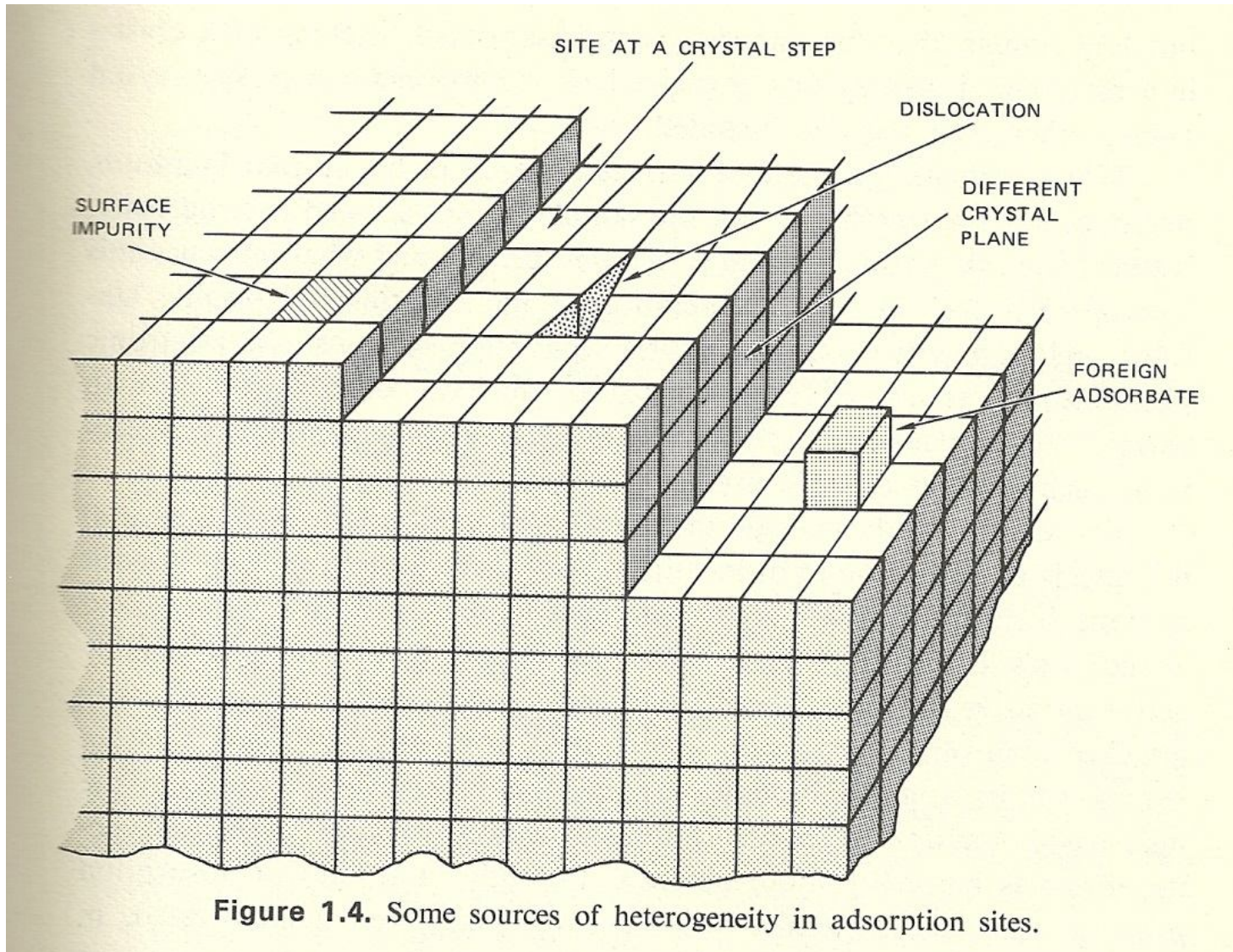


Intrinsic surface states: On ideal surfaces (perfect termination with 2D translational symmetry)

Extrinsic surface states: On surfaces with imperfections (e.g. missing atom)

3. Surface states and space charge

Sources of surface states:



3. Surface states and space charge

Surface states in the rigid band model

Intrinsic states

(e.g. Shockley states, Tamm states)

- Broken periodicity on the surface
- Undercoordinated surface atoms
- Surface termination with translational symmetry
- Relaxed, reconstructed or relocated surface

Extrinsic states

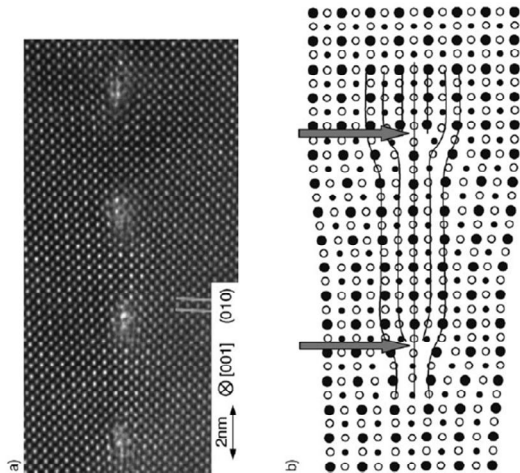
(e.g. adsorbates, defects, dislocations, surface enrichment of bulk impurities)

- Heterogeneous surface without translational symmetry

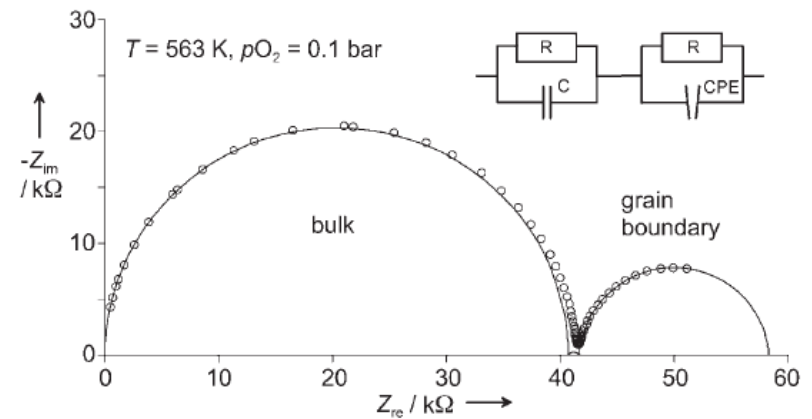
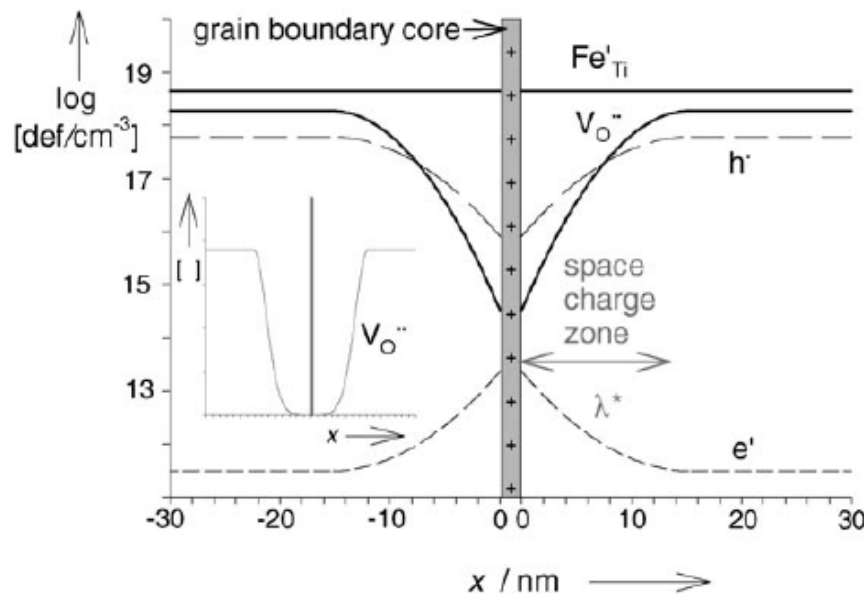
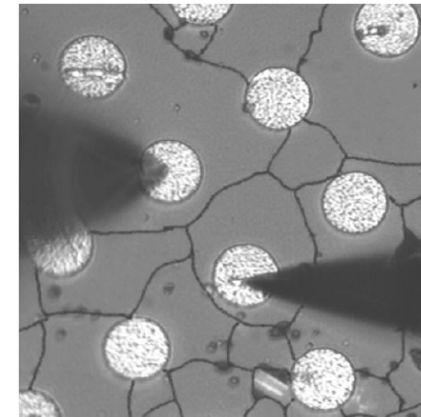
3. Surface states and space charge

Blocking grain boundary → formation of space charge region

SrTiO₃



Localized Impedance Spectroscopy



3. Surface states and space charge

Redox couples on surfaces forming surface states:

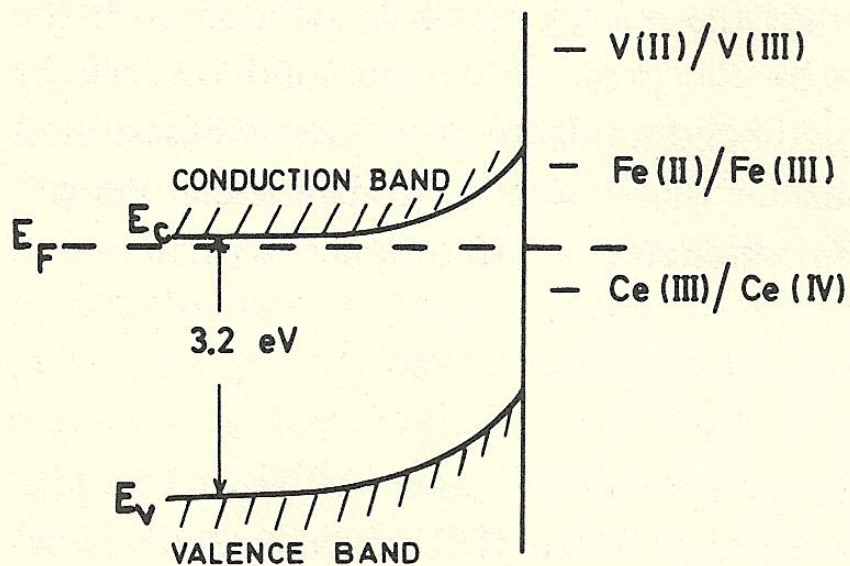
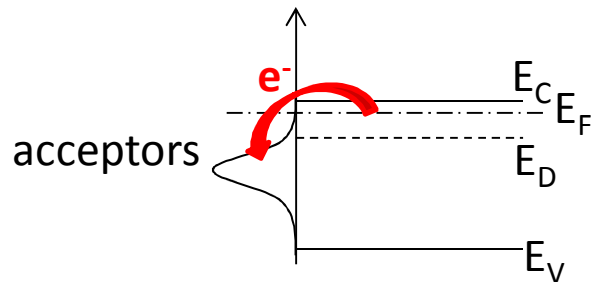


Figure 1.1. Surface states of chemical couples, to illustrate the expected energy level as the chemical properties of the species is varied. ZnO is used as the example semiconductor; the indicated position of the surface states relative to the bands is discussed further in the text.

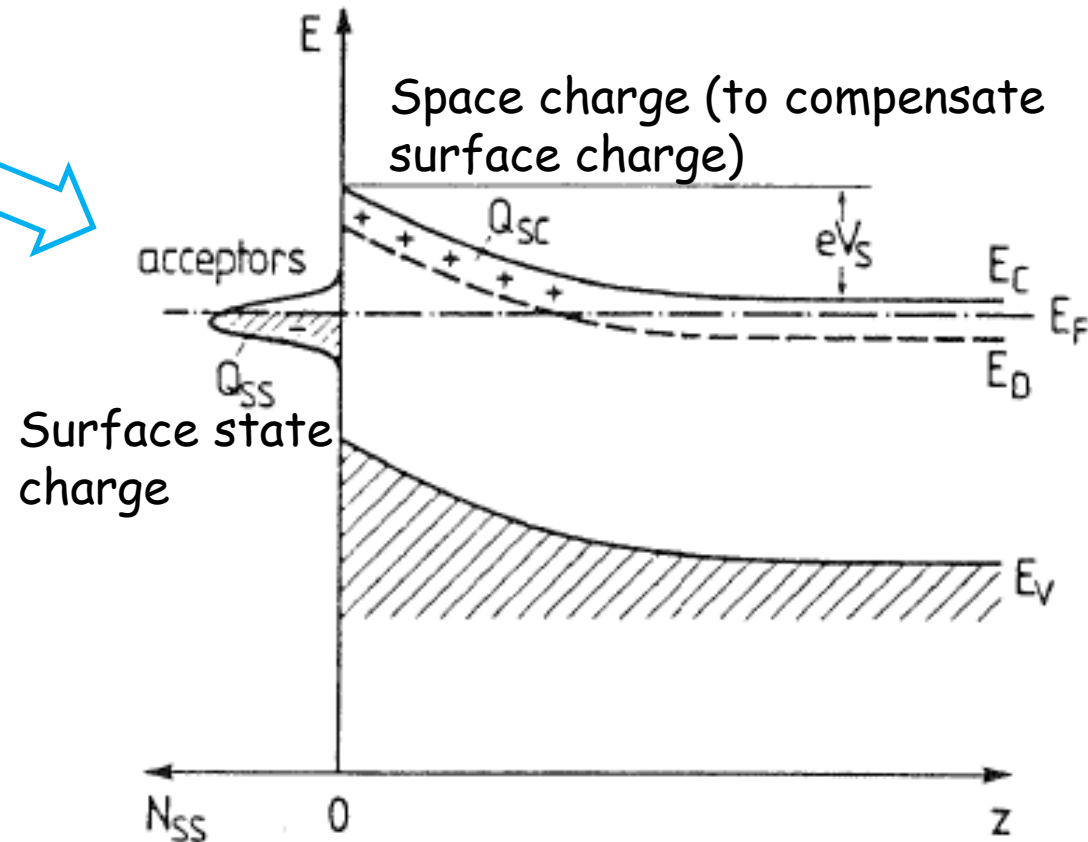
3. Surface states and space charge

Flatband situation (not equilibrated):



Build-up of negative charge in surface acceptor state is compensated by oppositely charged space charge region

Formation of (depletion) space charge layer:

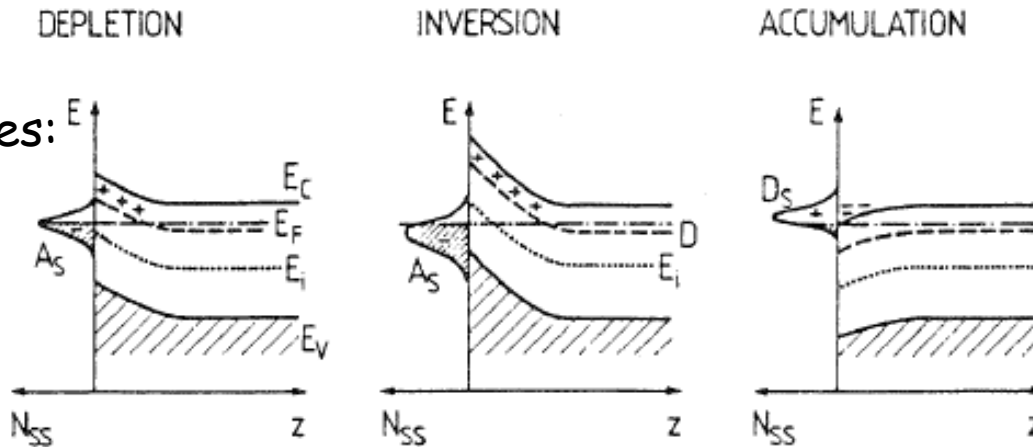


Neutrality condition $Q_{SS} = -Q_{SC}$ determines the Fermi level

3. Surface states and space charge

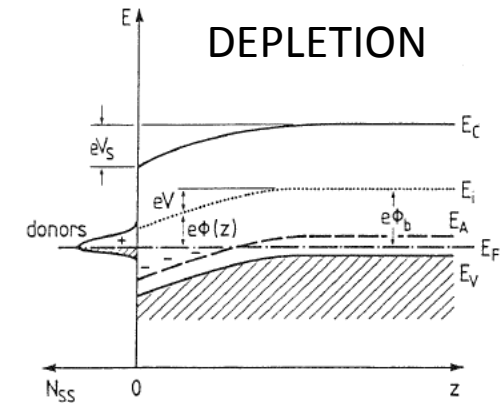
n-SC:

Band schemes:

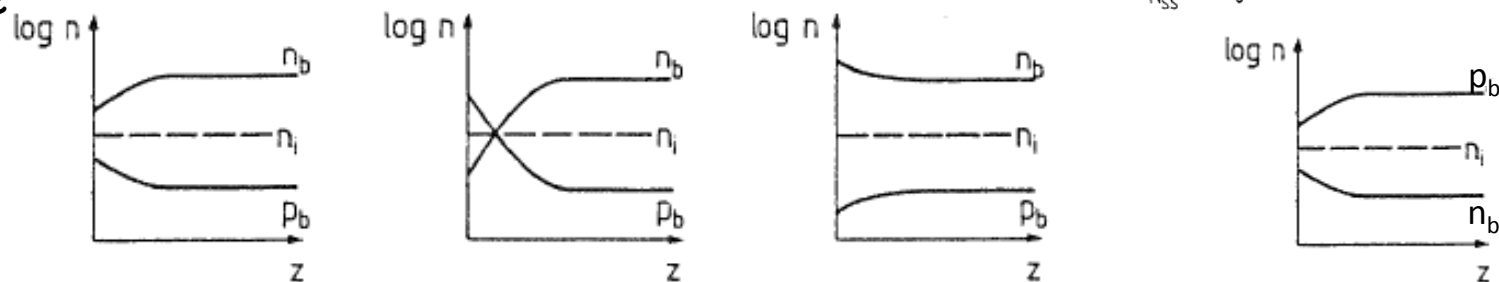


p-SC:

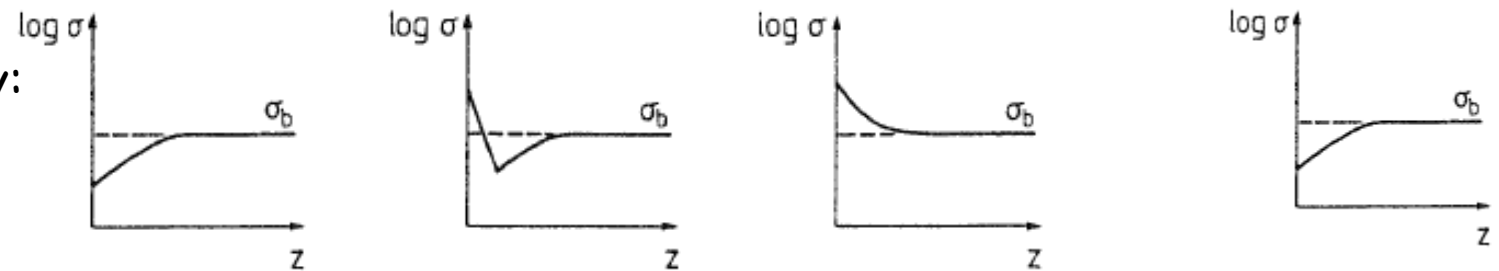
DEPLETION



Free charge carrier densities:



Local conductivity:



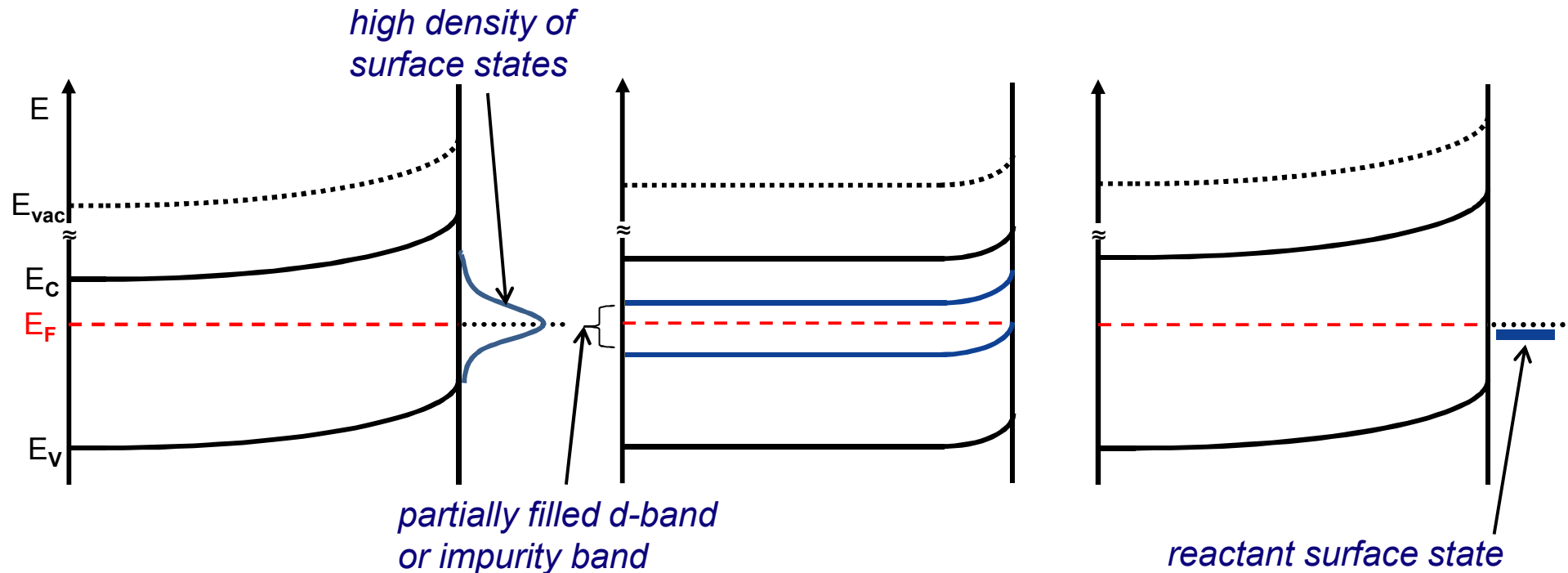
3. Surface states and space charge

Surface Fermi energy can be pinned by

1) Surface states

2) Partially filled narrow band (e.g. d-band)

3) Reactant



localized surface levels, e.g. Cr^{3+} on Cr_2O_3 , Mn^{2+}/Mn^{3+} on MnO , O^- , foreign surface impurities

in many transition metal oxides, e.g. Co_2O_3 , VO_2 , impurity band associated with vacancies (bonding \rightarrow surface molecule model)

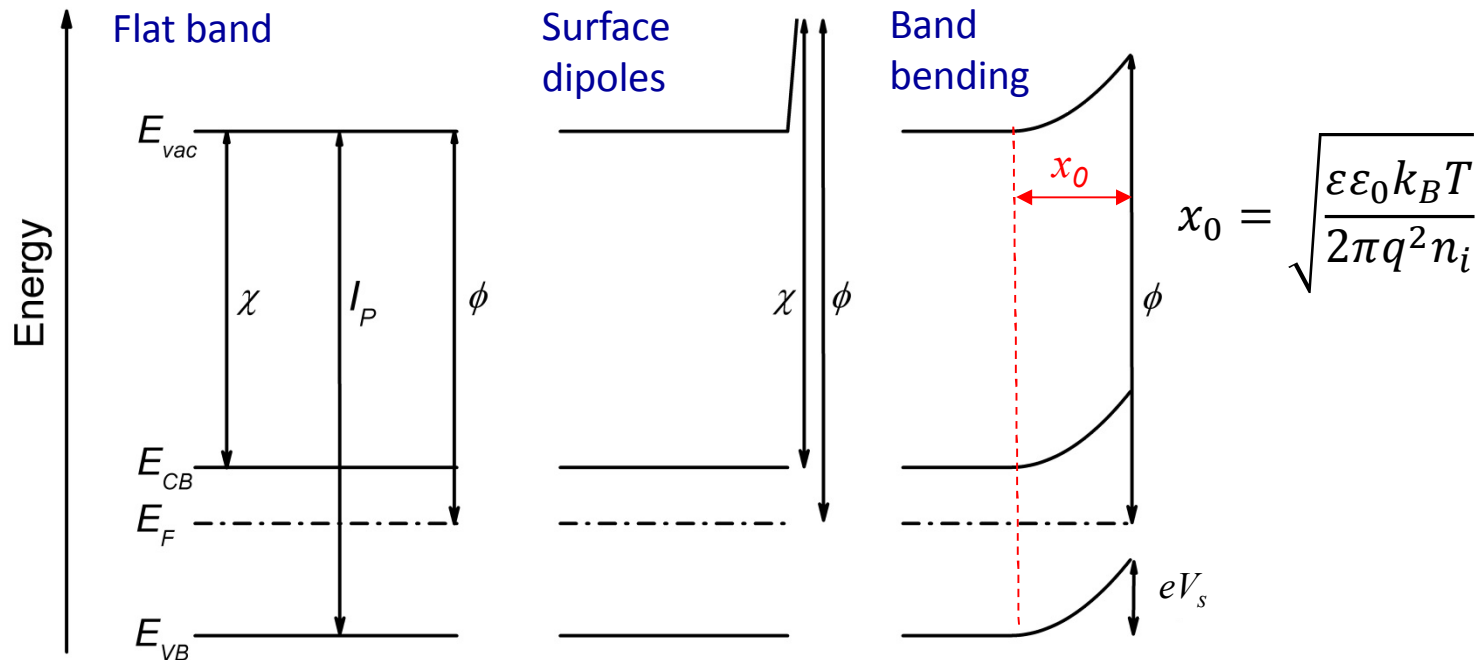
little local interaction, direct electron exchange with conduction band, e.g. ZnO

3. Surface states and space charge

Spectroscopic evidence? Semiconductor characterized by:

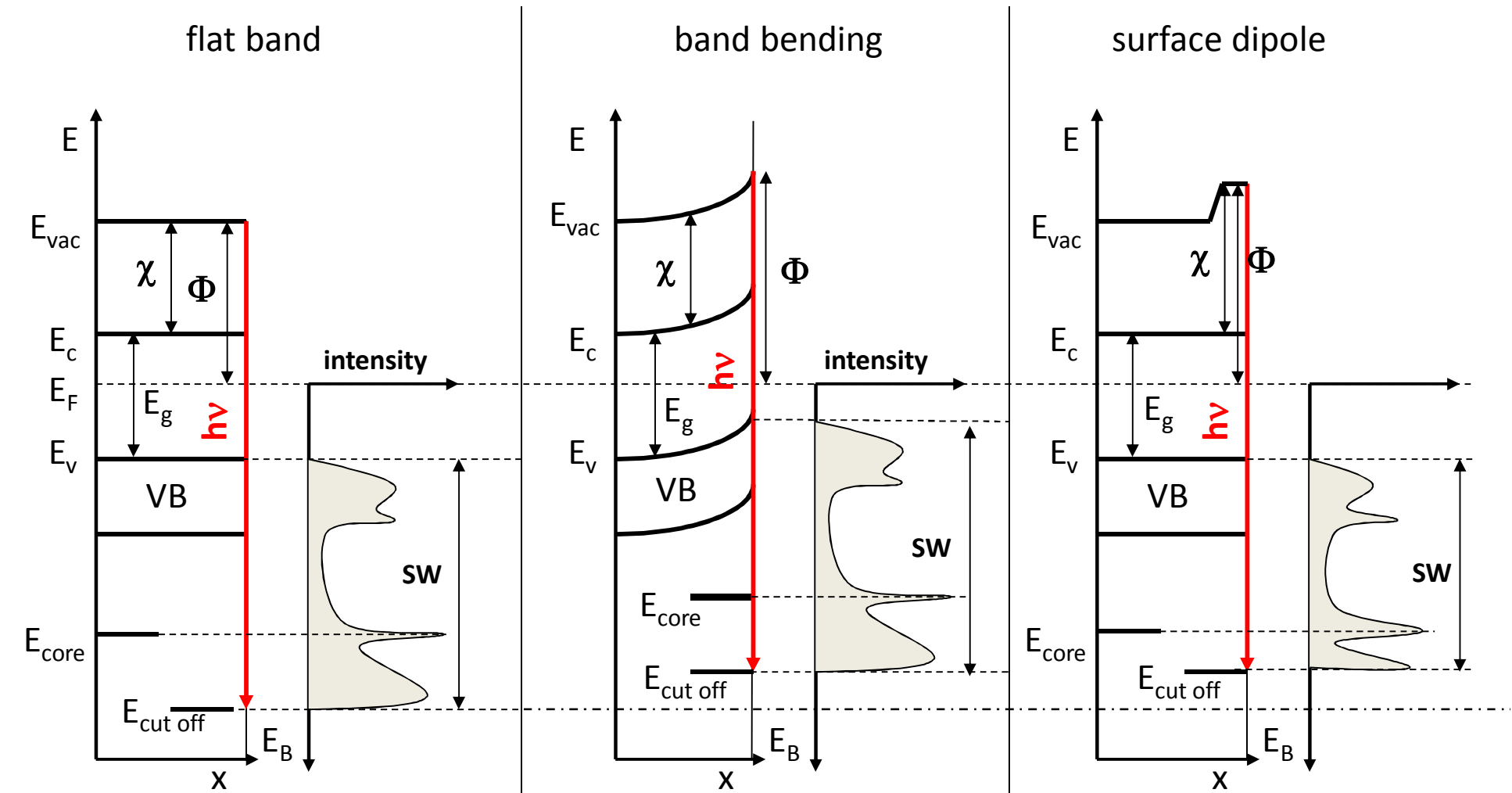
- the work function Φ
- the electron affinity χ
- the surface barrier eV_s (binding energy shift of valence band edge)
- the Debye length x_0 (width of depletion or accumulation layer)

-the conductivity: $\sigma(x) \approx en_e(x)\mu_e, n_e = N_{\text{eff}}^C \exp\left[-\frac{E_C(x) - E_F}{k T}\right]$



3. Surface states and space charge

Spectroscopic evidence? → Photoemission spectroscopy

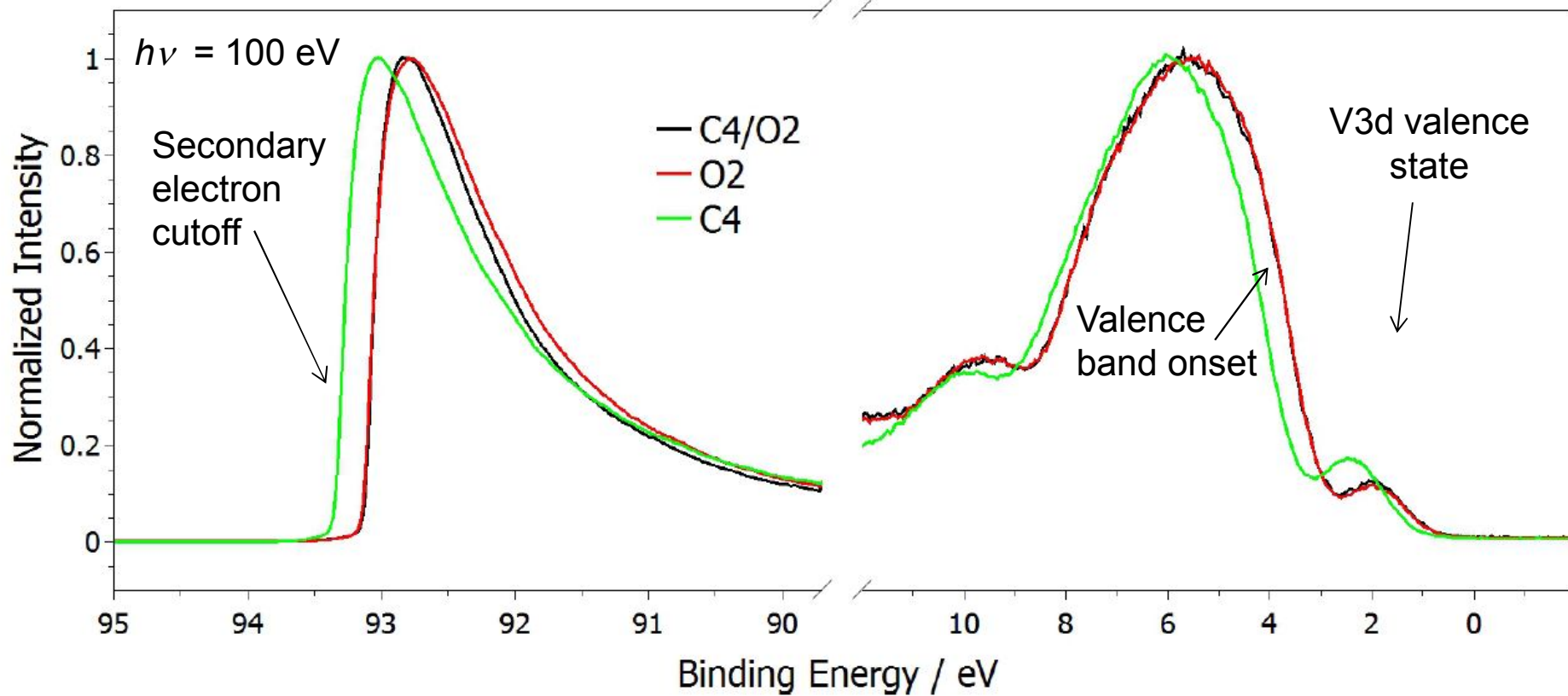
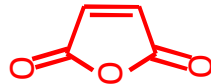
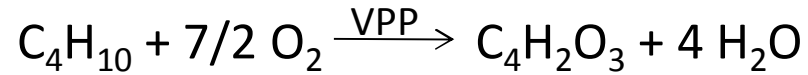


$$\text{Work function } \Phi = h\nu - E_{cutoff}$$

3. Surface states and space charge

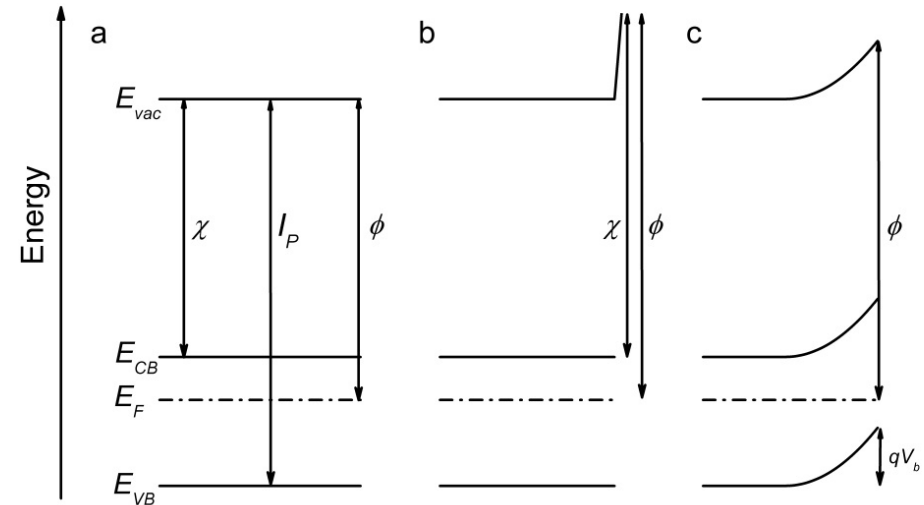
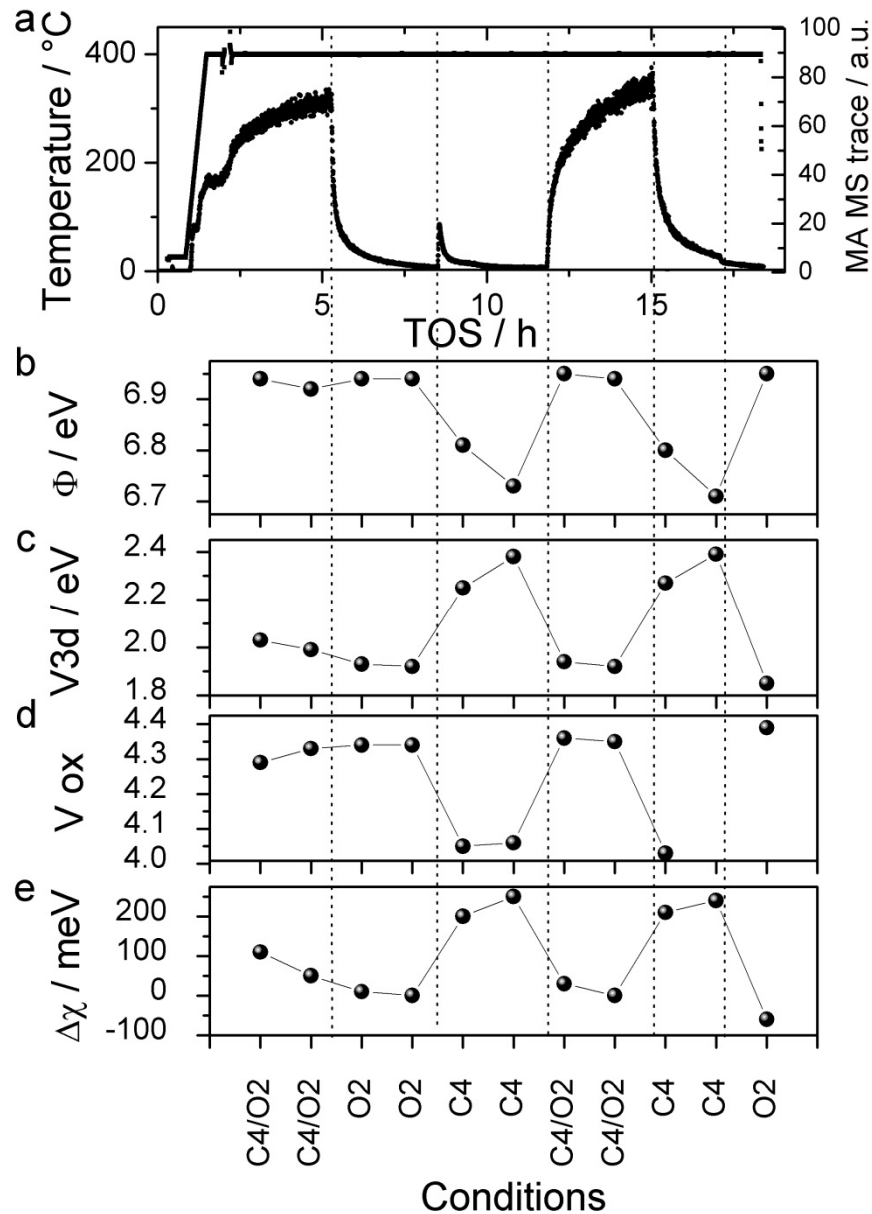
n-butane

maleic anhydride



$$\text{Work function } \Phi = h\nu - E_{\text{cutoff}}$$

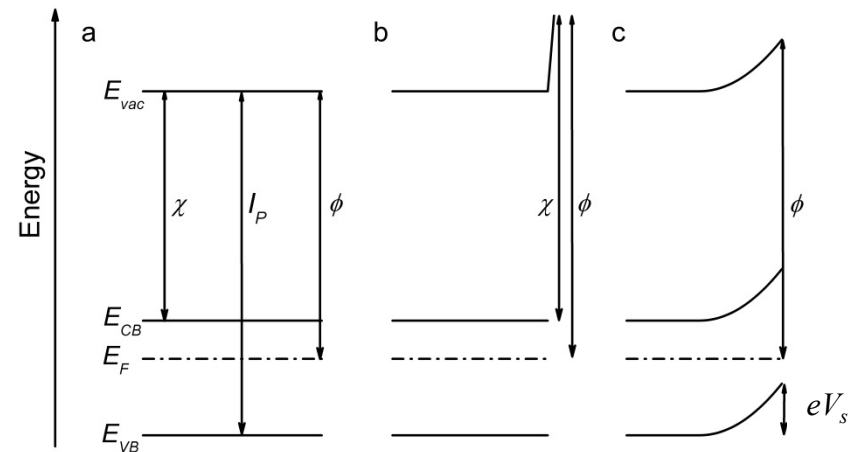
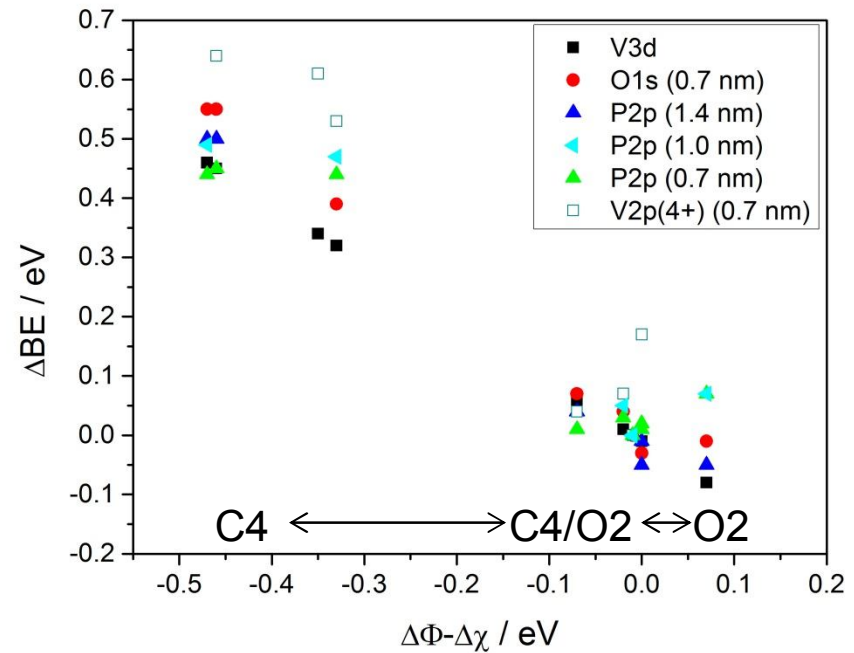
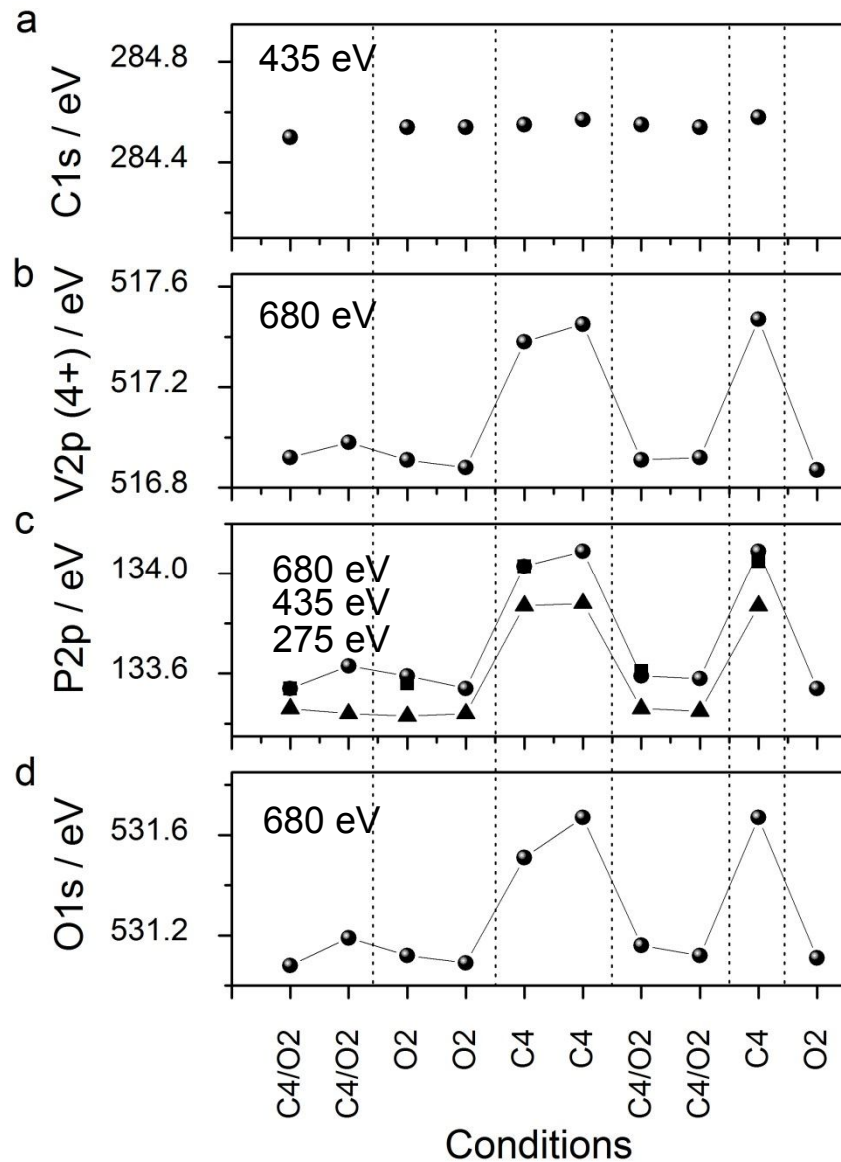
3. Surface states and space charge



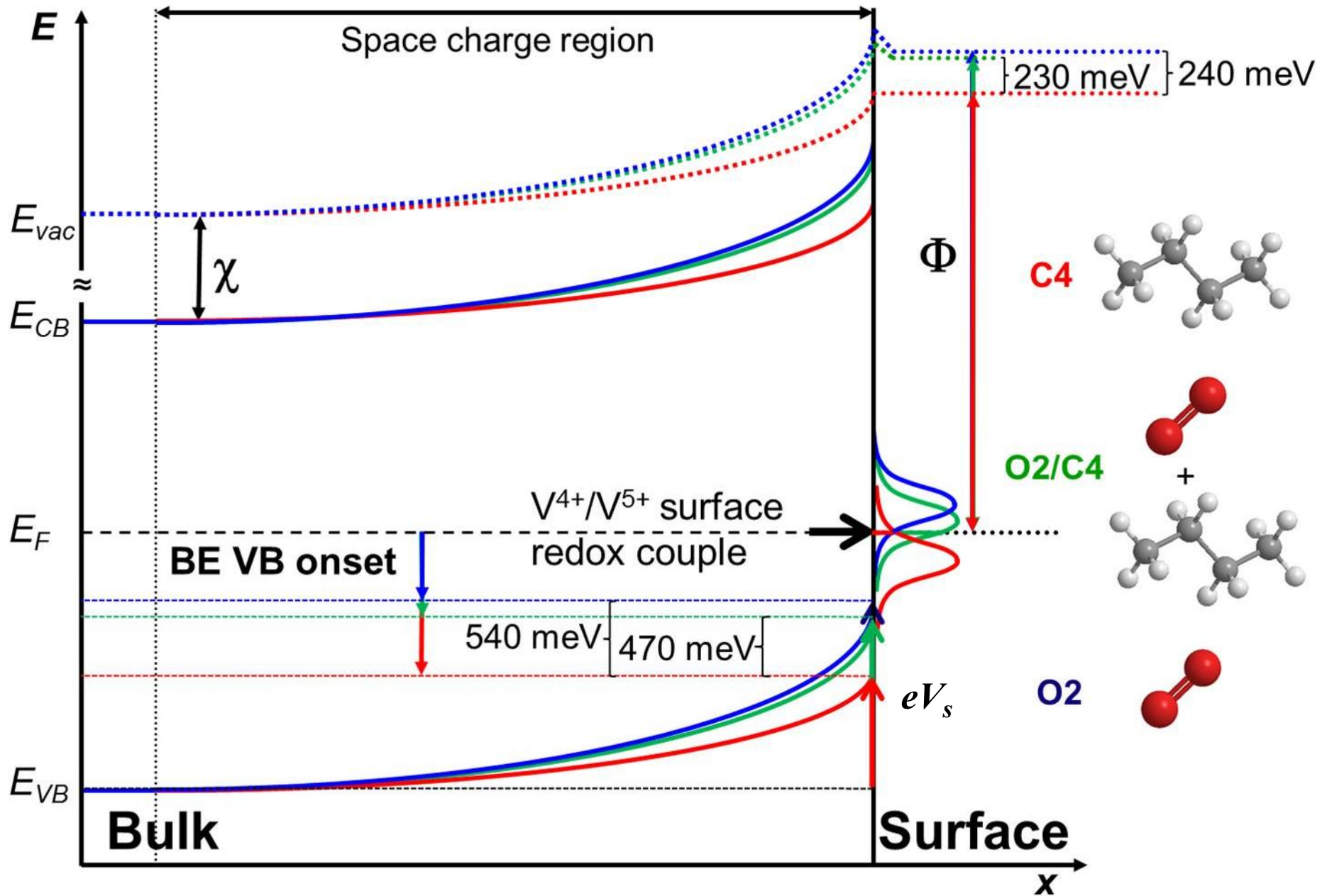
Work function $\Phi = h\nu - E_{cutoff}$
 O2 \rightarrow C4: $\Delta\Phi = -240$ meV
 $\Delta BE(V_{3d}) = 540$ meV

Electron affinity change $\Delta\chi = \Delta\Phi + \Delta BE(V_{3d})$
 $= +300$ meV

3. Surface states and space charge

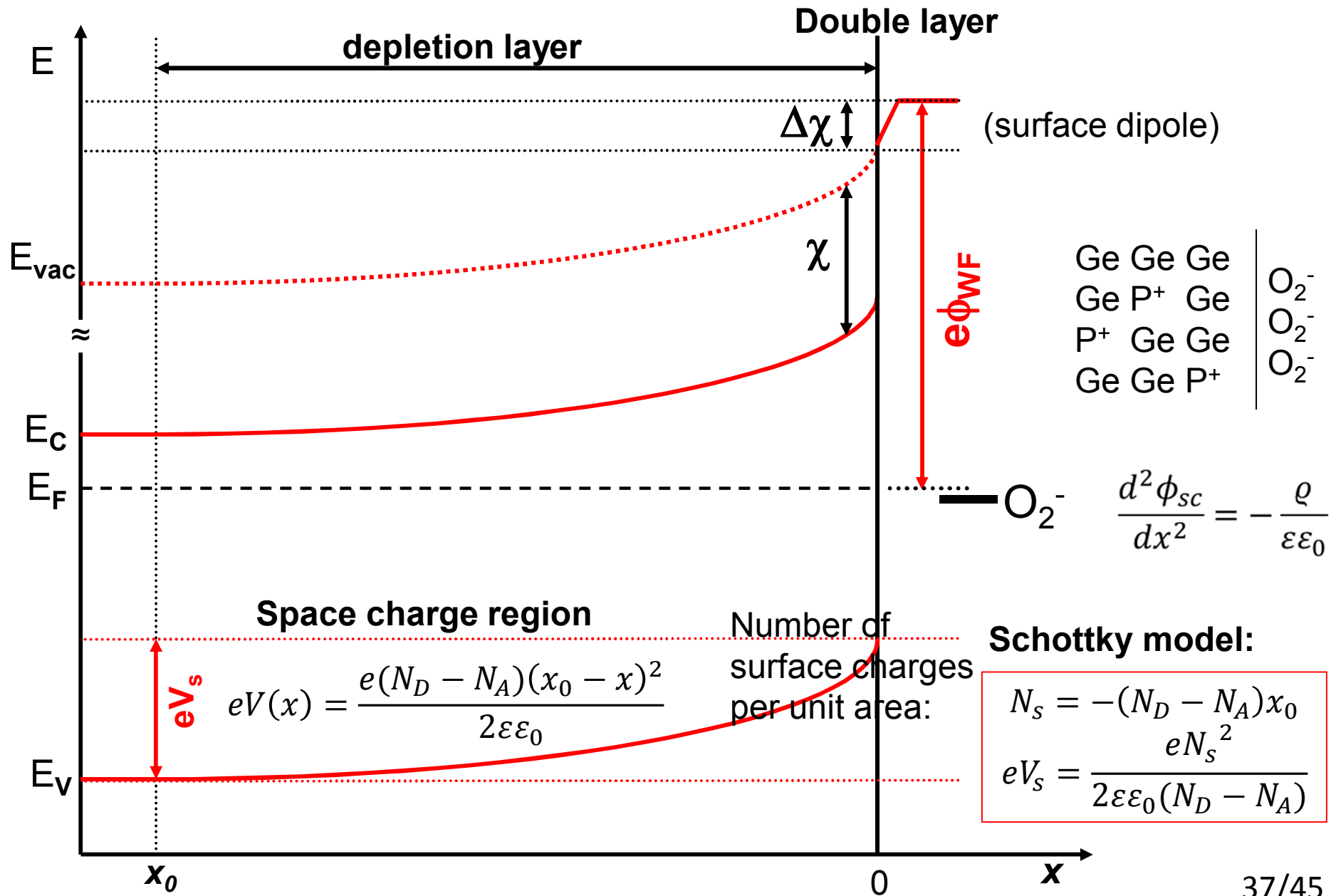


3. Surface states and space charge



Electron and hole transfer between bulk and surface: The consequences of space charge regions for adsorption and catalysis

4. Bulk-surface charge transfer



4. Bulk-surface charge transfer

Schottky model (example calculation):

$V_s = 1$ volt (typical surface barrier for adsorption)

Dielectric constant $\epsilon = 8$

Very pure material: $N_D - N_A = 10^{20} \text{ m}^{-3}$

$\rightarrow N_s = 3 \times 10^{14} \text{ m}^{-2}$; **1.5×10^{-5} monolayers** (e.g. of O_2^-)

Heavily doped semiconductor: $N_D - N_A = 10^{25} \text{ m}^{-3}$

$\rightarrow N_s = 1 \times 10^{17} \text{ m}^{-2}$; **5×10^{-3} monolayers** (e.g. of O_2^-)

Schottky relation:

$$eV_s = \frac{eN_s^2}{2\epsilon\epsilon_0(N_D - N_A)}$$

Depletion layer limits surface coverage to about $10^{-3} - 10^{-2}$ monolayers of equilibrium ionosorption! (Weisz limitation)

4. Bulk-surface charge transfer

Fermi energy pinning by redox couples

Surface state additives can be used to control the surface barrier (pin the Fermi energy) of a semiconductor or insulator

e.g. $10^{18}/\text{m}^2$ (0.1 monolayer) redox couple deposition (e.g. 1 FeCl_3 / 1 FeCl_2) on n-type semiconductor

→ Weisz limitation: max. $10^{16}/\text{m}^2$ charges can be transferred to surface; only 1% of deposited surface states would change their electron occupancy

→ e.g. 3×10^{-4} monolayers converted from Fe^{3+} to Fe^{2+} induces change of surface barrier by 0.42 V

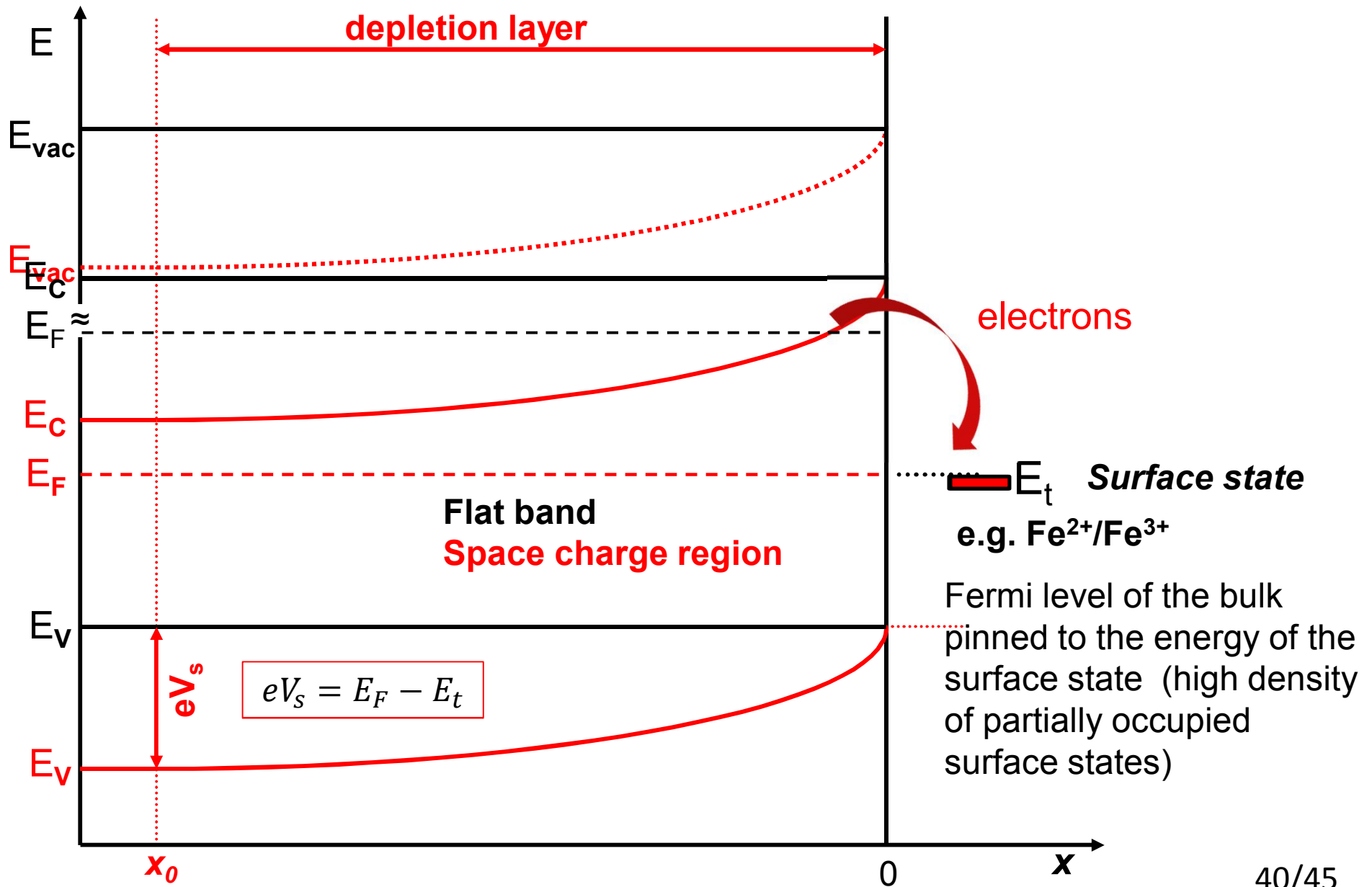
$$E_F = E_t + kT \ln \frac{[\text{Red}]}{[\text{Ox}]} \approx E_t$$

Fermi energy E_F is firmly pinned to surface state energy E_t

First approximation (ignoring local surface bonding of the adsorbate at the gas/solid interface) : $E_t \sim$ redox potential of the couple

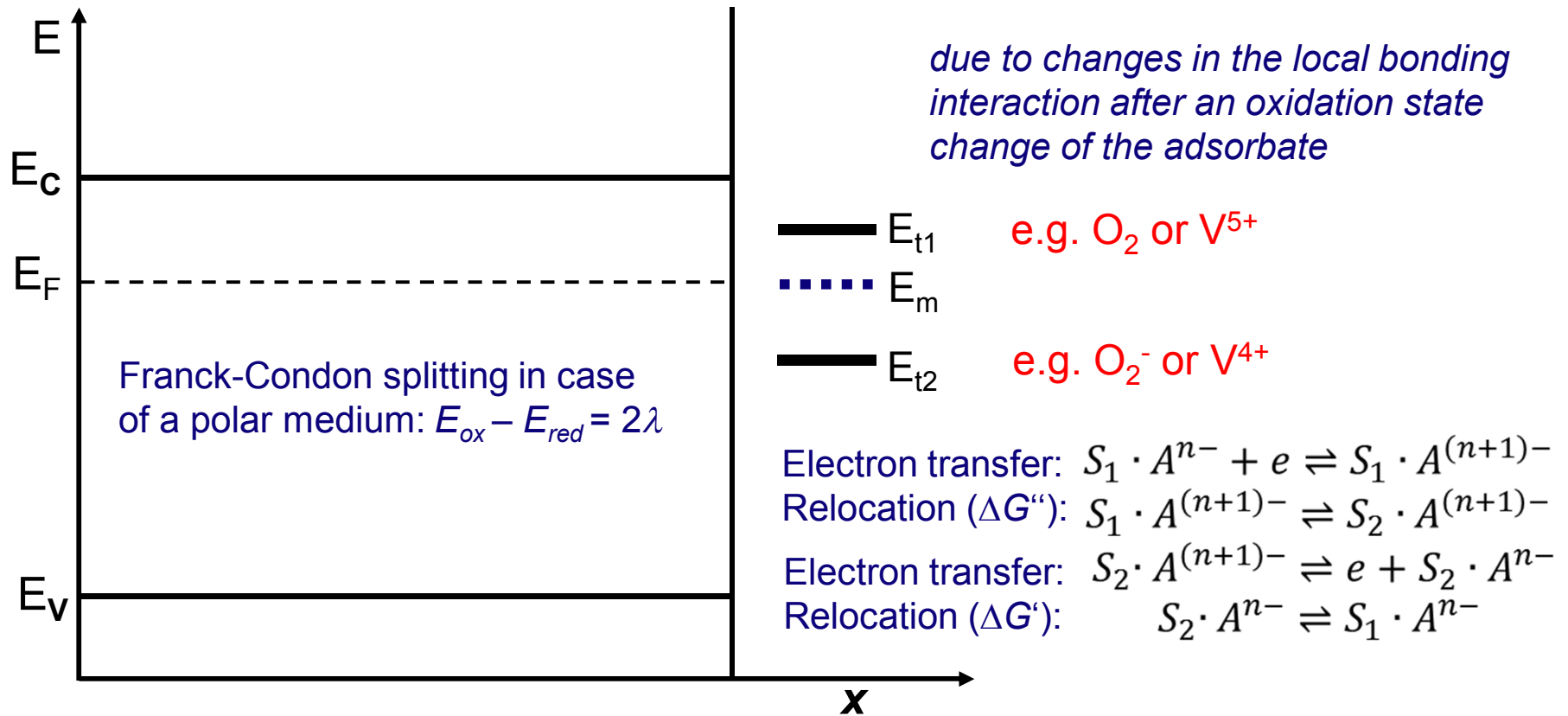
If gases react with redox couple, e.g. $\text{O}_2 \rightarrow$ adsorption/desorption of O_2 will control the surface Fermi energy

4. Bulk-surface charge transfer



4. Bulk-surface charge transfer

Frank-Condon splitting of unoccupied (E_{t1}) and occupied (E_{t2}) surface states



$$E_{t1} - E_{t2} = -(\Delta G'' + \Delta G')$$

$$E_F = \frac{1}{2}(E_{t1} + E_{t2}) + \frac{1}{2}(\Delta G'' - \Delta G') + kT \ln \left(\frac{[S_2 \cdot A^{(n+1)-}]}{[S_1 \cdot A^{n-}]} \right) \leftarrow \frac{\text{reduced species}}{\text{oxidized species}}$$

$$E_m = \frac{1}{2}(E_{t1} + E_{t2}) + \frac{1}{2}(\Delta G'' - \Delta G') \quad (\text{intermediate energy for an apparent single energy level})$$

4. Bulk-surface charge transfer

**Role of space charge layer in controlling electron transfer
→ charge transport kinetics**

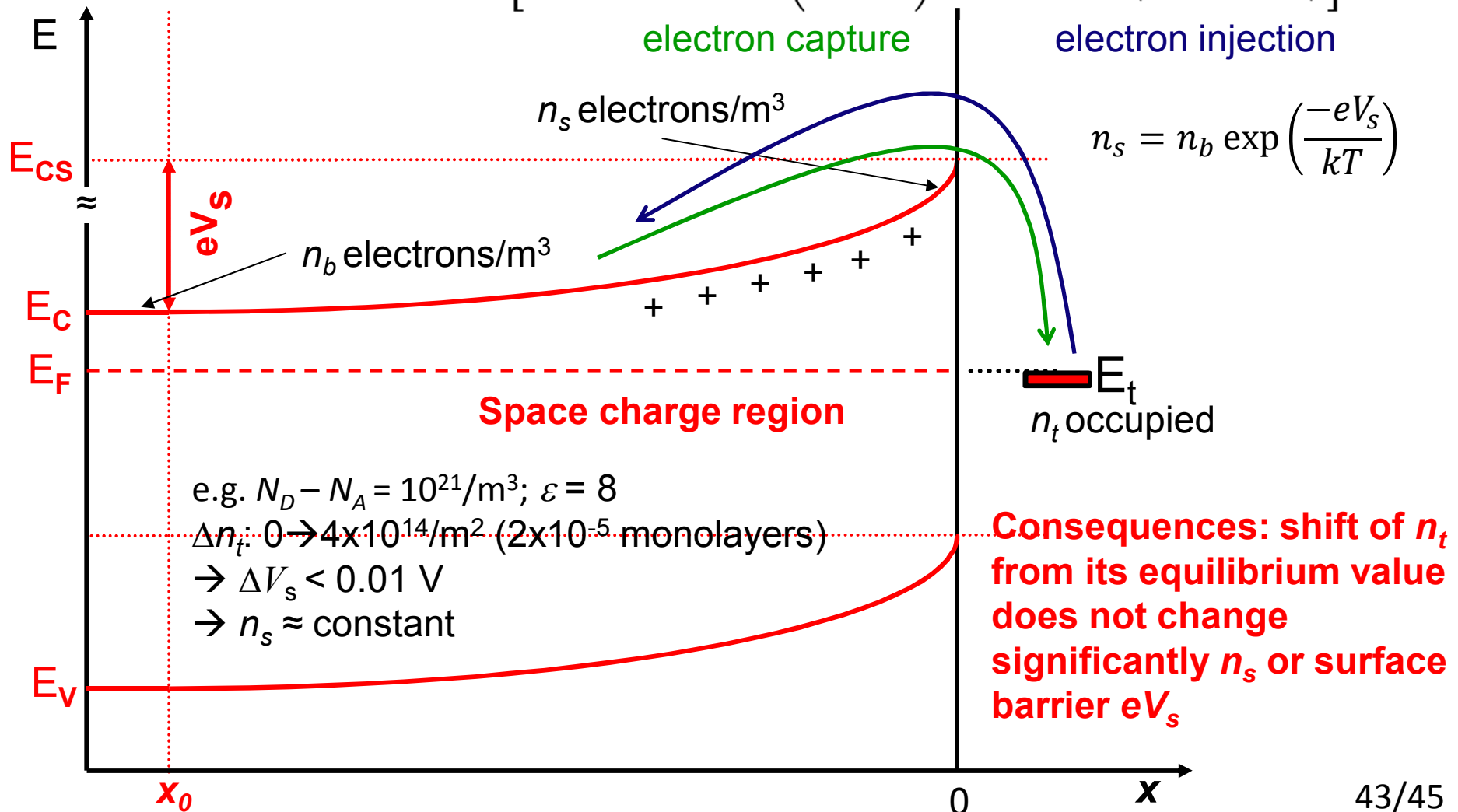
3 models:

- 1) Low and constant surface states density (surface free of volatile species)
- 2) High but constant surface states density
- 3) Energy level of surface states fluctuates with time (Franck-Condon effects, Gerischer model)

4. Bulk-surface charge transfer

1) Low and constant surface states density N_t low ($< 10^{15}/\text{m}^2$)

$$\frac{dn_t}{dt} = K_n \left[(N_t - n_t)n_b \exp\left(-\frac{eV_s}{kT}\right) - n_t N_c \exp\left(-\frac{E_{cs} - E_t}{kT}\right) \right]$$



4. Bulk-surface charge transfer

2) High but constant surface states density N_t very high ($> 10^{15}/\text{m}^2$)

$$\frac{dn_t}{dt} = K_n \left[(N_t - n_t)n_b \exp\left(-\frac{e V_s}{kT}\right) - n_t N_c \exp\left(-\frac{E_{cs} - E_t}{kT}\right) \right]$$

$$\rightarrow \frac{dn_t}{dt} = B' \left[\exp\left(-\frac{e V_s}{kT}\right) - \exp\left(-\frac{e V_{s0}}{kT}\right) \right]$$

V_{s0} : surface barrier at equilibrium
 B' : collection of constants

($N_t - n_t$ and n_t less sensitive to n_t than $\exp(-eV_s/kT)$)

$$\frac{dn_t}{dt} = B \left[\exp\left(\frac{-e\Delta V_s}{kT}\right) - 1 \right]$$

Rate of return
to equilibrium
for Δn_t :

$$\frac{d\Delta n_t}{dt} = B [\exp(-\beta \Delta n_t) - 1]$$

Rate decreases exponentially
with Δn_t

with $B = B' \exp(-eV_{s0}/kT)$
and β : constant

Consequences: a small change in surface state occupation induces a large increase of the surface barrier; electron transfer can be slowed to a negligible rate!

6. Literature

(Examples)

Adsorption/catalysis on semiconductors, e.g.:

- ✗ S. R. Morrison, "*The chemical physics of surfaces*", Plenum Press New York and London **1977** (semiconductor physics concepts of surfaces, includes chapter on heterogeneous catalysis)
- ✗ F. F. Volkenshtein, "*The electronic theory of catalysis on semiconductors*"; Pergamon Press **1963**.
- ✗ H. Lueth, "*Solid surfaces, interfaces and thin films*", Springer **2010**
- ✗ H. Lueth, "*Space Charge Layer*", Springer Verlag **2001**
- ✗ P. A. Cox, "*The electronic structure and chemistry of solids*", Oxford University Press **1989**
- ✗ R. Hoffmann, "*Solids and surfaces : a chemist's view of bonding in extended structures*" VCH **1988**

Fundamental textbooks, e.g.:

- ✗ N. W. Ashcroft, N. D. Mermin, "*Solid state physics*", Brooks/Cole Cengage Learning **2009**
- ✗ H. Ibach, H. Lueth, "*Solid-state physics : an introduction to principles of materials science*", Springer **2009**